HP-67 CALCULATOR PROGRAMS FOR THERMODYNAMIC DATA AND PHASE DIAGRAM CALCULATIONS

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May 25, 1978

Prepared for the U. S. Department of Energy under Contract W-7405-ENG-48



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LBL-5485

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*This report was done with support from the United States Energy Research and Development Administration. Any conclusions or opinions expressed in this report represent solely those of the author(s) and not necessarily those of The Regents of the University of California, the Lawrence Berkeley Laboratory or the United States Energy Research and Development Administration.

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THERMODYNAMIC HP67 CALCULATOR PROGRAMS

Introduction

Thermodynamic data and phase equilibria are greatly needed to meet current materials problems. Specific design criteria often require rather unique combinations of materials. Of the many combinations that might be of interest, data are available for only a small fraction. Useful combinations are often overlooked or great expense is incurred because an unnecessarily large number of combinations must be tested. The calculation of thermodynamic data and phase diagrams, even when no experimental data exist, is possible using spectroscopic levels of the gaseous atoms and chemical bonding theory (1-11). This was demonstrated for a large number of transition metal binary and multicomponent systems for which little or no data were available (2). Because of limitations of space for the large number of binary and multicomponent systems presented, projections of phase diagrams along the temperature axis were presented yielding only maximum extent of phase boundaries at the optimum temperatures although the text did discuss temperatures coefficients of the boundaries. The results of the calculations as a function of temperature will be presented in detail in subsequent publications.

The present report is a supplement to a tabulation of the thermodynamic and phase data for the 100 binary systems of Mo with the elements from H to Lr. The calculations of thermodynamic data and phase equilibria were carried out from 5000 K to low temperatures at which attainment of phase equilibria would not be practical. The results of these calculations are to be published as a special issue of the Atomic Energy Review, a publication of the International Atomic Energy Agency, Vienna. As only the results of the calculations are presented in the Atomic Energy Review issue, the present report is being used to present the methods of calculation used.

The thermodynamics involved is rather straightforward and the reader is referred to any of the advanced thermodynamic texts, e.g. Ref.(12), for derivations of the thermodynamic relations. The procedures for using chemical bonding theory to provide estimates of thermodynamic data have been outlined in previous publications (1-12).

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The calculations were largely carried out using an HP-65 programmable calculator and the programs have been tabulated (13). In this report, those programs have been put in form for use with the HP-67 calculator which allows for a great reduction in the number of programs required to carry out the calculations.

It is difficult in a compendium of calculation procedures of this type to present topics in a coherent manner, but an effort has been made to assemble programs into related groups. As the equations which serve as the basis for the calculations are discussed in detail in Ref.(13), the discussion of the background equations will be limited in this report.

CHAPTER I

Data Fitting Using the Chebyshev Polynomials

A. <u>Least-Square Fitting Using Chebyshev Orthogonal Polynomials</u>

The Chebyshev (Tschebycheff) polynomials, $T_n(x) = \cos(n\cos^{-1}x)$, are orthogonal over the continuous interval $0 \le x \le 1$ and they have been shown to be the most economical polynomial for expressing f(x) as a polynomial series with the minimum number of terms for a given accuracy.

The Chebyshev polynomial can be modified (14,15) to $C_n(\bar{x})$ which is orthogonal for discrete integer values of the variable, \bar{x} , from 0 to N with

$$C_{O} = 1 \text{ and } C_{n} = n! \Delta^{n} \left[\left(\frac{\overline{x}}{n} \right) \left(\overline{x+1-N} \right) \right] = \sum_{m=0}^{n} (-1)^{m} {n \choose m} {n+m \choose m} \frac{\overline{x!(N-m)!}}{(\overline{x}-m)!N!}$$
$$= \sum_{m=0}^{n} (-1)^{m} \frac{(n+m)!}{(n-m)!(m!)^{2}} \frac{\overline{x!(N-m)!}}{(\overline{x}-m)!N!} \cdot$$

The recurrence relation is

$$C_{n+1} = [(2n+1)(N-2\bar{x})C_n - n(N+n+1)C_{n-1}]/(n+1)(N-n)$$

$$C_0 = 1$$

$$C_1 = 1 - 2\bar{x}/N$$

$$C_2 = 1 + 6\bar{x}(\bar{x}-N)/N(N-1)$$

For N = 7, for example,

$$C_0 = 1, C_1 = 1 - \frac{2}{7}(\bar{x}), C_2 = 1 - \bar{x} + \frac{1}{7}(\bar{x})^2 C_3 = 1 - \frac{55}{21}(\bar{x}) + (\bar{x})^2 - \frac{2}{21}(\bar{x})^3.$$

The Chebyshev polynomials for discrete integer values are particularly useful for least square fitting of experimental data. Because of the orthogonality, cross terms are zero in the matrix used to solve the set of linear equations obtained by setting the partial derivatives of the squares of the deviations equal to zero.

Thus, the coefficients of $f(\bar{x}) = c_0 C_0(\bar{x}) + c_1 C_1(\bar{x}) + c_2 C_2(\bar{x}) + c_3 C_3(\bar{x}) + \dots$ are readily calculated without solution of a matrix by the relation $c_n = (f, C_n)/(C_{n,x}C_n)$ where

$$(f,C_n) = \sum_{x=0}^{N} f(x)C_n(\bar{x}) \text{ and } (C_n,C_n) = \sum_{x=0}^{N} [C_n(\bar{x})]^2 = \frac{(N+n+1)!(N-n)!}{(2n+1)(N!)^2}$$

For N+1 values of f(x) at N evenly spaced intervals of x, the values of $C_n(\overline{x})$ to be multiplied by f(x) are evaluated by the following program.

$$c_{n+1} = \frac{(2n+1)(N-2\bar{x})c_n - n(N+n+1)c_{n-1}}{(n+1)(N-n)}$$

$$c_0 = 1 \quad \text{and} \quad c_1 = (N-2\bar{x})/N$$

$$c_2 = \frac{3(N-2\bar{x})c_1 - (N+2)}{2(N-1)}$$

$$c_3 = \frac{5(N-2\bar{x})c_2 - 2(N+3)c_1}{3(N-2)}, \quad \text{etc.}$$

The integers, \bar{x} , which range from 0 to N, are related to x by $\bar{x} = (x-x_i)/I = x/I - r$, where x_i is the initial value of x and I is the interval spacing of x.

The values of $C_n(\bar{x})$ are stored in the registers starting with $C_1(1)$, $C_2(1) \cdots C_{n_{max}}(1)$, $C_1(2)$, $C_2(2) \cdots C_{n_{max}}(2)$, $C_1(3) \ldots$ In each cycle of calculations of $C_n(x)$, n is incremented until $C_{n_{max}}(x)$ has been calculated, then x is incremented and the cycle started again at n = 1. The calculation stops when x exceed N/2, as the second half of the values of $C_n(x)$ are the same as the first half except for change of sign for odd values of n. For n' = 3, there is sufficient storage to calculate the $C_n(x)$ values up to N = 15. For n' = 4, the maximum is N = 11. However, if storage is to be provided for subsequent (f, C_n) calculations, the maximum N values for n'=3 and 4 are 13 and 9, respectively.

n', the maximum value of n or the order of the polynomial used for fitting, is keyed in first. With N evenly spaced intervals of x, the number of data points to be treated, N+1, is keyed in next.

Directions:

(1) Insert tape Cheb C_n

(2) n' ↑ N+1 A 24 displayed at end of calculation The values of C_n(x̄) calculated in the order of increasing n for each value of x̄ and in order of increasing x̄ are stored in the registers in reverse order from index number 24 to 7. If it is desired to examine the C_n(x̄) values, R/S following the calculations of step 2 will display 24 C₁(1), 23 C₂(1), 22 C₃(1), 21 C₁(2), etc. Only the contents of registers with non-zero values are displayed. When no non-zero values remain, 6 followed by 24 will be displayed (See Ref. 23).

If it is desired to preserve the $C_n(\bar{x})$ values for future (f,C_n) calculations, key f W/DATA and insert data card. Values of $C_n(\bar{x})$ are also tabulated for several values of N and n in Ref.(13) and could be keyed in for calculation of (f,C_n) values by Prgm. Cheb (f,C_n) .

Prgm	. Cheb	Cn								and a support of the	
(2)	fLBLA	l	-	ST05	hR↓].		STOG	2	5	
	hSTI	0	STO4							13	
	fLBLO	RCL5	2	÷	RCL4	1	4	gx>y	GTO2	STO4	
	1	STO3	RCL4	2	Х	RCL5	÷	CHS	1	÷	
	fDSZ	STO(i)								35	
	fLBL1	RCL4	2	Х	RCL5		Х	RCL3	2	Х	
	1	+	Х	hx ≵y	RCL3	Х	RCL5	1	+	RCL3	
	+	Х	+	RCL3	RCL5		*	RCL3	1	+	
		fDSZ	STO(i)rcl6	RCL3	1	+	gx>y	GTOO	STO3	
	hR↓	hR↓	fISZ	RCL(i)	hx≉y	fDSZ	GTO1			82	
	fLBL2	2	STO÷5	gLBLfe	RCL5	fINT	STO4	2	4	hSTI	
	hRTN	6	STOG	fLBLC	RCL(i) fx≠0	GTOD	fDSZ	GTOC	fLBL D	
	hRCI	DSPO	hPSE	RCL(i)	gx=y	GTOfe	DSP2	f-x-	fDSZ	GTOC 112	
R	3	4		5	6	I		Valu	es of C _n	(\bar{x}) are	
	n=l	x=0		N	n'-1	Inde	ЭX	stor	ed in re	verse	
	to to N _N		N	N				order from E to A,			
	n'-1	int	$\frac{\mathbb{N}}{2}$	$\frac{N}{2}$				S9 t	o SO, an	d 9 to 7.	

Test:

(2) $n'=3 \uparrow No.$ of data pts.=14 A, 24 after 40 sec.

R/S 24, $C_1(1) = 0.85$; 23, $C_2(1) = 0.54$; 22, $C_3(1) = 0.08$; 21, $C_1(2) = 0.69$; 20, $C_2(2) = 0.15$; 19, $C_3(2) = -0.46$; 18, $C_1(3) = 0.54$; 17, $C_2(3) = -0.15$; 16, $C_3(3) = -0.69$; 15, $C_1(4) = 0.38$; 14, $C_2(4) = -0.38$; 13, $C_3(4) = -0.66$; 12, $C_1(5) = 0.23$; 11, $C_2(5) = -0.54$; 10, $C_3(5) = -0.47$; 9, $C_1(6) = 0.08$; 8, $C_2(6) = -0.62$; 7, $C_3(6) = -0.17$; 6, 24. The least square fit to the starting data with terms ranging in n up to the maximum value, n', yields the equation $f(x) = \sum_{n=0}^{n} [(f,C_n)/(C_n,C_n)] C_n(\bar{x})$, which has the advantage over equations fit by other least-square procedures in not weighting the different powers unequally. Unless weighting factors are added, all points will have equal weight. The equation can be expanded in powers of \bar{x} by substitution of

$$C_{n}(\bar{x}) = \sum_{m'=0}^{n} (-1)^{m'} \frac{(n+m')! \bar{x}! (N-m')!}{(n-m')! (m'!)^{2} (\bar{x}-m')! N!}$$

Each $C_n(\vec{x})$ will provide n terms in powers of \vec{x} ranging from $(\vec{x})^{\circ}$ to $(\vec{x})^{m'}$. Substitution of $\vec{x} = x/I - r$ will then convert each term to m' terms in x^m ranging from x° to $x^{n'}$. The final equation will be a polynomial $f(x) = \sum_{m=0}^{\infty} a_m x_{\cdots n}^m$. Each of these coefficients relate to the original (f, C_n) terms by $a_m = \sum_{m=0}^{\infty} b_{mn}(f, C_n)$. For n' = 3, there will be ten b_{mn} coefficients. For n' = 4, there will be fifteen b_{mn} coefficients to be evaluated, multiplied by the appropriate (f, C_n) values and summed to obtain the a_m values.

The \mathbf{b}_{mn} coefficients are evaluated by the equation

$$b_{mn} = \frac{1}{(C_n, C_n)} \frac{1}{(-1)^m} \sum_{m'=m}^{n} A_{m'n} R_{m'm}$$
where $(C_n, C_n) = \frac{(N+n+1)!(N-n)!}{(2n+1)(N!)^2}$
and $A_{m'n} = \frac{(n+m')!(N-m')!}{(n-m')!(n-m')!(N-m')!}$.
For $m = n$, $m'=m=n$, $R = 1$ and $A = \frac{(2n)!(N-n)!}{(m!)^2N!}$.
Thus $b_{mn} = \frac{(2n!)^2(N-n)!}{(n!)^2(N!)(-1)^n(C_n, C_n)}$ for $m=n$.
For $m = n-1$, $b_{n-1,n} = (b_{nn})(-\ln)(N/2 + r)$
For $m = 0$, $R_{m'0} = r(r+1) \dots (r + m' - 1)$ and $b_{0n} = \frac{1}{(C_n, C_n)} \sum_{m=0}^n A_{m'n}R_{m'0}$.
For $m = 1, m' = 3$, $R_{m'm} = 3r^2 + 6r + 2$.
For $m = 2, m' = 4$, $R_{m'm} = 6r^2 + 18r + 11$.

Values of b₀₀, b₀₁, b₁₁, b₀₂, b₁₂, b₂₂, b₀₃, b₁₃, b₂₃, and b₃₃ are tabulated for n=3 and N = 4, 7 and $\mathbf{8}$, and 9 and equations as a function of N are given in Table I-5 of Ref.(13).

When values of f(x) are given for x values that are evenly spaced at N intervals of magnitude I starting with an initial value x_i , Prgm. Cheb(f,C_n) will convert to variable $\bar{x} = x/I - x_i/I = x/I - r$ with \bar{x} varying by integral values from 0 to N and then will least-square fit $f(\bar{x})$ to a Chebyshev polymomial of order n' using the b_{mn} coefficients of Prgm. Cheb C_n corresponding to n' and N as discussed in Ref.(13). If followed by Prgm. Cheb-a, the Chebyshev polynomial is expressed as $f(x) = \sum_{n=1}^{\infty} a_n x^n$ upon insertion of I and $r = x_i/I$. The following program is designed for n' = 3, but can be adapted to n' = 4. i $\frac{N}{2}$ is integer value of $\frac{N}{2}$.

Directions:

(1)	Insert	Cheb	(f,C_n)	Prgm.
(2)	Insert	Cheb	Cn	Data Tape for $n'=3$ and appropriate N
(3)	f(0)	B		$\frac{\text{Display}}{f(0)}$
(4)	f(1) f(2)			f(1) f(2)
	•	•		
f	$(i\frac{N}{2}-1)$	•		$f(i\frac{N}{2}-l)$
	$(i\frac{N}{2})$	•		When N odd, $f(i\frac{N}{2})C_3(i\frac{N}{2})$; otherwise $\frac{N}{2}$
	$(i\frac{\mathbb{N}}{2}+1)$	•		$-f(i\frac{N}{2}+1)$
f	$(i\frac{N}{2} + 2)$	•		$-f(i\frac{N}{2}+2)$
	• • •	:		
f (5)f	(N-1) (N)	R/S E		-f(N-l) (f,C ₀) (f,C ₁)
<u>a</u> .		0		(f,C_2) (f,C_3)

Start at step 3 for each new set of data with same N. Step 5 must be keyed before starting again at step 3. For different N start at step 2.

Test:

- (2) Insert Cheb C_n Data tape for n'=3 and N+1=14 data points.
- (3) 4B; 15 R/S; 40 R/S; 85 R/S; 156 R/S; 259 R/S; after each of these steps, the input number used in the calculation is displayed;
 400 R/S 67.133; 585 R/S; 820 R/S; 1111 R/S; 1464 R/S; 1885 R/S;
 2380 R/S; after each of these steps, the negative of the input number is displayed;
- (5) 2955 E 12159, -7483, 2632, -408.

Prgn	n. Cheb(f,C _n)	*1.***1.2.4.4.45		a na mana singka ka ka ka ka sing a	3.9.5 Trings 3877					
Po LProcessory	gLBLfe	RCL5	fINT	STOT	2	λ_{\downarrow}	hSTI	hRTN			8
(3)	fLBLB	STOO	ST01	ST02	STO3	hRTN					14
(4)	gLBLfc	个	个	†	fGSB3	fGSB5	hx≵I	hR↑	hRTN	RCL4	
	fx=0	GTOfb	hR↓	GTOfe	fLBLfb	hR↓	GTOfd				31
	fLBL3	STO+0	fGSB4	STO+1	fGSB4	ST0+2	fGSB4	ST0+3	hRTN		40
	fLBL4	CLX	RCL(i)	fDSZ	Х	hRTN					46
	fLBL5)†	hx≵I	gDSZ(i) hRTN	fGSB6	hx≵l	hRTN			54
	fLBL6	hx ‡ I	fISZ	RCL5	RCL5	fINT	gx≠y	hR'I'N	hRCI	3	
	+	hSTI	hRTN								67
	gLBLfd	fGSB7	fGSB8	hR↑	hR1'N	GTOfd					73
(5)	fLBLE	fGSB7	STO+0	hR↓	ST0+1	hR↓	ST0+2	hR↓	ST0+3	gGSBf	e.
	RCLO	RCL1	RCL2	RCL3	gSTK	hRTN					89
	fLBL7	CHS	个	CHS	1	CHS	个	CHS	hRTN		98
	fLBL8	ST0+0	fGSB9	ST0+3	fGSB9	ST0+2	fGSB9	ST0+1	hRTN		107
	fLBL9	C.L.x.	RCL(i)	fISZ	Х	hRTN		113 m	emory	steps.	
R	0	1	2	3]			I			
	(f,C ₀) (f,C ₁) (f,C ₂) (f,(C ₃) In	.x. d.e.x.	<u>N</u> 2	C _n Inde	эх		
	C _n (x)	values	in E to	A, S9 t	to 0, ai	nd P9 -	to 7.				

Directions:

(1)	Insert Cheb-a tape
(2)	Insert (f,C ₀) to (f,C ₃) in registers $0-3$ if not carried over from Prgm. Cheb.(f,C _n).
(3)	1 † r A
(4a)	SST
(4b)	SST

- (4e) SST

(5) $x \to f(x); \to f(x+1); \to f(x+21), \text{ etc.}$

NOTE: After step 4c, the a_n values can be reviewed by keying another SST which will flash the four a_n values in the order a_0 to a_3 . The (f,C_n) values are retained in RO-3 and step 3 can be repeated after N/2 STO 5; for example, if incorrect values of I or r are used.

Display a0 a1

a2

az

Test:

- (2) 12159 STOO, 7483 CHS STO1; 2632 STO2; 408 CHS STO3; 6.5, STO5;
- (3) 200 \uparrow 1.5 A 0.625; (4) SST 0.00375; SST -1.25×10⁻⁵ SST 1.25×10⁻⁷
- (5) 300 B 4; E15, E40, E85; 1500 B 400, E585, E820.

Prgm	. Cheb-a		- <u></u>			******						
(3)	fLBLA	STOE	hR↓	CHS	STOC	RCL5	2	Х	STOD	3 X		
	STOA	STOB	fGSBl	RCL3	STOG	X	ST09	fGSB2	RCL3			
	STO8	2	STOA	STOB	fGSB1	RCL2	ST05	Х	STO+8	fGSB2	2	
	RCL2	Х	ST07	3	STOA	STOB	fGSBO	STOX6	RCLE	3		
	Х	6	+	RCLE	Х	2	+	fGSB3	2	ST01	3	
	RCLE	Х	1	+	fGSB3	+	1	STOB	fGSB3	-+-		
	RCLC	*	RCL6	hSTI	Х	ST0+7	1	STOA	fGSBl	RCL1		
	Х	ST0+7	fGSB2	RCL1	Х	ST06	RCLD	1	+	hl/x		
	RCLO	Х	ST0+6	2	STOA	fGSBO	STOX5	RCLE	fGSB3	1		
	+	2	STOB	hR↓	RCLE	1	+	RCLE	Х	STO1		
	fGSB3	+	RCL5	Х	sto+6	3	STOA	RCL4	fGSB3	1		
	STOB	+	RCLE	fGSB3	÷	3	STOB	hR↓	RCLE	2		
	+	RCL4	Х	fGSB3	+	hRCI	Х	STO+6	rcl6	hRTN		
(4)	RCL7	rcl8	RCL9	gSTK							134	
	fLBL2	fGSB1	RCLC	Х	RCLA	Х	RCLD	2		RCLE		
	+	Х	hRTN								147	
	fLBL1	fGSBO	fGSB3	RCLC	RCLA	hy^X	÷	hRTN			155	
	fLBLO	RCLD	hN!	gx^2	RCLD	RCLA	-	hN!	÷	RCLD		
	RCLA	+	l	+	hN!	÷	RCLA	2	X	1		
	+	Х	hRTN								178	
	fLBL3	RCLD	RCLB	-	hN!	Х	RCLD	hN!	÷	RCLA		
	RCLB	+	hN!	Х	RCLA	RCLB	-	hN!	*	RCLB		
<i>(</i>)	hN!	gx ²	÷. •	hRTN							202	
(5)		hR↑	RCLC		fLBLB	<u>↑</u>	<u>۲</u>	1	RCL9	Х		
	RCL8	+	Х	RCL7	+	Х	RCL6	+	hRTN		221	
D	<u></u>	-	<u>^</u>	2	1.				6	-	<u></u>	-
R	0	1	2	3	4		5		6	7	8	9
	(f,C_0)	(f,C_1)	(f,C_2)	(f, C_3)) $r^2 + r$		N Z	(f,C ₃)	aر	ap	az
		L		-)			$(\mathbf{f}, \mathbf{C}_2)$	(f,C	3)/(C3,C3)	-1-	han	J
						(f,C	5 2 (f,C ₂) 2)/(C ₂ ,C	(2)	a _o			
R	А	B (C D	E		I						
τı	ы			L								
	n	m' -:	I N	r	4	or 5						
					(f,Ca)/(C ₃ ,	C ₃)					
							~					

The three programs $Cheb-C_n$, $Cheb-(f,C_n)$, and Cheb-a can be combined in several ways. There is enough space on one tape to combine the first two programs for a total of 204 memory steps if the 8 steps of gLBLfe are not duplicated; thus there would be no need to put the $C_n(x)$ values on a data tape. However, if N+1, the number of data points, is frequently used, it would save time to have stored the $C_n(x)$ values on a tape rather than to recalculate each time. If a data tape is used for the $C_n(x)$ values, a portion of Prgm. Cheb-a can be added to the end of Prgm.Cheb-(f,C_n).

This would provide enough space for the following additions for calculation of df(x)/dx and $\int_{0}^{x} f(x)dx$ upon keying xC and xD, respectively, or following xB with R/S to obtain df(x)/dx and another R/S to obtain the integral.

	hR↑					3			RCL8	2
	Х	÷	Х	RCL7	+	hRTN	hR↑			
	fLBLD	↑	↑	RCL9	4	÷.	X	RCL8	3	utern F
	+	Х	RCL7	2	*	÷	Х	RCL6	÷	X
Biancongon and services	hRTN						an anana an a			

It is unusual to have thermodynamic data of sufficient accuracy to warrant more than the four constant fit of Prgm. Cheb-a. Prgm. Cheb- C_n already can accommodate quartic or higher power fits. Prgm. Cheb- (f,C_n) and Cheb-a could be expanded to a quartic fit if desired.

For most applications of Prgms. Cheb C_n , Cheb (f,C_n) , and Cheb-a, the same number of data points, N+1, would not be repeatedly used and there would be no need to store the constants from the Cheb C_n calculations. Then Cheb C_n and Cheb (f,C_n) would be combined on the same tape and the instructions would be simplified to the following:

(1) Insert Cheb C_n -Cheb(f, C_n) tape	Display
(2) n'=3 ↑ N+l A	24
(3) f(0) B	f(0)
(4) f(1) R/S	f(1)
f(2) R/S	f(2)
• •	
f(N-1) R/S	f(N1)
(5) f(N) E	$(f,c_0),(f,c_1),(f,c_2),(f,c_3)$
(6) Insert Cheb-a tape	
(7) I † r A	^a 0
(8) SST SST SST	al,a2,a3
(9) x $B \rightarrow f(x)$; $E \rightarrow f(x+1)$; $E \rightarrow f(x+1)$	21), etc.
TEST: (Continuation of Prgm. Cheb C_n	test at top of pg. 0)

(7) 10 + 1 A 1; (8) SST .1, SST .01, SST .001; (9) 10 B 4, E15, E40.

Chapter II

Fitting of Experimental Data by An Analytical Function

Two types of operations will be considered. The first type requires a very accurate fit in a limited range of x for interpolation purposes, with no smoothing of the data. The second type smooths the data to fit some predetermined functional form by a least-square procedure. The first type will be discussed first and illustrated with some examples.

A. <u>Interpolation Fit to $f(x) = \Sigma a_n x^n$ </u>

Program Interp. GK fits four (x,y) pairs with the x values at evenly spaced intervals of magnitude I to a cubic polynomial, or any three pairs to a quadratic equation, or any two pairs to a linear equation, which can be used for interpolation purposes. The calculations are outlined in Ref.(13). In addition, the program is specifically designed to accept values of $-(G^{\circ}-H_{Std}^{\circ})/RT$ or $-(G^{\circ}-H_{Std}^{\circ})/T$ for each of the reactants and products of a chemical reaction at two, three or four temperatures and fit the resulting $-(\Delta G^{\circ}-\Delta H_{Std}^{\circ})/RT$ values to an interpolation equation which can then be combined with ΔH_{Std}° or $\Delta H_{Std}^{\circ}/R$ for the reaction to calculate ln K or K, the equilibrium constant of the reaction, at desired temperatures.

Directions:

(1) Insert tape Interp.GK	
<u>2 Pt. Fit</u>	Display
(2a) $y_1 \uparrow y_2 A$	y1-y2
(3a) $x_1 \uparrow x_2 R/S$	ao
(4a) SST	al
(5) x E	У
<u>3 Pt. Fit</u>	
(2b) y ₁ ↑ y ₂ ↑ y ₃ B	у1-у5
(3b) $x_1 \uparrow x_2 \uparrow x_3 R/S$	a ₀
(4b) SST SST	a ₁ ,a ₂
(5) x E	У
<u>4 Pt. Fit</u>	
(2c) $y_1 + y_2 + y_3 + y_4$ C	I ³ a ₃
(3c) I ↑ x ₁ R/S	ao
(4c) SST SST SST	^a 1, ^a 2, ^a 3
(5) x E	У

For the reaction aA + bB = mM + nN + oO, the values of $g = -(G^{\circ}-H_{Std}^{\circ})/RT$ (a positive number) for the reactants and products are keyed in as follows:

D.

J1-J5

(21.)	fe to init	iate a comp	lete set of data	$\frac{\text{Display}}{0.00}$
(b)	gA↑aX	g _B ↑ЪХ	D	0.00
(c)	g _M ↑ mX	g _N ↑nX g	; ₀ ↑oX R/S	$-(\Delta G^{\circ}-\Delta H_{Std}^{\circ})/RT$, 0.00

Steps b and c are repeated at each temperature. If only one mole of a reactant or product appears in the equation, e.g. n=1, omit n X and merely key g_N \uparrow . After the values at each of the temperatures have been entered, the following steps are carried out depending upon whether a 2, 3 or 4 pt. fit is desired.

(d) f a	<u>Display</u> -(ΔG°-ΔHStd)/RT at 4th temperature or zero if less than 4 temperatures treated.
(e4) For 4 pt. fit: key C and then continue with steps 3c,4c,and 5.	I ³ a ₃
(e ¹ ₃) For 3 pt. fit: key f b	$-(\Delta G^{o}-\Delta H_{Std}^{o})/RT$ at third temperature
(e ₃ ") B	<u>д</u>]-д5
Continue with steps 3b, 4b, and 5.	
(e2) For 2 pt. fit: key f b twice.	$-(\Delta G^{o}-\Delta H_{\mathrm{Std}}^{o})/\mathrm{RT}$ at second temperature

(e½) A

Continue with steps 3a, 4a, and 5.

After the values of $-(\Delta G^{\circ} - \Delta H_{Std}^{\circ})/RT$ have been fit to either a linear, quadratic or cubic equation for interpolation purposes, the following steps yield values of lnK or K.

(6)	$-\Delta H_{std}/R$	STO 6	-AHStd/R
(7)	Τfd		ln K
(8)	SST		K

NOTE 1: A new set of data cannot be treated unless step d has been carried out followed by step a. The program will accommodate an equation with three different products or four products if there is only one mole of the last product value entered. There can be two different reactants or three if there is only one mole of the last reactant value entered. However, any number of products and reactants can be accommodated if + is keyed in after the second reactant entry and every subsequent one and if + is keyed in after the third product entry and every subsequent one.

- NOTE 2: The program can be used for $-(G^{\circ}-H^{\circ}_{Std})/T$ and ΔH°_{Std} as well as for the dimensionless quantities used to illustrate the displays, but step 7 will display R lnK instead of lnK and it must be divided by R before obtaining K by step 8. R in appropriate units can be stored in register D for use with step 7 or RCL D STO:0 STO:1 STO:2 STO:3 will convert the equation for $-(\Delta G^{\circ}-\Delta H^{\circ}_{Std})/T$ to the dimensionless $-(\Delta G^{\circ}-\Delta H^{\circ}_{Std})/RT$ form: so that it is unnecessary to divide by R each time step 7 is carried out. Of course, the appropriate $-\Delta H^{\circ}_{Std}$ or $-\Delta H^{\circ}_{Std}/R$ must be used.
- The values of $-(\Delta G^{\circ} \Delta H_{Std}^{\circ})/RT$ obtained at each temperature are NOTE 3: stored in the secondary registers starting with SO during the calculations following step d. Thus, a set of values at four temperatures is available for repeat fits using two, three or four of the temperatures. For each fit, after the first, key fPZS and go to step d and then continue with either step e4, e3, or e2 depending upon which fit is desired. Step d returns the values to the secondary registers and they are available for repeated use. Although the program will not fit more than values for four temperatures at a time, steps a-c can be carried out repeatedly for up to ten temperatures. Any two, three, or four of these values can be recalled and stored in the stack as one would do to initiate steps 2a, 2b, or 2c. After fPZS to safeguard the values for future use, one would then proceed with steps 2 to 4 for the particular fit chosen. Also as discussed in Note 4, this procedure requires keying of hSFO if steps a to c are to be used again. Steps a to c can also be used to evaluate $\Delta H_{Std}^{2}/R$ of the reaction from the enthalpies of formation of the reactants and products. One should use fPZS appropriately to insure that values of $-(\Delta G^{\circ}-\Delta H_{Std}^{\circ})/RT$ or values of a_n that are wanted for subsequent calculations are not erased by the ΔH calculations.
- NOTE 4: Flag 0 is used to control the index numbers that regulate the storing of values during steps a to c. The program sets FO when inserted and subsequent operation of the program maintain FO in the proper set or cleared position as needed. However, if the use of steps a to c is not subsequently followed by step d, the next operation of steps a to c will not start storage in RO but will store beyond the last registers used.

<u>Test</u>:

- (2b) 1.978 + 2.536 + 3.25 B 0.558;
- (3b) 0.3 + 0.4 + 0.5 R/S 1.240;
- (4b) SST 0.120 SST 7.800; (5) 0.4 E 2.536;
- (2c) 1.552 * 1.978 * 2.536 * 3.25 C 0.004;
- (3c) 0.1 [†] 0.2 R/S 1.000
- (4c) SST 2.000 SST 3.000 SST 4.000; (5) 0.4 E 2.536

cont'd.

 $C(gr) + 2Cl_2(g) = CCl_4(g)$ (a) fe 0.00 (b_1) 1.16 \uparrow 49.85 \uparrow 2 X D 0.00; (c_1) 68.1 R/S - 32.76, 0 500 (b₂) 2.78 \uparrow 55.43 \uparrow 2 X D 0.00; (c₂) 81.31 R/S - 32.33, 0 (b₃) 4.19 \uparrow 58.85 \uparrow 2 X D 0.00; (c₃) 90.01 R/S - 31.88, 0 (b₄) 5.38 \uparrow 61.34 \uparrow 2 X D 0.00; (c₄) 96.53 R/S - 31.53, 0 1000 1500 2000 K (d) f a - 31.530; (e4) C - 0.020; (3c) 500 ↑ 500 R/S - 33.050; (4c) SST 3.60×10^{-4} SST 5.20×10^{-7} SST - 1.60×10^{-10} ; (5) EEX 3 E - 32.330; 1.98719 STOD, $-\Delta H_{o}^{o}/R = 25 \times 10^{3} \text{ RCLD} \div = 12581 \text{ STOG}$ RCLD 1.98719 STO:0 STO:1 STO:2 STO:3 K 8.676 (7) 500 fd 5858 (8) SST 750 fd 0.392 1.479 2.50×10-2 -3.689 EEX3 fd 4.73×10-4 1500 fd -7.656 6.94×10-5 2EEX3 fd -9.576 fP≠S f a -31.530; f b - 31.880; (2b) B - 0.430; (3b) 500 ↑ EEX3 ↑ 1500 R/S - 33.170; (4b) SST 8.0×10⁻⁴ SST 4.0×10⁻⁸; RCLD 1.98719 STO÷0 STO÷1 STO÷2 (7) 500 fd 8.676 5858 750 fd 0.395 1.485 1500 fd -7.656 4.73×10-4 fP≠S fa - 31.530; fb - 31.880 fb - 32.330; (2a) A - 0.403; (3a) 500 \uparrow EEX 3 R/S - 33.190; (4a) SST 8.6×10⁻⁴ (5) EEX 3 E - 32.330; RCLD 1.98719 STO÷0 STO÷1 1.487 (7) 750 fd 0.397 1500 fd -7.666 4.69×10⁻⁴

									· •	
Prgm.	Interp.G	K								
(d)	gLBLfa	RCLO	RCL1	RCL2	RCL3	fP≵S	hSF0	hRTN		8
(e,f)	fLBLfb	hR↓	hR'IN			`				
(2a)	fLBLA	STO4		0	ST02	ST03	hR↓	hRTN		11
(3a)	ST05	_	•	ST01	RCL5	X	CHS	RCL4	÷	I 9 STOO
(4a)	hRTN	RCL1	•	DTOT	1011/	21	Ond	11011-	•	
				amor	Later		cmol		0	31
(2b)	fLBLB	hx≵y		ST05	hR↓	hLSTx	ST04		0	
())) () () () () () () () ()	ST00	STO3	hR↓	hRTN						44
(3b)	ST09	hR↓	STO8	hR↓	ST07	RCL8	-	<u>.</u>	RCL8	RCL9
	+	STOB	Х	RCL7	RCL8	+	RCL5	Х	RCL9	RCL8
		STOA	<u>.</u>	-	RCL9	RCL7		÷	ST01	RCL5
	RCLA	÷	_	CHS	RCLB	• •	ST02	rcl8	Ε	CHS
(4b)	RCL4	+	STOO	hRTN	RCL1	RCL2				9 0
(2c)	fLBLC	hR↑	STOC	_	ST07	hR↓	hLSTx	_	ST05	hR↓
(20)	hLSTx		ST04	RCL7	3	÷	RCL4	+	RCL5	
	2	÷.	hRTN	1(01)	<u> </u>	•	10017		101)	
(2-)				2	hy ^x		CIIIO O	ם זאם	DATA	113
(3c)	STOB	hR↓	STOA	3			ST03	RCLB	RCLA	+
	Х	3	X	CHS	RCL5	RCL4	2	X	-	2
	÷	RCLA	gx2	÷	+	ST02	RCL4	RCLA	÷	RCLA
	RCLB	2	Х	+	RCL2	X		RCLA	RCLB	+
	RCLB	Х	3	Х	RCLA	gx ²	+	RCL3	Х	-
	STO1	0	STOO	RCLB	E	CHS	RCLC	+	ST00	hRTN ₁₇₃
(4c)	RCL1	RCL2	RCL3							1/6
(7,8)	gLBLfd	Ε	rcló	hR↑	÷	+	hRTN	ge^X		184
(5)	fLBLE	± ↑	0 ↑	<u>↑</u>	RCL3	Х	RCL2	+	Х	RCLL
	+	x	RCLO	+	hRTN	21	110436-		43.	
()		0	10000	, 个	111(1.1V	hF?0	hSTI	hRTN		199
(a)	gLBLe		т +	T STO(i		hRTN	HOTT	TTU T IV		207
(Ъ)	fLBLD	+					aton	1 0770	e	2.1.3
(c)	+	+	+	RCL(1) —	STO(i)	f1SZ	hCFO	f-x-	fe
	hRTN							224 m	emory	steps.224
2 pt.	R O		2 3	4	5					
(2a-4a	a) a ₀	al	0 0	д5	x2					
			<u> </u>	4		6	r7	8		<u> </u>
3 pt.			2 3		5	0	7			9
(2b-4t	o) 0	a_{l}	a ₂ 0	<u>у</u> 5	у3 - д5		хJ	х	2	x ³
	a	*****						c		-
	А	В								
	x3-x5	x2+x	3							
), mat				4	5	7				
4 pt.			2 3		-	7				
(2c-4a			a ₂ a3	<u>д5-д</u> 1	<u>дз-д</u> т д	4-y1				
	A	В	C							
	I	хl	Уl							
R () to 3		4	5	6	7		8	9	
							10		9	
	ch a _n is			used	$\Delta H_{Std}^{o}/F$			0		
	Rn	Д2				уц-у	l or	x2	or x3	
R	A	В	С		I					
	x3-x2	x_2+x_3	0	I	ndex					
C	orI		or y	/1						
			-AHŠta),	KT for	2 to 4	tempera	tures i	in orde	er of t	increasing
	temperatu	re.								

B. Least-Square Fitting of Data to an Analytical Function

Least-square fitting of data to an equation y = f(x) is not a routine process but requires careful consideration of the variations of errors in y as a function of x. (15,16,17) For example, if it were desired to obtain the values of c and d in the expression $y = cx^2 + dx^3$ that best represent a set of data, one could least-square a variety of functions of y. The use of the unweighted function would tend to heavily weight values of y at large x. As just one alternative example, one could least-square $y/x^2 = c + dx$ and obtain, in general, quite different values of y at low x than for the previous procedure.

One should carefully consider the magnitude of errors in y as a function of x before selecting the appropriate procedure. One should apply appropriate weighting to off-set any bias of the least-square procedure as well as to attempt to correct for systematic errors. (17)

The first example will be the fitting of a set of data (x,y) to a three constant equation, $y = a + bf_1(x) + cf_2(x)$. To minimize $[y-a-bf_1(x)-cf_2(x)]^2$, one takes partial derivatives with respect to a, b and c and sets them equal to zero to obtain

ia + $\Sigma f_1(x)b$ + $\Sigma f_2(x)c$ = Σy $\Sigma f_1(x)a$ + $\Sigma [f_1(x)]^2b$ + $\Sigma f_1(x)f_2(x)c$ = $\Sigma y f_1(x)$ $\Sigma f_2(x)a$ + $\Sigma f_1(x)f_2(x)b$ + $\Sigma [f_2(x)]^2c$ = $\Sigma y f_2(x)$

where the sums are over the i values of x. Directions:

(1) Insert tape LSf(x) (2) GTO f d W/PRGM, key in $f_1(x)$, SST SST, Key in $f_2(x)$, RUN (3) A (4) $x_i B$ repeat 4 & 5 for all i (5) $y_i R/S$ repeat 4 & 5 for all i (6) R/S (7) SST SST (8) x C y

NOTE: Return to step 3 for each new set of data. To add more data after step 7, RCL3 STOX5 and start at step 4 again.

Test	<u>;</u> :										
f](3	() = lnx	$f_{2}(x)$	$= x^{-1}$								
(2)			SST SST	h l/x RU	JN; (3)	A					
(41)) 1	в 0.000	; (5') 2	0.0 E	R/S 20.	000					
(4")) 10	B 2.303	; (5") 1	5.605 F	R/S 1.	561					
(4")			; (5") 1								
(6)			(7) SS				3) 10 (C 15.6	05		
<u>.</u>	h. LSf(x)										
(3) (4) (5) (6)	flbLA flbLB flbL1 RCLΣ+	fCLREG † STO+5 STO1	fP≵S ↑ X hR↓	fCLREG ↑ STO+6 STO2	hRTN fe hR↓ fP ∠ S	hR↑ hLSTx RCL9	fd X PCI 8	Σ+ STO+7 RCL7	hLSTx hRTN RCL5	hRTN fP ≵ S	5 1 5 2 4
(0)	STO8 STO÷5	hR↓ ÷ ST08	STO9 RCL8 hx≵y	hR↓ STO4	STO4 RCL1 RCL2	hR↓ RCL2 gx ²		RCL1 ÷	gx ² X RCL9	RCL3 RCL4	
	STO9 RCL1	RCL6 X RCL0	RCL7 hR^{\uparrow} RCL9	RCL5 - X	RCL2 STOO RCL8	X RCL4 hRCI	hx≵y RCL9 X	-	hSTI RCL8 hx≠y	RCL5 gx ² ÷	
	STOB X STOA	RCL4 RCL1 hRTN	X RCLB RCLB	RCLO X RCLC	hx≵y +	RCL3	RCL8 ÷	÷ RCL5	STOC hx≵y	RCL2	11
(7)	fLBLC + gLBLfd gLBLfe	↑ RCLA	↑ + f _l (x) f ₂ (x)	fe hRTN] 1	RCLC hRTN hRTN	X 136 m	hx≵y		RCLB + $f_1(x)$.	X +fo(x)	13:
	-	6.]						kana Ana ang ang ang ang ang ang ang ang ang a	
R	0 1 t Σf ₁ (x) Σf ₂	2 3 (x) i		$\sum_{i=1}^{5} (\mathbf{x}) \Sigma \mathbf{y}$ $\frac{1}{i} \Sigma$		(x) Σy:	7 f ₂ (x) 1	$\sum_{n=1}^{8} \left[\mathbf{f}_{1}(\mathbf{x}) \right]^{2}$		
				đ	i	У		1. 164 AUX - 18 AUX -	S		r
RS	μ Σf _l	(x) Σ[:	$f_{1}^{5}(x)]^{2}$	6 Σf ₂ (x)	$\Sigma(f_2(x$)] ² Σf _l	} (x)f ₂	9 (x) i	A B a b	C I c u	
q =	$\frac{1}{i}(\Sigma f_1(x$)) ² - Σ	$(f_1(x))^2$, r=	$\frac{1}{i}(\Sigma f_2($	x)) ² -	$\Sigma(f_2)$	x)) ² ,			
t =	$\frac{\Sigma y \Sigma f_{l}(x)}{i}$	<u>)</u> - Σyf	1(x),	u =	ΣyΣf ₂ (x)/i -	Σyf ₂ (:	x)			

Linear Regression

For a simple linear regression to fit a set of (x,y) to f(y)=bf(x)+a, the following program is adapted from program Stat 1-22 of the HP-65 User's Library (18). f(x) may be 1/x and f(y) may be ln y or more complicated functions. In the equations for the coefficient of determination and for the standard errors, f(x) will be taken as x and f(y) will be taken as y. The coefficient of determination is given by

$$r^{2} = \frac{[\Sigma xy - (\Sigma x \Sigma y)/i]^{2}}{[\Sigma x^{2} - (\Sigma x)^{2}/i][\Sigma y^{2} - (\Sigma y)^{2}/i]}$$

The standard error estimate of y on x is given by

$$s_{y \cdot x} = \left[\frac{\sum (y - \hat{y})^2}{1 - 2}\right]^{\frac{1}{2}} = \left[\frac{\sum y^2 - a\sum y - b\sum xy}{1 - 2}\right]^{\frac{1}{2}}$$

The standard error of a is given by

$$s_a = s_{y \cdot x} \left[\frac{\Sigma x^2}{i\Sigma x^2 - (\Sigma x)^2} \right]^{\frac{1}{2}}$$

The standard error of b is given by

$$s_{b} = s_{y,x} [\Sigma x^{2} - (\Sigma x)^{2}/i]^{-\frac{1}{2}}.$$

Directions:

(1) Insert tape LR

(2) Key GTO fb W/PRGM, key f(x) SST SST, key f(y) SST SST, key $f^{-1}(y)$ RUN. If f(x) = x or f(y) = y, nothing need be keyed in for f(x) or for f(y) and $f^{-1}(y)$.

(3) fa

(4)	x _l † y _l E	1
	x ₂ ↑ y ₂ E	2
	• • •	•
	x _n ↑ y _n E	n
(5)	А	a
	SST	Ъ
(6)	В	r^2
(7)	xC	ŷ
(8)	D	s _{y•x}
	R/S	sa
	SST	sb

NOTE 1: For data at regular intervals, after fLBLE add hLSTx $f^{-1}(x)I + hx \neq y$ gLBLe and after fLBLC add hR $\uparrow I_1 + \uparrow \uparrow \uparrow$ gLBLc where I is the interval between x values and $f^{-1}(x)$ converts f(x) to x. Step 4 is changed to $x_1 \uparrow y_1$ fe, $y_2 \to y_3 \to \dots \to y_n$ E and step(5) is changed to $x_1 \uparrow \uparrow \uparrow fc \rightarrow y_1$, $C \rightarrow y_2$, $C \rightarrow y_3 \dots \to C \rightarrow y_n$.

Prgm. L	<u>IR</u>									
fLBLE	fd	hx≵y	fb	Σ+	hRTN					6
gLBLa	fCLREG	· ·	hRTN							10
fLBLA	ſ₽ ≵ S	RCL8	RCL4	rcl6	Х	RCL9	÷		STOC	
RCL5		gx^2	RCL9	÷		STOD		STOB	fP≵S	
fx	RCLB	Х	-	STOA	hRTN	RCLB	2			3 7
fLBLB	RCLB	RCLC	Х	fP≑S	RCL7	rcl6	gx^2	RCL9	÷	
	•	f₽≵S	hRTN				-			51
fLBLC	fb	RCLB	Х	RCLA	+	fGSBO				59
fLBLD	fP≵S	RCL7	RCL6	RCLA	X		rcl8	RCLB	Х	
	RCL9	2	-		f√x	hRTN				7.6
RCL5	RCL9	fP≵S		RCLD	÷	fvx	Х	ł	个	
RCLD	fvx	÷.	hx≵y	hRTN	hx≵y					9.2
	(x)]hRT									
gLBLd[f	(\mathbf{y})]hRT	IN IN								
I.TRTO[1	,−1 (y)]h	RTN								
R	4		5	6	5	7		8	9	
		7	$2(f(x))^2$	Υ.	2/~~)	$\nabla(e)_{rr}$	112	$\Sigma f(x)f(y)$	i	
	$\Sigma f(x)$	2	.(I(X))	Δ3	с(у)	2(1(y))	ZI (X)I (Y)	1 I	
R	A	nanya karangkang agan, Matantaka k	В			C			D	
11	А	L			- ()				-	(721)21.
	a	L	Ъ		$\Sigma f(\mathbf{x})$)f(y)-($\Sigma f(x) \Sigma$	f(y)/i Σ(f	·(x)) [—] -	$(\Sigma f/(x))^2/i$

Test:

- (2) GTO fb W/PRCM hl/x SST SST fLN SST SST ge^x RUN
- (3) fa (4) 1300 \pm 1.47×10⁻² E 1, 1400 \pm 2.63×10⁻² E2, 1500 \pm 4.5×10⁻² E3, 1600 \pm 6.96×10⁻²E
- (5) A, a=4.112 SST, b = -10835 (6) B 0.9998
- (7) 1300 C 1.466×10⁻², 1600 C 6.99×10^{-2}
- (8) D, $s_{lny+1/x} = 0.0124$; R/S $s_a = 0.08$; SST $s_b = 745$.

Least-Square Fit of $f(y) = af_1(x) + bf_2(x) + cf_3(x)$

Prgm. $LSf_1f_2f_3$ is an extension of Prgm. LSf(x).

a =
$$(A\Sigma y f_1 - q\Sigma f_1 f_2 + r\Sigma f_1 f_3)/D$$

b = $(q\Sigma f_1^2 - B\Sigma y f_1 + s\Sigma f_1 f_3)/D$
c = $(-r\Sigma f_1^2 - s\Sigma f_1 f_2 + C\Sigma y f_1)/D$
where A = $\Sigma f_2^2 \Sigma f_3^2 - (\Sigma f_2 f_3)^2$
B = $\Sigma f_1 f_2 \Sigma f_3^2 - \Sigma f_1 f_3 \Sigma f_2 f_3$
C = $\Sigma f_1 f_2 \Sigma f_2 f_3 - \Sigma f_1 f_3 \Sigma f_2^2$
D = $A\Sigma f_1^2 - B\Sigma f_1 f_2 + C\Sigma f_1 f_3$
q = $\Sigma y f_1 \Sigma f_3^2 - \Sigma y f_3 \Sigma f_2 f_3$
r = $\Sigma y f_2 \Sigma f_2 f_3 - \Sigma y f_3 \Sigma f_2 f_3$
s = $\Sigma y f_3 \Sigma f_1 f_2 - \Sigma y f_2 \Sigma f_1 f_3$

If f_1 , f_2 , and f_3 are not functions of x directly but are functions f(x') of x' = f'(x), the program provides for conversion of x to x'. For example, if a solidus or liquidus curve is to be expressed as a function of T_{MP} -T rather than T, f'(x) would be keyed in as T_{MP} -CHS. After insertion of a set of data, additional data can be added after steps (5)-(7); follow with steps (5)-(6) for new values of a, b, and c.

Directions:

- (1) Insert tape $LSf_1f_2f_3(x)$
- (2) Key GTO fa W/PRGM, key in f'(x) SST SST SST, key in f₁(x) SST SST, key in f₂(x) SST SST, key in f₃(x) SST SST, key in f(y) SST SST, key in f⁻¹(y) RUN. If any f(x) = x or f(y) = y, key in nothing for that function.
- (3) A (before a set of data to clear registers).Display(4) $x_i \uparrow y_i$ B repeat for all i. $f_1(x)f_3(x)$ (5) Ca(6) SST SSTb,c(7) x D \hat{y} (8) (x_1-100) STOE $E \rightarrow y_1$, $E \rightarrow y_2$, etc for
- (8) (x_1-100) STOE $E \rightarrow y_1$, $E \rightarrow y_2$, etc for x at 100 intervals.

$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$									
	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	LSf _l f ₂ f ₃ (x)								
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	5 4 2								
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$									
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	9								
$\Sigma f_1^2 \qquad \Sigma f_2^2 \qquad \Sigma f_2^2 \qquad \Sigma f_1 f_1^2 \qquad \Sigma f_2^2 \qquad \Sigma f_1 f_2^2 \qquad \Sigma f_2^2 \qquad \Sigma f_1 f_2^2 \qquad \Sigma f_1 f_2^2 \qquad \Sigma $	$\Sigma f_1^2 \qquad \Sigma f_2^2 \qquad \Sigma f_1 f_2$ $A \qquad B \qquad C \qquad D \qquad E \qquad I$ $f(y) \qquad f_1(x) \qquad f_2(x) \qquad f_3(x)$ $A \qquad B \qquad C \qquad D$ $a \qquad b \qquad c \qquad r \qquad q$ $T:$ $GTO \ f \ a \ W/PRGM \ 3 \ EEX \ 3 \ - \ CHS \ SST \ $	i f _l f2								
f(y) $f_1(x)$ $f_2(x)$ $f_3(x)$ A B C D a b c r q TEST: (2) GTO f a W/PRGM 3 EEX 3 - CHS SST SST SST SST gx^2 SST SST 3 h SST SST h 1/x SST SST h 1/x RUN (3) A (4) 1800 $+ 2.2894 \times 10^{-4}$ B 2.07×10 ¹² ; 1900 $+ 2.7465 \times 10^{-4}$ B 1.4	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	9 i								
f(y) $f_1(x)$ $f_2(x)$ $f_3(x)$ A B C D a b c r q TEST: (2) GTO f a W/PRGM 3 EEX 3 - CHS SST SST SST SST gx ² SST SST 3 h SST SST h 1/x SST SST h 1/x RUN (3) A (4) 1800 $+ 2.2894 \times 10^{-4}$ B 2.07×10 ¹² ; 1900 $+ 2.7465 \times 10^{-4}$ B 1.4	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1								
a b c r q TEST: (2) GTO faW/PRGM 3 EEX 3 - CHS SST SST SST SST gx ² SST SST 3 h SST SST h 1/x SST SST h 1/x RUN (3) A (4) 1800 + 2.2894×10 ⁻⁴ B 2.07×10 ¹² ; 1900 + 2.7465×10 ⁻⁴ B 1.4	a b c r q T: GTO f a W/PRGM 3 EEX 3 - CHS SST SST SST SST SST gx ² SST SST 3 h y ^x SST SST h 1/x SST SST h 1/x RUN									
 TEST: (2) GTO f a W/PRGM 3 EEX 3 - CHS SST SST SST SST SST gx² SST SST 3 h SST SST h 1/x SST SST h 1/x RUN (3) A (4) 1800 ↑ 2.2894×10⁻⁴ B 2.07×10¹²; 1900 ↑ 2.7465×10⁻⁴ B 1.4 	T: GTO f a W/PRGM 3 EEX 3 - CHS SST SST SST SST SST gx ² SST SST 3 h y ^x SST SST h 1/x SST SST h 1/x RUN									
 (2) GTO f a W/PRGM 3 EEX 3 - CHS SST SST SST SST SST gx² SST SST 3 h SST SST h 1/x SST SST h 1/x RUN (3) A (4) 1800 ↑ 2.2894×10⁻⁴ B 2.07×10¹²; 1900 ↑ 2.7465×10⁻⁴ B 1.4 	GTO f a W/PRGM 3 EEX 3 - CHS SST SST SST SST SST gx^2 SST SST 3 h y^x SST SST h 1/x SST SST h 1/x RUN									
SST SST h 1/x SST SST h 1/x RUN (3) A (4) 1800 \uparrow 2.2894×10 ⁻⁴ B 2.07×10 ¹² ; 1900 \uparrow 2.7465×10 ⁻⁴ B 1.4	SST SST h 1/x SST SST h 1/x RUN									
(3) A (4) $1800 + 2.2894 \times 10^{-4}$ B 2.07×10^{12} ; $1900 + 2.7465 \times 10^{-4}$ B 1.4										
$2 \pm 23 + 3.333 \times 10^{-4} \text{ B} 4.1 \times 10^{-11}$ 5.1234×10 ⁻⁴ B 4.1×10 ¹¹										
	c, 0.992 (6) SST 1.016×10 ⁻³ SST 9.92×10 ⁻⁷ ; (7) 2200 D 5.1235×10 ⁻¹	¹ ;								
3.333×10^{-4} E 4.100×10 ⁻⁴ E 5.1235×10 ⁻⁴ .										

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The common practice of high precision fitting of experimental data with a power series with a large number of terms is not followed here. The experimental error inherent in most data do not warrant high precision fits and the use of a large number of terms that work against one another makes it impossible to assign physical significance to any When the data have been made available at even intervals, the term. use of orthogonal functions such as the Chebyshev polynomials of Chapter 1 to fit $f(y) = \sum_{n=1}^{\infty} a_n [f(x)]^n$ offers many advantages. The orthogonality makes the solution of matrices simple and a change in the number of terms does not change the coefficients of the earlier terms. However, there are often theoretical reasons to expect a better fit with a mixture of functions such as lnx with powers of x as well as inverse powers of x. Then $f(y) = a + bf_1(x) + cf_2(x)$ or $f(y) = af_1(x) + bf_2(x) + cf_3(x)$ can be fit using the previous programs. The next section describes ways of using additional functions without increasing the number of independent parameters.

Fit of Data to Equations with More Than Three Constants

The solution of simultaneous equations with more than three variables requires 4×4 and higher matrices. If the data are available at even spaced x values, a much simpler procedure involves the use of orthogonal polynomials such as the Chebyshev polynomials. Such a simpler procedure is to use interpolation programs to provide data for evenly spaced x values and to fit the data with a Chebyshev polynomial as illustrated in Chapter I.

Sometimes, it is desired to fit data to a four constant equation, but it may not be advisable to fit an equation with four simultaneously adjustable constants, as errors in the data can yield quite unreasonable constants when the four terms are allowed to work against one another in an attempt to fit the scattered data. A procedure for fitting the four constants with only three being independently adjusted will be illustrated for heat capacity data and for values of enthalpy increments available from drop-calorimeter measurements. The equation $C_p = a + bT + cT^{-2}$ is often adequate for many substances for a considerable temperature range above room temperature and can be fit as described for a three constant fit. The negative cT^{-2} term represents the drop in heat capacity at lower temperatures due to quantized vibrational levels. The 'a' term corresponds to the classical Dulong and Pettit value. The 'bT' arises from electronic, thermal expansion and anharmonicity contributions. These contributions also provide higher order terms which together with a contribution from vacancy formation necessitate a dT^2 term at higher temperatures. The enthalpy equation corresponding to $C_p = a + bT + cT^{-2} + dT^2$ is $H_T - H_{TStd} = a(T - T_{Std}) + \frac{1}{2}b(T^2 - T_{Std}^2) - c(T^{-1} - T_{Std}^{-1}) + (d/3)(T^3 - T_{Std}^3)$. As the enthalpy data often have approximately the same percent error over the temperature range of study, the expression to be fitted by least-square procedures is often

 $(H_{T}-H_{T_{Std}})/(T-T_{Std}) \approx a + \frac{1}{2}b(T+T_{Std}) + c/TT_{Std} + (d/3)(T^{2}+TT_{Std}+T_{Std}^{2}).$

If both the 'c' and 'd' terms are important, there will be a minimum value of dC_p/dT at $(T^*)^{l_1} = -3c/d$ with the contribution of the dT^2 term becoming very small at lower temperatures and with the contribution of the cT^{-2} term becoming very small at higher temperatures. A first approximation to T* can be obtained from inspection of the data. A revised value can be selected upon examination of the deviations between observed and calculated C_p values near T*. A practical way of treating the data is to replace the constant d by $-3c/(T^*)^{l_1}$ to obtain $C_p = a + bT + c[T^{-2}-3(T^*)^{-l_1}T^2]$. One could also replace c to obtain $C_p = a + cT + d[T^2 - \frac{1}{3}(T^*)^{l_1}T^{-2}]$. The equations are equivalent. Program LSf(x) can be used with $f_1(x) = T$ and $f_2(x) = T^2 - 3(T^*)^{-l_1}T^2$ if the constant d is retained or with $f_2(x) = T^2 - \frac{1}{3}(T^*)^{l_1}T^{-2}$ if the constant d is retained. Program LSC^{*}_p is given here with the constant c retained.

Directions:

(1)	Insert tape LS Cp	Display
(2)	Т* А	Display T*
(3)	T B (repeat for all data	Т
(4)	T B $C_{p} R/S$ repeat for all data	$C_p f_2(x)$
	R/S	a
(6)	SST SST SST	b,c,d
(7)	ТС	cp
(8)	Ti ↑ Hi-H298 fa	T_{i}
(9)	T D	H-Hi
(10)	R/S	H-H298

NOTE: If $T_i = 298.15$, key $H_i - H_{298} = 0$ in step 8. Start at step 2 for each new set of data. To add more data after step 6, RCL3 STOX5 and start with step 3 again. <u>Test</u>:

(2)		(3) 50								.6×10	- 5;
	1500 B; 4.7056 R/S - 1.0×10 ⁻⁴ ; 2EEX 3 B; 6.975 R/S - 2.8×10 ⁻⁴										
(5)) R/S 0.999; (6) SST 0.001001, SST -99932; SST 0.9998×10 ⁻⁶										
(7)		1.350;								R/S I	179.9
Prgm	. LS C*				 Addition of the second sec second second sec						
(2)	fLBLA	fCLREG	fP≉S	fCLREG	STOE	hRTN					6
(3)	fLBLB	个	1	个	fe	hR↑	Σ+	hLSTx	hRTN		15
(4)	ST0+5	X	sto+6	hR↓	hLSTx	Х	ST0+7	hRTN	(,	rtana	23
(5)	The ne	X xt 91 st CHS	eps are	unchang	ged from	Prgm.	LSf(x) throu	igh STOA	steps	25-115/
	RCLC	CHS	3	Х	RCLE	4	hy^X		STOD	hx≿	y.
(6)	bRTN	RCLB	RCLC	RCLD			•				128
	gLBLfe	RCLE		4	hy^X	3	hx≵y	+	CHS		
	hx≹y	gx^2	Х	hLSTx	hl/x	÷	hRTN				143
(7)	fLBLC	个	个	1	RCLD	Х	RCLB	+	Х	RCL	
	hR↑	gx^2	+	+	RCLA	+	hRTN				160
(8)	gLBLa	STOO	hR↓	hSTI	hRTN						165
(9)	fLBLD	†	个	个	hRCI	+	Х	hRCI	gx2	+	
	RCLD	Х	3	4	RCLC	hRCI	÷.	hR↑	÷	+	
	hR↑	hRCI	+	RCLB	Х	2	÷	+	RCLA	+	
	hR↑	hRCI	-	Х	hRTN	RCLO	+	hRTN	203 me	emory	steps.
Regi	sters a	re the s	ame as i	for LSf(x) exce	pt for	Hi-H2	98 in C) and for	· A-I	•
	А		В	С		D		E	I		
	a		Ъ	с		ď.		T*	u		
									Ti		

Program LS C_p^* can be readily adapted to the fit of $({\rm H_T-H_T}_1)/({\rm T-T}_1)$ values to the constants a, b, c and d.

$$(H_{T}-H_{T_{1}})/(T-T_{1}) = a' + b'T + c'f(T)$$
 where

$$a' = a + \frac{1}{2}bT_{1} - cT_{1}^{2}(T^{*})^{-l_{4}} = a + \frac{1}{2}bT_{1} + \frac{1}{3}dT_{1}^{2}$$

$$b' = \frac{1}{2}b - cT_{1}(T^{*})^{-l_{4}} = \frac{1}{2}b + \frac{1}{3}dT_{1}$$

$$c' = -c(T^{*})^{-l_{4}} = \frac{1}{3}d$$

$$f(T) = T^{2} - (T^{*})^{l_{4}}/TT_{1}$$

Directions: Insert tape LS H* (1)Display T* (2)T* ↑ T; ↑ H;-H298 A Т (3)Τ В repeat for all data $f_2(T)(H-H_i)/(T-T_i)$ (4) H-H-i R/S R/S (5)đ (6) SST SST SST c,b,a (7)т С Cp (8)T D H-H; (9)R/S H-H298 Test: (2)740 ↑ 298.15 ↑ 0 A 740 500 B; (4) 180 R/S - 1570828; EEX 3 B; 1246.5 R/S - 10220; (3)1500 B; 3130 R/S 4113513; 2EEX3 B; 6030 R/S 12391015; $R/S 9.9997 \times 10^{-7}$ (6) SST - 99953, SST 0.001000, SST 1.00048; (5)(7)500 Cl.35; 500 D 180; R/S 180 Prgm. LS H* hR↓ fPzS hRTN ST00 **fCLREG** hSTI hR↓ STOE (2)fLBLA fCLREG fLBLB 个 ϯ ↑ fe hR↑ Σ+ hLSTx hRTN (3) 19 ÷ (4) ST00 CLx hRCI RCLO hx≵y hR↑ hx≵y _ ST0+6 hR↓ hLSTx Х STO+7 hRTNST0+5 Х 36 The next 91 steps to STOA are unchanged from Prgm.LSf(x) except for (5)changing hSTOI and hRCI to STO6 and RCL6 at steps 73 and 90. RCLC 3 Х STOD RCLB RCLC hRCI Х RCLB ST00 ST01 2 Х STOB hR↓ hRCI Χ RCLA hx≵y 4 hy^X Х CHS STOC STOA RCLE RCLC CHS hR↓ RCLB RCLD hRTN hR↓ hR↓ hx_{zy} 164 hLSTx f√x RCLD (8)fLBLD RCL1 hx≵y Х hLSTx fe hx≁y Х 3 ÷ hx≠y hR↓ ÷ RCLO + hR↑ hRCI Х hRTN fP≵S RCLO ϯ hRTN fP≵S 192 ÷ CHS 4 hyX ÷ gLBLfe RCLE ϯ hRCI hx₹y gx^2 hRTN hLSTx +206 gx² hx≵y Χ (7)fLBLC ↑ ϯ hl/xRCLC Χ RCLB RCLA +hRTN 220 memory steps. + All registers are the same as for LSf(x) except that u is stored in R6, H_i-H_{298} in SO, a' is added to RO, b' is added to R1 and A to I are as follows: С Ι D В Ε А ሞ* Ъ' c' đ T_{i} a' Ъ С a

Program LSf(x) and the related programs fit data to a three constant equation. To least-square fit the two constant $y = bf_1(x) + cf_2(x)$ equation, it is necessary to minimize $\Sigma[y - bf_1(x) - cf_2(x)]^2$ to obtain the two equations: $b\Sigma f_1^2 + c\Sigma f_1 f_2 = \Sigma y f_1$ and $b\Sigma f_1 f_2 + c\Sigma f_2^2 = \Sigma y f_2$. Simultaneous solution yields $b = (\Sigma y f_1 \Sigma f_2^2 - \Sigma y f_2 \Sigma f_1 f_2)/D$ and $c = (\Sigma y f_2 \Sigma f_1^2 - \Sigma y f_1 \Sigma f_1 f_2)/D$ where $D = \Sigma f_1^2 \Sigma f_2^2 - (\Sigma f_1 f_2)^2$. With $f_1(x) = x^0$, the above equations would also fit to a equation with a constant term plus a function of x. However, program SD-03A of the HP-67 Standard Pac(18)covers several of the commonly used simple functions such as y = a + bx, $y = ae^{bx}(a>0)$, y = a + blnx, and $y = ax^b(a>0)$.

A program based on the above equations for a two constant fit is given here to deal with properties of binary solutions as a function of molefraction, x, which varies from 0 to 1. The properties C_p , H, S and G are expressed⁽¹²⁾ as functions of x in the following manner:

 $\bar{Y}_1 = bx_2^2 + cx_2^3$, $\bar{Y}_2 = (b + \frac{3}{2}c)x_1^2 - cx_1^3$, and $Y = x_1x_2(b + \frac{1}{2}c + \frac{1}{2}cx_2)$. In many solutions, the change from an excess of one component to an excess of the other component changes the character of the interaction and the constants of the function. (19) Thus, it is often expedient to fit each portion of the composition range independently with the change often at the midpoint, but frequently at other compositions depending upon the orbitals being utilized as one metal is added to another. When a single equation is not used over the entire composition range for partial molal quantities, a constant of integration of the Gibbs-Duhem equation must be evaluated at the overlapping composition for \bar{Y}_1 at high x_2 or for \bar{Y}_2 at high x_1 to assure continuity of the functions. The programs given below provide for evaluation of the constant of integration when necessary.

Because of occurrence of substantial systematic errors in many data, complications in the character of the interactions, and inherent bias of the least-square process when applied to the functions of interest, suitable weighting of the data is very important. A variety of weighting procedures are used. Key A makes a least-square fit to $\overline{Y}_1 = bx_2^2 + cx_2^3$, which heavily weights values at high x_2 . Key B uses $\overline{Y}_1/x_2 = bx_2 + cx_2^2$ to reduce the bias of the least-square process toward high x_2 values. Key C uses $\overline{Y}_1/x_2^2 = b + cx_2$ to further reduce the weight at high x_2 . Components are assigned 1 and 2 so that the experimental \overline{Y} is \overline{Y}_1 . When integral quantities are available, key D fits $Y/x_1x_2 = b + \frac{1}{2}c + \frac{1}{2}cx_2$. When a single equation for Y_1 is not used over the entire composition range, the integration of the Gibbs-Duhem equation ⁽¹²⁾ to obtain \overline{Y}_2 at high x_1 requires evaluation of a constant of integration which is obtained from $d = \frac{1}{4}[(b'+c')-(b''+c'')]$ where the single primes apply to the constants for the $x'_2 = 0.5 - 1$ range and the double prime applies to the constants for the $x''_2 = 0 - 0.5$ range. When the change in equations is at a value of x_1 not at the midpoint, $d = x_1^2 \{[b' + c'(\frac{3}{2} - x_1)] - [b'' + c''(\frac{3}{2} - x_1)]\}$ which is stored in RD for evaluation of \overline{Y}_2 and Y at low x''_2 .

To calculate \overline{Y}_1 , \overline{Y}_2 or Y given b and c, store b in RB and c in RC and steps 6 to 8 will yield the desired quanties. If necessary, d should also be stored in RD.

Directions:

(1)	Insert Prgm. LS Y	
(2)	fa	Display
	$\begin{array}{c} \mathbf{T}_{0} \text{fit } \overline{Y}_{1} = bx_{2}^{2} + cx_{2}^{3} \\ \hline x_{2} A \\ \overline{Y}_{1} R/S \end{array} \right\} \text{ repeat for all data}$	
(3a)	$x_2 A$ (repeat for all data	x_2^2 $\overline{Y}_1 x_2^3$
(4a)	\overline{Y}_{\perp} R/S ()	$\overline{\mathbf{Y}}_{1}\mathbf{x}_{2}^{3}$
(5)		b,c
(6)	x ₂ E	Ī
(7)	R/S	Ϋ́2
(8)	R/S	Y
	Tofit $\overline{Y}_1/x_2 = bx_2 + cx_2^2$	
(ЗЪ)	$\begin{array}{c c} \hline Tofit \ \overline{Y}_1/x_2 = bx_2 + cx_2^2 \\ \hline x_2 & B \\ \overline{Y}_1 & R/S \end{array} \begin{array}{c} repeat \ for \ all \ data \end{array}$	x ₂
(4b)	\bar{Y}_{1} R/S)	Ī _l x ₂
	Steps 5-8 as above	
	To fit $\overline{Y}_1/x_2^2 = b + cx_2$	
(3e)	$\left. \begin{array}{c} x_2 & C \\ \overline{Y}_1 & R/S \end{array} \right\}$ repeat for all data	^x 2 Ÿ1/x2
(4e)	Ϋ́ _l R/S)	Ī1/x2
	Steps 5-8 as above	
	To fit $Y/x_1x_2 = b + \frac{1}{2}c + \frac{1}{2}cx_2$	
(3d)	To fit $Y/x_1x_2 = b + \frac{1}{2}c + \frac{1}{2}cx_2$ $x_2 D$ Y R/S repeat for all data	хl
(4d)	Y R/S)	Y/x _l x ₂
	R/S	$b + \frac{1}{2}c, \frac{1}{2}c$
(5d)	R/S	c,b
	Steps 6-8 as above	

- NOTE 2: Once constants b, c, and, if necessary, d are stored, steps 6-8 can be repeated for any values of x in range of fit.
- NOTE 3: For calculations of \overline{Y}_2 at low x_2 , d is needed if the original \overline{Y}_1 data were not fit to a single equation for the entire composition range as explained in the text. For such a situation, follow step 6 with step 6' after evaluation of data for $x_2^{"} = 0$ to 0.5 range. Step (6'): (b'+c') fe will store d = $\frac{1}{4}[(b'+c')-(b"+c")]$ in register D. When the break in the data treatment is not at x = 0.5, use equation for d in the text.

Ϋ́ Prgm. 5 (2)**f**CLREG hRTN gLBLfa **f**CLREG f₽żS gx2 (3a) fLBLA ↑ Χ hLSTx Σ+ hLSTx hRTN 13 (4a)Х ST0+1 hx≵y hLSTx Х ST0+2 hRTN GTOO 21 gx^2 hLSTx ∑+· hLSTx hRTN (3b) fLBLB 27 (4b)ST0+1 Х SII:0+2 hRTN GTOO 32 (3c) fLBLC 1 Σ+ hR↓ hRTN 37 gx² hLSTx (4c) \uparrow hR↓ hx≵y ÷ ST0+2 hR↑ 4 ST0+1 hRTN GTO0 49 CHS hRTN (3d) fLBLD fGSBC \uparrow ↑ 1. 57 ST0+1 ÷ hRTN (4d) ST0+2 ÷ hx_Ży hx-y 64 STO4 hR↓ fP≵S RCL8 RCL7 ſ₽≵S ST05 hR↓ (5)fLBLO RCL5 RCL2 RCL4 RCL5 Х ST03 Х RCL1 Х _ RCL3 gx^2 ÷ RCL4 STOC RCL2 ST00 RCL1 RCL5 Χ **** ÷ (5d)RCL4 STOB X RCLO RCLC hRTN 43443 f-x-•---STOB hLSTx 2 Х STOC hRTN f-x-RCLB 112 $\bar{gx^2}$ (6)STOA Х hLSTx RCLC Х RCLB +**fLBLE** ST09 hRTN 123 gx^2 ST08 RCLC 2 ÷ RCLA 4 (7)1 RCL9 12ed Х RCLD +hRTN 137 (8)RCL9 Х RCL8 χ + hRTN hx≵y 144 4 ÷ STOD hRTN (61) gLBLfe RCLB RCLC + 153 0 1 2 4 5 8 9 R 3 Σf2 Σf_1^2 D Σyfı Eyf2 Σfjf2 ХŢ X2 5 8 А В С D RS 7 Σf_2^2 Σf_1^2 b+1/2 $\frac{1}{2}c$ Σfjf2 đ b+ex2 Ъ С

The following values can be used to test Prgm. $\overline{Y}.$

x2	0.1	0.2	0.3	0.4	0.5	0.6	0.7	0.8	0.9
Ϋ́l	4	24	72	160	300	504	784	1152	1620
Ϋ́2	1134	1024	882	720	550	384	234	112	30
Ϋ́	117	224	315	384	425	432	399	320	189

As these values are derived from $\overline{Y}_1 = bx_2^2 + cx_2^3$ with b = 200 and c = 2000, any of the weighting procedures of Prgm. \overline{Y} will yield b = 200 and c = 2000.

CHAPTER III

Standard Thermodynamic Calculations

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A. Thermodynamic Properties of Gases

The several programs tabulated here are designed to calculate the translational, electronic, vibrational and rotational contributions to the thermodynamic functions of gaseous molecules using the equations given in Chapter 27 of Ref.(12). For gaseous atoms, Prgm. E-A provides the electronic and translational contributions to the thermodynamic functions. It can also be used to calculate only the electronic contribution for either atoms or molecules.

Program D-LT calculates the translational, rotational and vibrational contributions for diatomic or linear polyatomic molecules. Program Cor calculates the corrections to the results of Prgm. D-LT if the effects of anharmonicity, centrifugal stretching and vibration-rotation interaction are to be included. When it is desired to obtain moments of inertia for the rotational calculations from molecular constants, Prgm. ID can be used to calculate B_e values for diatomic or linear molecules, moments of inertia or products of moments of inertia for a variety of molecules including bent triatomic, symmetric top CBA₃, octahedral BA₆, etc. Program NL uses the product of the moments of inertia of non-linear molecules to calculate the translation, rotational and vibrational contributions.

The spectroscopic constants used should be expressed in units of cm^{-1} , the unit conventionally used in spectroscopic publications. All the thermodynamic quantities are calculated in dimensionless form, which is most suitable for direct calculation of equilibrium constants. However, the programs provide for multiplication by R in appropriate units to yield the functions in the units desired.

In identifying quantities, the subscripts e, r, t and v are used to refer to electronic, rotational, translational and vibrational contributions, respectively. The superscript ^o for quantities including translational contributions indicate that the values are for the gaseous standard state, although the superscript is not always shown, as all values are for the gaseous standard state. Diatomic or Linear Polyatomic Gases

Program D-LT uses the rotational constants $B_{\rm e}$ and $\alpha,$ the symmetry number σ , the vibrational constants we and $x_e w_e$, and the molecular weight The calculations using the equations on pp. 420-30 of Ref.(12) proceed Μ. as follows with $y = hc(B_e - \frac{1}{2}\alpha)/kT$ and $u = hc(w_e - 2x_ew_e)/kT$. $C_{r,t}/R = \frac{7}{2} + y^2/45$, $S_{r,t}/R = \frac{7}{2} - \ln y\sigma - y^2/90 + \frac{3}{2}\ln M - 3.665 + \frac{5}{2}\ln T$, $-(G-H_{o})_{r,t}/RT = S_{r,t}/R + y/3 + y^{2}/45 - \frac{7}{2},$ $(H_{298}-H_{0})_{r,t}/R = 298.15(\frac{7}{2} - \frac{1}{3}y_{298} - \frac{1}{45}y_{298}^{2}),$ $-(G-H_{298})_{r,t}/RT = -(G-H_0)_{r,t}/RT + (H_{298}-H_0)_{r,t}/RT,$ $(H-H_0)_{r,t}/RT = \frac{7}{2} - \frac{1}{3y} - \frac{1}{45y}^2$ $(H-H_{298})_{r,t}/R = [(H-H_{o})_{r,t}/RT - (H_{298}-H_{o})_{r,t}/RT]T$ with $-(H_{298}-H_{o})_{r,t}/RT$ obtained by subtracting $-(G-H_{298})_{r,t}/RT$ from $-(G-H_0)_{r,t}/RT$ stored in the stack. $C_v/R = u^2 e^{u}/(e^{u}-1)^2$, $-(G-H_O)_v/RT = -ln(1-e^{-u})$, $(H-H_{O})_{v}/RT = u/(e^{u}-1), S_{v}/RT = (H-H_{O})_{v}/RT - (G-H_{O})_{v}/RT,$ $(H_{298}-H_{O})_{v}/R = (uT)/(e^{u298}-1), -(G-H_{298})_{v}/RT = -(G-H_{O})_{v}/RT + (H_{298}-H_{O})_{v}/R$ $(H-H_{298})_v/R = T[(H-H_0)_v/RT - (H_{298}-H_0)_v/RT].$ The various contributions to S/R, Cp/R, -(G-H₂₉₈)/RT, (H-H₂₉₈)/R and (H₂₉₈-H₀)/R are summed and are finally stored in registers 5 to 9, respectively.

Directions:

(la) Insert Tape D-LT

(lb)	Insert Data Tape D-LT	
	<u>Diatomic Gas</u>	Display
(2)	Μ ↑ Β _e ↑ α ↑ σ Α	1.5 lnM-3.665
(3)	w _e ↑ x _e w _e B	uT
(4)	ТС	C _p /R
(5)	R/S	S°/R
		-(G°-H ₂₉₈)/RT
		(H ₀ -H ₂₉₈)/R
(6)	SST	(H ₂₉₈ -H ₀)/R
(7)	R D	Cop
(8)	R/S	S°
		-(G°-H ₂₉₈)/T
		-(G°-H ₂₉₈)/T H°-H ₂₉₈
cont'	d.	

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cont'd. from p. 1

	<u>Triatomic Linear Gas</u>	Display
(2)	Μ↑ B _e tαtσ A	1.5 lnM-3.665
(3)	we't xewe'B	u'T
(3a)	w <mark>"</mark> ↑x _e w <mark>e</mark> fb	u"T
(3b)	w [™] ↑ x _e w [™] f b	u"T (doubly degenerate bend)
(4T)	ТЕ	C _p /R
(5-8)	same as for diatomic gas.	4

- NOTE 1: After constants are entered, start at step 4 or 4T for each temperature. After completing calculations for one molecule, go back to step 1b before entering constants for new molecule. However, entering f P≠S RCL 2 STO A RCL 3 h STI f P≠S or entering 3.665 STO A 1.4388 h STI will replace step 1b and allow starting at step 2.
- NOTE 2: w'_e '' must be doubly degenerate bending frequency. If α or $x_e w_e$ are not known, key zero.
- NOTE 3: Step 7 can follow either step 4, 5 or 6, but step 5 can only follow step 4, step 6 can only follow step 5, and step 8 can only follow step 7.

Test for Diatomic

- (2) 114.938 ↑ .314 ↑ .003 ↑ 1 A 3.452;
- (3) 700 † 7 B 987.017; (4) EEX 3 C 4.423
- (5) R/S 32.981, 29.975, 3005.472; (6) SST 1080.765
- (7) 1.98719 D 8.789. (8) R/S 65.539, 59.567, 5972.444

Test for Linear Triatomic

- (2) 113.936 **†** .211 **†** 0 **†** 2 A 3.438;
- (3) 700 ↑ 0 B 1007.160;
- (3T')650 † 0 fb 935.220;
- (3T[•])150 ↑ 0 fb 215.820;
- (4T) EEX 3 E 7.342; (5) R/S 38.822, 33.874, 4947.9
- (6) SST 1527.743; (7) 1.98719 D 14.590; (8) R/S 77.146, 67.314, 9832.4.

Prgm	. D-LT										
Data	Tape:	298.15	in E, S	3.665 in	n A an	d RS2,	and 1.	.4388 in	n I and	. RS3.	
Tape	D-LT										
(2)	fLBLA	STOC	hR↓	2	*	-	hRCI	Х	STOB	hR↑	
	fLN	1	۲	5	Х	RCLA		STOA	hRTN		19
(3)	fLBLB	fGSBO	STOD	hRTN		-	-				2 3
(3T)	gLBLfb		hF?2	GTO1	STOO		hRTN				30
	fLBL0	2	X 1. TOITINT		hRCI	Х	hRTN				37
(4丁)	fLBL1 fLBLE	STO1 hSFO	hRTN fGSBC	RCLO	fGSB2	ST00	fGSB3	RCL1	fGSB2	ST01	40
(41)	fGSB3	fGSB3	RCLO	fGSB2	RCL1	STO0	hR↓	STOL	hCFO	GTO4	
	fLBL2	RCLD	hx≵y	STOD	hR↓	hRTN					66
(4)	fLBLC	fGSB5	RCLB	hx≉y	÷	ST02	gx^2	4	5	*	
	ST03	3	•	5	ST06	÷	RCLC	RCL2	X	fLN	
	CHS	RCL6	+	RCL3	2	*	-	RCLA	+	hRCI	
	fLN	2	٠	5	Х	+	∑+	hLSTx		3	
	+	RCL3	+	STO3	+	RCL6		RCLB		÷	
	RCLE X	÷ CHS	↑ STO9	gx ² hRCI	5 ÷	÷ +	+ ST07	RCL6	RCL3	RCLE	
	rcl6	СпБ +	hRCI	X	STO8	.1.	0101	-	TOT O		141
	fLBL3	RCLD	hRCI	÷	STO2	gx ²	hLSTx	ge ^x	ST03	Х	
	hLSTx	1		STO4	gx2	÷	RCL3	RCL4	÷	fLN	
	ST05	RCL2	RCL4	*	STOG	+	∑+	RCL5	RCL2	RCLD	
	RCLE	÷	ge ^x	1	-		ST05	÷	ST0 + 7		
	RCL6	÷	hRCI	Х	ST0+8	RCL5	hRCI	Х	ST0 + 9	hF?0	192
(hRTN				_						1 9 2
(5)	fLBL4	$RCL\Sigma +$		hx≵y	STOG	flBL6	hRTN	hR↓	f-x-	RCL7	
(6)	f-x-	RCL8	hRTN	RCL9	amor	00.0	1 -	1.087	1. To 177 by		206
(7)	fLBL5 fLBLD	fP≵S STOX5	0 CITOY	STO4 5 STOX7	STO6	fP ≵S	hR↓	hSTI GTOG	hRTN		215
(1)	ענידד	PIOYO	P.TOY	D DIONI	PLOYO	RCDD	RCL6	GIUO			223
R (0 1	2	3)†		5	6		7	8	9
			y ² /45		7(G.	-H298	Σ	298 5 ^H 298 ^{-H} 0
u	"T u"T	У	y /45	e -1	-1n(1-e ^{-u})	3.5	2	RT	$L = \frac{1}{R}$	Z R
			v . v ²		(H ₂	98-н _о).	v u				
u	'T u"T		$\frac{y}{3} + \frac{y}{45}$			RT	e ^u -	1			
	tro tilen		u			~ /~	50 (D				
u	"T u"'T	u	eu		Σ	S/R	$\Sigma C_{p}/R$				
R	52 S	53	S4 1	36	R	A	В	С	D	E	un la construction de la
	65 1.438					,665	уT	σ		298.15	
3.0	U) I.430	JU 40	/n Lop	(11)			4 ×	~	u"T		
					(3 1nM) .665)			u ‴ T		Т
					(-3	.665/			u'T		

Calculation of Contributions from Anharmonicity, Centrifugal Stretching and Vibration-Rotation Interaction

Program Cor provides the corrections to be applied to the thermodynamic functions calculated for a diatomic molecule by the harmonic-oscillator-rigidrotator approximation. The initial spectroscopic values are used in units of cm⁻¹. The final results are first displayed in dimensionless form with the quantities divided by R, the gas constant. Insertion of R = 1.98719cal/K mol or 8.3144 J K⁻¹mol⁻¹ will yield the values in either set of units. However, the dimensionless values are the most suitable if the values are to be used for subsequent equilibrium calculations; steps 7 and 8 can be omitted if quantities in dimensions are not needed. The calculations are based on the following equations on pp. 430-7 of Ref.(12).

$$\begin{split} & w_{O} = w_{e} - 2x_{e}w_{e}, & B_{O} = B_{e} - \frac{1}{2}\alpha, & x = w_{e}x_{e}/w_{O} \\ & u = (hc/k)w_{O}/T = 1.43879 w_{O}/T, & \gamma = B_{e}/w_{e}, & \delta = \alpha/B_{O} = (B_{e}/\alpha - 0.5)^{-1} \\ & -\left(\frac{(F-H_{O})}{RT}\right)_{cor} = [\delta+2xu/(e^{u}-1)]/(e^{u}-1) + 8\gamma/u \\ & \left(\frac{H-H_{O}}{RT}\right)_{cor} = \{\sigma ue^{u} + [2ue^{u}-(e^{u}-1)](2xu)/(e^{u}-1)\}/(e^{u}-1)^{2} + 8\gamma/u \\ & C_{cor}/R = \{\delta u^{2}e^{u}(e^{u}+1) + [2ue^{u} + u-2(e^{u}-1)](2ue^{u})(2xu)/e^{u}-1)\}/(e^{u}-1)^{3} + 16\gamma/u \end{split}$$
 $S_{cor} = (H-H_O)_{cor}/T - (F-H_O)_{cor}/T$. Display Directions: (1) Insert tape Cor (2) $w_e \uparrow x_e w_e \uparrow B_e \uparrow \alpha A$ $(H_{298}-H_0)_c/R$ $(S_{298})_{c}/R$ (3) SST $(C_{298})_{c}/R$ (4) SST -(G-H298)c/RT (5)ТВ (6) SST will flash the following quantities $(H-H_{298})_c/R$ S_c/R sequentially C_c/R -(G-H₂₉₈)_c/RT -(G-H298)c/T (7) R C (8) SST (H-H₂₉₈)_c S_C -(G-H₂₉₈)_c/

NOTE: Steps 3 and 4 can be skipped if corrections to S298 and C298 are not needed. Once step 2 has been run, start at step 5 for each temperature. Steps 7 and 8 can be omitted if dimensionless quantities are adequate.

Prgm	. Cor		*********		and a de de angen a se aff i 12 de la demo					
(2)	flbla	ST06	hR↓	ST07	hR↓	1	ŕ	2	Х	hR↑
(-)		CHS	1	•)4	3	8	8	hx≵y	X
	STOA	hR↓	hLSTx	÷	STOB	hR↓	8	hx≵y	⊥⊥∧≮y ÷	RCL7
	X	STOD	hLSTx	RCL6	:		5	True y	hl/x	STOC
	2	9	8	•	1	5	fGSBB	RCL7	RCLO	X
(3,4		hRTN	RCL9	RCL8			1,000.0	1001	110100	Λ
(5)	fLBLB	STOO	RCLA	hxty	÷	STOL	↑	ge^{x}	ST02	1
(>)		ST03	-	2	Х	RCLB	X	ST05	RCLC	+
	RCL3	÷	RCLD	RCL1	÷.	+	hSTI	RCL1	RCL2	X
	STO4	2	X2	RCL3		RCL5	X	RCL4	RCLC	X
	+	RCL3	gx ²	+	RCLD	RCL1	÷	+	STO7	RCL4
	2	X	RCL1	+	RCL3	2	Х	-	RCL4	X
	2	Х	RCL5	Х	RCLC	RCL3	2	+	X	RCL4
	Х	RCL1	X	+	RCL3	3	hyx	÷	RCLD	2
	Х	RCL1	÷	+	STO8	RCL7	RCLO	Х	RCLE	
	ST05	hR↓	hRCI	RCL7	+	STO9	hx≵y	RCLE	RCLO	÷
(6)	hRCI	+	STOG	RCL5	hR↓	hRTN	gSTK			
(7)	fLBLC	STOX5	STOX6	STOX8	STOX9	RCL5	RCL9	RCL8	RCL6	hRTN
(8)	gSTK						172 me	emory st	ceps.	
-								Ť		
RО	1 2	3) ₄	5		6		7	8	9
T	l 2 u e ^u	e ^u -1	ueu	2xu		α		^B e	C/R	S/R
Ŧ	u c	e -1	uc	(e ^{u_} H-H ₂	1)		· /1	H-H)	0/1	0/11
					<u>98</u>	(G-H29	8) (1			
				R		RT		RT		
-	А	В		С	D		E		I.	
R				r	0		H298-H	o (G	∃_ _{HO})	
	uT	x		δ	8	γ	R		RT	

Test:

2359.6 + 14.46 + 2.01 + .0187 A 0.181; SST 0.001; SST 0.001 2×10³ B 0.007; SST 23.7, 0.019, 0.030, 0.007 1.98719 C 0.0147; SST 47.1, 0.038, 0.060, 0.0147

Moments of Inertia

The moments of inertia of molecules are needed to calculate the rotational contributions to thermodynamic properties using Prgms. DLT and NL. When there is more than one value, Ig is the moment about the axis of greatest symmetry. I" indicates a doubly degenerate moment. The symmetry number, σ , is tabulated along with the equations for I. If m is the atomic weight of A and r is the internuclear distance in Å, mr² must be divided by 10^{16} N_o = 6.02209×10^{39} to obtain units of g cm². Program ID calculates I×10³⁹ in g cm², B_e = $h/8\pi^2$ cI in cm⁻¹ for linear molecules, and the product of the three principal moments of inertia, $D \times 10^{117}$, for non-linear molecules. For symmetric top molecules and unsymmetric linear molecules, m₁ is atomic weight of A, m₂ of B and m₃ of C.

Diatomic A₂, $\sigma = 2$ I" = $\frac{1}{2}mr^2$ Linear ABA, $\sigma = 2$ I" = $2mr^2$ Planar BA₃, $\sigma = 6$ I_S = $3mr^2$ Tetrahedral BA₄, $\sigma = 12$ I = $\frac{8}{3}mr^2$ Trigonal Bipyramid BA₅, $\sigma = 6$ I_S = $3mr^2$ Uterated and the state of the state of

$$I_{S} = 3m_{1}r_{12}^{2}(1-\cos\theta)$$

$$I'' = m_{1}r_{12}^{2}(1-\cos\theta) + \frac{m_{1}(m_{2}+m_{3})r_{12}^{2}(1+2\cos\theta)}{3m_{1}+m_{2}+m_{3}}$$

$$+ \frac{m_{3}r_{23}}{3m_{1}+m_{2}+m_{3}} \left[(3m_{1}+m_{2})r_{23} + 6m_{1}r_{12} \left(\frac{1-2\cos\theta}{3}\right)^{\frac{1}{2}} \right]$$

If $m_3 = 0$, the equations give I for the symmetric top $B \xrightarrow{A} A$, $\sigma = 3$. If $m_1 = 0$, I" becomes the doubly degenerate moment of inertia of diatomic C-B, $\sigma = 1$.

If $\theta = 0$, I" becomes the moment of inertia of a linear triatomic molecules with atomic masses m_3 , m_2 and m_1 , if m_1 is divided by three at the initiation of the program.

Directions:	Dianlar
(1) Insert Prgm. ID	Display
(2) m † r A	$mx^2/6.02209$
Diatomic A ₂ (3D) R/S	$1" \times 10^{39}$, Be
Linear ABA (3LT)B	1"×10 ³⁹ , B _e
(3P) C	15×10 ³⁹ 1"×10 ³⁹ D×10117
(3T) D	1×10 ³⁹ D×10117
Trigonal Bypyramid BA5 (3TB)E	1 _S ×10 ³⁹ 1¼×10 ³⁹ D×10 ¹¹ 7
Octahedral BA6 (3 OC) f a	1×10 ³⁹ D×10117
Bent Triatomic BA ₂ (3BT)m/m _B ↑ θ f b	I ₁ ×10 ³⁹ I ₂ ×1039 I ₃ ×1039 D×10 ¹¹ 7
(2a) $\frac{\text{Symmetric Top CBA}_3}{r_{12} \uparrow m_1 \uparrow \theta \text{ f c}}$ (3a) $r_{23} \uparrow m_3 \uparrow m_2 \text{ R/S}$	I _S ×10 ³⁹ I"×10 ³⁹ D×10117
(2b) $\frac{\text{Symmetric Top BA}}{r_{12} \uparrow m_1 \uparrow \theta f c}3$ (3b) 0 $\uparrow m_2 R/S$	15×10 ³⁹ 1"×10 ³⁹ D×10117
$\begin{array}{c} \underline{\text{Linear CBA}} \\ (2c) & r_{12} \uparrow m_1 \text{ f d} \\ (3c) & r_{23} \uparrow m_3 \uparrow m_2 \text{ R/S} \end{array}$	0 1"×10 ³⁹ Be
(2d) $r_{23} \uparrow m_3 \uparrow m_2$ f e	I"×10 ³⁹ ^B e

Prgm.	ID		**************************************		1997 P			ndrasan B. ii Midana Ababasa Biyarintan Mikaanan Tikan				
(2)	flBL STO4	А	gx ² ÷	X STOA	6 hRTN	٠	0	2	2	0	9	14
(3D) (3LT)	fLBL		5 2	fGSB0 fGSB0	GTO1. fLBL		2	•	7	9	9	18
(3P)	3 flBL		2 3	hx≵y fGSB0	÷ STOO	STO1 1	hRTN •		fGSBO	GT02		34 43
(3T) (3TB)	fI,BL fLBL		8 3	↑ fGSB0	3 STOO	÷ 3	fGSBO •	↑ 5 ₂	GTO2 fGSB0	GTO2		51 60
(3 OC)	gLBL fLBL		4 RCLA	fGSBO hx ≵y	↑ X	fLBL2 f-x-	STO1 hRTN	gx ²	Χ	ST03	hRTN	76
3BT)	gLBL ÷ 2	fb	fCOS l X	1	+ CHS +	STO8 2 ÷	hx≵y X fGSB0	STO7 fGSBC STO1	h 1/x STOO 2	2 RCL8 RCL8	+ RCL7	
2a-d)	fGSB gLBL		STO2 fCOS RCLA	X STOD X	X hR STOG	STO3 STO5 2	hRTN hx∹≿y		fGSBA	1	RCLD	112
3a-c)	+ fLBL RCL		STO7 STOB	n hx≵y hx≵y ST08	STOO STOC RCLC	2 hRTN + +	X RCLA STO9	÷ X	2 RCL5 RCL7	X 3 X	l X RCL6	137
	+ 6 hRCI		hx ⇒ y X X	hSTI RCLE RCL9	hR↓ X ÷	RCL7 RCL8 RCL4	3 hRCI ÷	÷ X +	f√x + STOl	RCL5 RCLC f-x-	X Z gx ²	
2c) 3c)	RCLO gLBL fGSB		X 3 RCL1	STO3 ÷ fGSB1	hRTN O hRTN	gGSBfc	hRTN				0	191 197 201
2d)	gLBL gGSB	fe	STOB hRCI	CLX RCLC	hR↓ RCLB	STOC fGSB3	CLX RCL1	hR↓ fGSBl	hSTI hRTN	hR↓	hR↓	219
0	1	2	3	ц	5	6		7	8		9	
ı" I _S	B _e I"		6 D	.022	ml	m _l r ₁₂ ² (1-	cosθ)	m mB	l+co	sθ'	3m1+m2+	m3
Il	I ₂	I3	D				1+	2cosθ	3m1+m	2		
A		В		С	*************	D	E		I			
mr .02209	9	m2		m3		cosθ	rj	2	^r 23			
3LT) 1 3T) 1 3 OC) 2a) 2 2b) 2 2c) 2	$ \begin{array}{c} \text{LO} & \uparrow & 2 \\ \text{B} & 13.28 \\ \text{O} & 17.71 \\ \text{f} & 26.9 \\ \text{f} & 26.9 \\ \text{O} & 10 \\ \text{O} & 10 \\ \text{O} & 10 \\ \text{O} & 10 \\ \text{O} & 3 \\ \end{array} $, 0. , 55 57, † 60 † 60 f d	21 57.05 18 75 f c f c 0.00	5.05 6.64 6.64	0.16	(3P) C (3TB) E (3BT) . (3a) 3 (3b) 0	19.93, 1 † 60 † 30 † † 20 R	9.96, 23.25 f b ll 20 R/ /S 8.6	1978.0 , 10 76	9.50 30,3.3 , 2230 .25		52

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Non-Linear Gases

Program ID calculates $D \times 10^{117}$, where D is the product of the three principal moments of inertia, for a variety of non-linear molecules. The following equations from Ref.(12) are used by Prgm. NL to calculate the rotational and translational contributions to the thermodynamic functions of non-linear molecules. The symmetry numbers, σ , are available in Ref.(12) and are also listed in the discussion of Prgm. ID. $C_{r,t}^{o}/R = 4, (H_{298}^{o}-H_{0}^{o})_{r,t}/R = 4(298.15), (H^{o}-H_{298}^{o})_{r,t}/R = 4(T-298.15)$ $S_{r,t}^{o}/R = 4 \ln T + (\frac{3}{2}) \ln M + \frac{1}{2} \ln (D \times 10^{117}) - \ln \sigma - 1.1823$ $4 + \ln 64\pi^{5} + \ln (82.057k^{3})/(h^{6}N_{2}^{5}10^{58.5}) = -1.1823$

 $-(G^{\circ}-H_{298}^{\circ})_{r,t}/RT = S_{r,t}^{\circ}/R - 4(1-298.15/T).$ The vibrational contributions are calculated using the same equations as used for Prgm. D-LT.

The program requires that the singly degenerate vibrational frequencies be inserted first followed by doubly degenerate frequencies and finally by triply degenerate frequencies. The number of frequencies of each type must be entered. i_1 is the number of singly degenerate frequencies, i_2 is the number of doubly degenerate frequencies, and i_3 is the number of triply degenerate frequencies.

Directions:

	Insert tape NL $M \uparrow \sigma \uparrow D \times 10^{117} A$ Values of w in cm ⁻¹ are entered on stack so that singly degenerate values are inserted first (lowest on the stack) followed by doubly and then trip degenerate values. Groups of 1,2, but not over 3 values on the stack are enter and followed by B. A maximum of 10 values	w2 [™] 1 ^D →0 w" ¹ [™] D→0
	can be entered. $i_1 \uparrow i_2 \uparrow i_3 C$	$10^{2}i_{\perp} + i_{2} + 10^{-2}i_{3}$
(5)	T D	-(GºH° ₂₉₈)/RT
(6)	R/S	(H°H ² 98)/R
		S°/R
		C°/R
(6')	SST	(H ₂₉₈ -H°)/R

cont'd.

Directions, cont'd.

- (7) R STO B
- (8) E
- (9) R/S

- R --(G°-H°298)/Т H°-Н298 S° С°р
- NOTE 1: Calculations at each temperature start at step 5 once the constants have been inserted. Step 7 need not be repeated unless new constants are entered in step 2.
- NOTE 2: If dimensionless quantities are not wanted, step 6 can be deleted. SST after either step 6 or step 9 will display value of $(H_{298}^2-H_0^2)/R$.

NL Test:

- (2) 133.936 ↑ 2 ↑ 202.5 A 0.000;
- (3) 700 ↑ 150 B 0.000; (4) 2 ↑ 0 ↑ 0 C 200.00;
- (5) EEX 3 D 35.772; (6) R/S 4588.2; 40.360, 6.835, SST 1466.9;

(7) 1.98719 STO B; (8) E 71.086; (9) R/S 9117.6, 80.204, 13.583.

Prgm. NL											
(2)	fLBLA	fCLREG	fP≵S	0	hR↓	f√x'	-	fLN	CHS	hx≵y	
	fLN	1	•	5	Х	1	•	1	8	2	
	3		+	STOD	CLX	1	0	hSTI	CLx	hRTN	
(3)	fLBLB	1	•	4	3	8	8	Х	STO(i)	fISZ	
	CLx	hR↑	CLx	hR↓	hR↓	fx=0	hRTN	GTOB			48
(4)	fLBLC	EEX	2		÷	hx≵y	EEX	2	Х	+	
	STOA	hRTN									60
(5)	fLBLD	STOC	fLN	4	Х	RCLD	+	ST05	4	ST06	
	RCLC	Х	2	9	8	٠	1	5	STOE	4	
	Х	ST09	58 *	STO8	RCLC	÷		STO7	RCLA	EEX	
	2	÷	fINT	f≠0	fa	RCLA	EEX	2	* *	gFRA	
	f=0	GTO8	EEX	2	Х	fINT	fb	fLBL8	RCLA	gFRA	C
	EEX	2	Х	ſ≠O	fc	GTO7					116
	gLBLfa	1	0	hSTI	+	STO4	fLBLO	fGSB3	gx=y	hRTN	126
	GTOO										127
	gLBLfb	STO+4	fLBL1	fGSB4	fGSB3	gx=y	hRTN	GTO1			135
	gLBLfc	STO+4	fLBL2	fGSB4	fGSB4	fGSB3	gx=y	hRTN	GTO2		144
	fLBL3	fGSB4	fISZ	hRCI	RCL4	hRTN		w.			150
	fLBL4	RCL(i)	RCLC	<u>.</u>	STOO	gx2	hLSTx	ge^{X}	STO1	Х	160
	hLSTx	1		ST02	gx^2	÷	sto+6	RCL1	RCL2	÷	
	fLN	STO3	RCLO	RCL2	* 	ST01	+	STO+5	RCL3	RCLO	
	RCL(i)	RCLE	*	ge^X	1	-	÷	STO3	+	STO+'	
	- hRTN	RCLL	+	RCLC	Х	ST0+8	RCL3	RCLC	Х	STO+	9 201
(6,6') (8,9)	fLBL7 fLBLE	RCL7 RCLB	hRTN STOX5	rcl8 stox6	f-x- STOX7	rcl5 stox8	f - x- GT07	RCL6	hRTN	RCL9	211 218

continued...

R	0 u	l e ^u	2 e ^u -1	3 -ln(l-e ^{-u})	4 Index limit		$\Sigma = \frac{\binom{7}{(G-H_{2}98)}}{RT} \sum_{R}^{H-H_{2}98} \sum_{R}^{\frac{9}{H_{2}98}-H_{0}} R$
		u e ^u -1		(H ₂₉₈ -H ₀) _v RT			
	A)i ₂ •0i	E B E B	ξ Τ	$\begin{pmatrix} \mathbf{S}_{r,t}/\mathbf{R}\\ -\mathbf{h}_{1nT} \end{pmatrix}$	E 298.15	1 Index	RSO-9 are available for up to 10 values of uT.

cont'd. from pg. 38

Calculation of Electronic Level Contributions to Thermodynamic Functions

Program A can be used to calculate the complete thermodynamic functions for gaseous atoms including translational and electronic contributions or it can be used to just calculate the electronic level contributions for either atoms or molecules. Degeneracies and energies of up to seventeen levels can be stored for the calculations. Although the program is intended for calculating the contribution of electronic levels to thermodynamic functions, it can be used for any internal energy levels. For example, if the vibrational or rotational levels are sufficiently perturbed to introduce significant error even after using Prgm. Cor for anharmonic and centrifugal stretching corrections, the partition function can be calculated, seventeen levels at a time, to evaluate the vibrational or rotational contributions.

.1

The spectroscopic unit, cm^{-1} , is used for the energies. The thermodynamic functions are calculated in dimensionless units, which are most appropriate for equilibrium calculations, but insertion of the gas constant R in appropriate units will provide the thermodynamic functions in the same units. The calculations follow the equations on pp. 422-4 of Ref.(12).

After energy levels and degeneracies have been stored in the preliminary steps, calculations at various temperatures commence with step(5). There are several choices in the sequence of steps following step(5). If the interest is primarily in the complete thermodynamic functions for an atom, step (5) would be followed by step(6"a) to yield $-(G^{\circ}-H_{OB}^{\circ})/RT$ and then by step(6"b) if the remainder of the functions are desired in dimensionless If $-(G^{\circ}-H_{298}^{\circ})/T$ is desired, step (6"a) would be followed by step(7) form. and then by step(8) if the other functions are of interest. If values for all functions are wanted both with dimensions and in dimensionless form, the sequence after step(5) would be (6"a), (6"b), (7) and (8). If only the electronic contributions are wanted, step (5) would be followed by step(6')and then by steps (7) and (8) for values expressed in the units of the value of R introduced in step(7). If, after obtaining the electronic contributions of an atom through step(6'), the complete thermodynamic functions should also be wanted, replacement of the directions of step(6"a) by hR↓ hR↓ T STO D M D will allow calculation of the complete functions without repeating step(5). The restrictions are summarized as follows: step(8) only after (7), step(6"b) only after (6"a), step(7) after (6'),(6"a), or (6"b), step (6') only after (5), and step(6"a) only after (5) except after (6') with insertion of the special step given above.

In addition to the program tape, a data tape is used to store the following constants: 3.665 STO A, 298.15 STO B, 1438.8 STO C. These constants can be keyed in manually or by tape. If the constants have not been inserted before step 2, integers will be lacking in all of the displays of step 2.

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Directions: (1) Insert tape A and Data tape A $\frac{\text{Display}}{1438\cdot8\epsilon_1\cdot g_1}\times 10^{-2}$ $(2-1)g_1 \uparrow \varepsilon_1 A$ 1438•8E2•B2×10-2 $(2-2)g_2 \uparrow \epsilon_2 R/S$ 1438·8E_i·g_i×10⁻² $(2-i)g_i \uparrow \varepsilon_i R/S$ i up to 17 (3) B (H₂₉₈-H₀)_e/R (4) R/S (S₂₉₈/R)_e (C₂₉₈/R)_e ~(G-H₂₉₈)_e/RT (5) T C (H-H₂₉₈)_e/R S_e/R C_e/R -(G-H₂₉₈)_e/RT (6') R/S (not to be followed by 6a" without step given in text) (H°-H298)/R S°/R (6") M D (6") R/S Cp/R -(G°-H208)/RT

(7)	RΕ	H-H298
(8)	R/S	N
		-(G-H ₂₉₈)/Т

- NOTE 1: If a number of levels are inserted together at an average value of ε_i , the total g must not exceed 99.
- NOTE 2: If values of ε_i larger than 69 502 cm⁻¹ are to be used, change constant in register C to 143.88 and change 3 at step 95 to 2 to allow values up to 695 000 cm⁻¹. With this change,ε values are truncated beyond .01 cm⁻¹.
- NOTE 3: If values for more than seventeen energy levels are inserted, the display will read Error. With seventeen levels, steps (3) and (5) each require 45 sec calculation time. With fewer levels, the time is cut as the calculation stops when a zero degeneracy is encountered. If calculations for a new atom or molecules is started by inserting new level values, one can start at step 2 if the same number or a greater number of levels is inserted: otherwise, one must reinsert

the data card. If only the electronic contributions are desired, up to twenty levels can be stored if three pairs of fGSBO hRTN are added after step 36, RCLC at step 38 is replaced by 1438.8, RCLB at step 54 is replaced by 298.15, the 19 of steps 112-113 is changed to 22 after at least fifteen steps of the fLBLD subroutine have been deleted to accommodate the insertions. Notes, cont'd.

- NOTE 4: The dimensionless values for functions at 298 K given by steps (3) and (4) can be multiplied by R by the following procedure. After step (4), key R in the desired units followed by E to display $(H_{298}-H_{0})_{e}$. Then R/S will display the electronic contributions to S₂₉₈, C₂₉₈, and $-(G_{298}-H_{0})/T$. If the total $H_{298}^{0}-H_{0}^{0}$ is wanted for a gaseous atom, it can be obtained by adding 1481.20 to the electronic contribution in cal/mol.
- NOTE 5: If $-(G-H_O)/RT$ is desired, it can be obtained after step 5 or any of the following steps by keying h RCI. Keying of g e^x would then give Q, the partition function.

Test with Si gas at 5000 K

(2) 1	个	О А		0.010
3	↑	77.12 R,	/S	110960.030
5	↑	223.16 R,	/S	321082.050
5	↑	6298.86	R/S	9062799.050
1	↑	15394.37	R/S	22149419.01
5	↑	33326	R/S	47949448.05
9	↑	39860	R/S	57350568.09
3	↑	40992	R/S	58979289.03
15	个	45303	R/S	65181956.15
3	↑	47284	R/S	68032219.03
5	ϯ	47352	R/S	68130057.05
15	↑	48161	R/S	69294046.15
9	↑	49128	R/S	70685366.09
3	↑	49400	R/S	71076720.03
21	↑	49966	R/S	71891080.21
5	ϯ	50189	R/S	72211933.05
10	↑	50535	R/S	72709758.10

- (3) B 162.716
- (4) R/S, 2.108, 0.176; 1.98719 E 323.347 R/S 4.190, 0.350, 3.105
- (5) 5 EEX 3 C 2.279
- (6') R/S 844.589, 2.448, 0.277, 2.279
- (7) 1.98719 E 1678.4; (8) R/S 4.864, 0.550, 4.528
- or after (5) (6"a) 28.086 D 12599.214; (6"b) R/S 27.579, 2.777, 25.059
- (7) 1.98719 E 25037.031; (8) 54.804, 5.518, 49.796 cal.

Prgm. A											
(2-1 to 3)	fLBLA	3	hSTI	hR↓	fGSBO	hRTN	fGSBO	hRTN	fGSBO	hRTN	
(2-4 to 8)	fGSBO	hRTN	fGSBO	hRTN	fGSBO	hRTN	fGSBO	hRTN	fGSBO	hRTN	
(2-9 to 13)		hRTN	fGSBO	hRTN	fGSBO	hRTN	fGSBO	hRTN	fGSBO	hRTN	
(2-14 to 16)		hRTN	fGSBO	hRTN	fGSBO	hRTN					36
(2-17)	fLBLO	RCLC	X	fINT	hx≵y	EEX	2	*	+	STO(i)	
(a)	fISZ	hRTN	GTO9	CITAR		dillon	fGSB2	ACCD2	Ъ₽₩	STOE	49
(3) (4)	fLBLB STOD	fGSBl hRTN	0 RCL1	STOE f-x-	RCLB RCL2	STOD hRTN	TCOBS	fGSB3	1117 1	DIOR	
(5)	fLBLC	STOD	fGSB1	fGSB2	GTO3	TTT T T					65 70
	fLBL1	0	STOO	STOL	STO2	3	hSTI	hRTN			78
	fLBL2	RCL(i)	^ ^	gFRAC	EEX	2	X	fx=0	hRTN	hx≵y	
	fINT	EEX	3	÷	RCLD	÷	1	hR↓	CHS	geX	
	Х	ST0+0	hR↑	Х	STO+1	hSTx.	Х	ST0+2	fISZ	ĩ	
	9	hRCI	gx>y	hRTN	GT02				-		113
	fLBL3	RCL2	RCLO	*	RCL1	RCLO	÷	STOL	gx^2		
	ST02	RCL1	RCLD	Х	RCLE		RCLO	fLN	hSTI	RCL1	
((.))	+	STOL	RCLE	RCLD	÷.	hRCI	+	ST00	hRTN		142
(6')	hR↑	STOD	CLX	RCL2	hx≵y	gSTK	hRTN	 T	0		149
(6 <mark>"</mark>)	fLBLD	hR↑ CEO+O	hx≵y cmoc	fLN X	hLSTx	RCLD	X RCLA	fLN	2 1	• 个	
	5 hRCI	STO+2 +	STOC hSTI	hR↓	hx ≵y ↑	+ ↑	RCLC	 +	ST0+1	hR↓	
	RCLB	RCLD		RCLC	X	÷	STO+0	hR↓	RCLD	RCLB	
		RCLC	X	+	STOD	1	4	3	8		
(6")	8	STOC	hR↓	fLBL4	hRTN	RCL1	f-x-	RCL2	f-x-	RCLO	
	hRTN										210
(7,8)	fLBLE	STOX0	STOX1	STOX2	RCLD	Х	STOD	GTO4			218
Data Tape:	3.665	STO A,	298.15	STO B, I	1438.8	STO C	•				
R O		1.	2	Regi	sters 3	to 9	have				
Q		ą'	Q"	1438	•8Ei•gi	×10-2	with i	= 1 to	7		
	`	v	·		* **						
-(G-H ₂₉₈	$\frac{3}{2}$ or	'/Q	C/R								
RT	Q	/ %	0/10								
		S/R									
RS 0 to 9 h	nave 143	38•8e _i •g	i×10 ⁻²	with i :	= 8 to 3	17					
R	A	В		С	D		E		I		
			ר 1	38.8	Т		^H 298- ^H	n 3	to 20		
3.	.665	298.15	14	JU.U		0	<u>90</u> R		everal	times	
				2.5	H-H29	8	τι		$(G-H_{O})$		
				/	R			-	RT RT		
			14	38.8					U.T.		

B. Calculation of Thermodynamic Functions from ${\rm C}_{\rm p}$ or H Values

Heat capacity or enthalpy data may have been fit to $C_p/R = a+bT+cT^{-2}+dT^2$ or $C_p/R = a_0 + a_1T + a_2T^2 + a_3T^3$ using the programs of Chapter II. Programs CHSG and Cheb CHSG use the constants of these equations to calculate C_p/R , $(H-H_{298})/R$, S/R and $-(G-H_{298})/RT$.

The C_p equations generally will have a limited temperature range of application fixed by either the range of existence of the phase or by the range for which the equation gives a sufficiently accurate fit. Thus, the programs provide for joining at the temperature T_i at which a new range starts and (H_i-H₂₉₈)/R and S_i/R, calculated in the lower temperature range, are inserted to allow extension into the higher temperature range. The equations for Prgm. CHSG are based on $C_p/R = a + bT + cT^{-2} + dT^2$.

$$(H-H_{i})/R = (T-T_{i})[a+\frac{1}{2}b(T+T_{i})+c/TT_{i}+\frac{1}{3}d(T^{2}+TT_{i}+T_{i}^{2})] (S-S_{i})/R = aln(T/T_{i}) + (T-T_{i})[b+\frac{1}{2}c(T+T_{i})/T^{2}T_{i}^{2}+\frac{1}{2}d(T+T_{i})] -(G-H_{298})/RT = (S-S_{i})/R + S_{i}/R - [(H-H_{i})/R+(H_{i}-H_{298})/R]/T.$$

The constants a, b, c and d are stored in registers A, B, C and D. If the data have been fit to an equation with fewer constants, store 0 for the unused constants. The program has been illustrated with dimensionless functions; it operates equally well if the constants yield C_p rather than Then S_i and H_i -H₂₉₈ should be inserted in steps 3 and 4 and the C_{D}/R . displays of steps 5 and 6 will be the indicated displays multiplied by R. If it is desired to convert the dimensionless displays of steps 5 and 6 to values in cal or j, R in the appropriate units should be stored in register Ε. Then step 7 will display the values multiplied by R. If the displays of steps 5-6 are already in dimensions, step 8 will divide by R from register E to display dimensionless values.

Directions

(1)

Insert Prgm. CHSG; store constants in registers A-D if not carried over from previous calculation Display

 $^{\mathrm{T}_{\mathbf{i}}}_{\mathrm{S}_{\mathbf{i}}/\mathrm{R}}$ T_i S_i/R (2) А (3) SST (H₁-H₂₉₈)/R (H₁-H₂₉₈)/R (H-H₂₉₈)/R, S/R, -(G-H₂₉₈)/RT C_p, H-H₂₉₈, S, -(G-H₂₉₈)/T (4) $(\bar{H}_{i}-H_{298})/R$ SST; 0 SST if T_{i} = 298 (5) Τ (6) SST SST SST (7)С to multiply by R (8) D to divide by R

							1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 -				
Prgm.	CHSG										
(2-4)	fLBLA	ST05	hRTN	STO4	ST03						
(5)	fLBLB	^	ST02	RCL5	+	ST00	Х	RCL5	gx^2	+	
	RCLD	Х	3	÷	RCLC	RCL5	•	RCL2	÷	+	
	RCLO	RCLB	Х	2	÷	+	RCLA	+	RCL2	RCL5	
		ST01	Х	RCL3	+	ST06	RCLC	2	÷	RCLO	
	Х	RCL2	RCL5	Х	gx^2	÷	RCLB	+	RCLD	2	
	÷	RCLO	Х	+	RCL1	Х	RCL2	RCL5	÷	fLN	
	RCLA	Х	+	RCL4	+	ST07	hx≠y	RCL2	÷	-	
	STO9	RCLD	RCL2	gx ²	Х	RCLC	hLSTx	÷	+	RCLB	
(6)	RCL2	Х	+	RCLA	+	STO8	hRTN	RCL6	RCL7	hR↑	
(7)	fLBLC	RCL9	RCLE	Х	hSTI	RCL8	hLSTx	Х	RCL6	hLSTx	
()	X	RCL7	hLSTx	Х	hRCI	gSTK	RTN		DOT	1 7 000	112
(8)	fLBLD	RCL9	RCLE	÷	hSTI	RCL8	hLSTx	÷	RCL6	hLSTx	129
	•	RCL7	hLSTx	•	hRCI	gSTK	RTN				129
R O	1	2	3	4 5	6	7	8	9			
T+T			-H ₂₉₈		н−н ₂	~~ S		-(G-H	$\left[0,08\right) $		
]1	- <u></u>	<u>R</u>	$\frac{S_{i}}{R}$ T_{i}		$\frac{98}{R}$ $\frac{S}{R}$	$\frac{C_p}{R}$	$\frac{-(G-H)}{F}$	2 <u>90'</u>		
			11		21			Ŧ			
	А	В	С	D	Ε	I	-H298)				
	a	Ъ	е	đ	R	-(G	<u>-п2987</u> т				
m											

<u>Test</u>:

10 STO A, EEX CHS 3 STO B, 1 CHS EEX 5 STOC, EEX CHS6 STOD, 1.98719 STOE (2) 298.15 A; (3) 10 SST; (4) 0 SST (5) EEX 3 B 11.900; (6) SST 7563, SST 22.747, SST 15.183; (6) C 23.648, 15029, 45.202, 30.172 (7) D 5.888, 3806, 11.447, 7.641 From $C_p/R = a_0 + a_1T + a_2T^2 + a_3T^3$, $(H-H_i)/R = a_0(T-T_i) + \frac{1}{2}a_1(T^2-T_i^2) + \frac{1}{3}a_2(T^3-T_i^3) + \frac{1}{4}a_3(T^4-T_i^4)$ $(S-S_i)/R = a_0ln(T/T_i) + a_1(T-T_i) + \frac{1}{2}a_2(T^2-T_i^2) + \frac{1}{3}a_3(T^3-T_i^3)$ $-(G-H_{298})/RT = (S-S_i)/R + S_i/R - [(H-H_i)/R + (H_i-H_{298})/R)]/T$ a_0 to a_3 will be in R6 to R9 if evaluated by Prgm. Cheb-a. Directions:

(1) Insert Prgm. Cheb CHSG	and a_0 to a_3 in R6-8 if not already there
(2) T _i A	Display T ²
(3) _{Si} /R SST	S _i /R
(4') (H _i -H ₂₉₈)/R SST	(H ₁ -H ₂₉₈)/R
(4") R SST	R
(5) T B	C _p /R
(6) SST SST SST	(H-H ₂₉₈)/R, S/R, -(G-H ₂₉₈)/RT
(7) C repeats display of st R has been stored in re	eps 5 & 6 multiplied by R if gister E by step 4".
constants gave C _p in di	eps 5 & 6 divided by R if the a _n mensions of R and therefore steps d quantities multiplied by R.
	SST will flash in turn the four values if a review is wanted.
Prgm. Cheb CHSG	

				i/R	(Hi-H29			Т				
2	А	B Ti	-	С	D	E		I				
	C _p /R (H-	H298)/R	S/R	-(G-H ₂	298)/RT	-(G-H	298)/T					
3	0	1	2	T.	3 ³ _T ³	Ţ	2 ⁴ -T ²	5 T-T _i	6 a ₀	7 a _l	8 a2	9 ^a 3
	*	RCL2	RCLE	÷	RCL4	gSTK	hRTN					135
(8)	fLBLD	RCL3	RCLE	 	STO4	RCL0	RCLE	÷	RC	L1	RCL	E
()	X	RCL2	RCLE	X	RCL4	gSTK	hRTN			and other		118
6) 7)	fLBLC	nk↑ RCL3	gSTK RCLE	х	STO4	RCLO	RCLE	х	RC	ГЛ	RCL	
6)	X RCL2	RCL7 hR↑	+ ~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	hRCI	Х	RCL6	+	ST00	hR	'1'N	RCL	1 101
	ST02	hx _₹ y		STO3	RCL9	hRCI	Х	RCL8		+	hRC	
	+	hRCI	RCLA	÷.	fLN	RCL6	Х	÷		LC	+	
	<u>.</u>	RCL4	RCL8	Х	2	<u>+</u>	+	RCL5	RC	L7	X	
	+	RCLD	+	STOL	hRCI	*	RCL3	RCL9		X	3	
	hRCI	14 14	hy ^X	RCLB	gx ²		RCL9	C X		÷ 4	+ +	
	- RCLA	STO4 RCLB	RCL7 X	X	2 STO3	÷ RCL8	+ X	hRCI 3		3 •	hy ^X +	
5)	fLBLB	hSTI	RCLA		ST05	RCL6	X	hRCI	g	х ²	RCL	
2-4	•	STOA	gx ²	STOB	hRTN	STOC	STOD	STOE		2		8

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Test:

 $C_p = a_0 + a_1T + a_2T^2 + a_3T^3$ with Prgm. CHSG

1.50365 STO 6, 8.19918×10⁻³ STO 7, -4.13509×10⁻⁶ STO 8, 9.21717×10⁻¹⁰ STO 9;

(2) 1400 A; (3) S₁₄₀₀ = 16.792 cal K⁻¹ SST; (4) H₁₄₀₀-H₂₉₈ = 7307 cal SST, 1.98719 SST; (5) 1400 B 7.407, SST 7307, SST 16.792, SST 11.573; 2800 B 12.276, SST 20 320, SST 23.057, SST 15.800; SST reviews the four values, D divides them by R.

The programs of Chapters I and II have provided least-square fits of enthalpy or heat capacity data to four constant equations. Programs CHSG and Cheb CHSG calculate values of the other thermodynamic functions from the four constant heat capacity equations. Using eight values of the C_p of molybdenum solid at 200 K intervals from 1400 to 2800 K,

 $C_p = 1.5037 + 8.1992 \times 10^{-3}T - 4.1351 \times 10^{-6}T^2 + 9.2172 \, 10^{-10}T^3$ was obtained which fit the original values with an average deviation of 0.009 or less than 0.1%. The same values were fit by Prgm. LS C_p^* , which uses only three independent constants, to obtain

 $C_p = 9.5466 - 3.8026 \times 10^{-3} T - 1.360 \times 10^{5} T^{-2} + 1.6987 \times 10^{-6} T^{2}$. This equation fits the original values with an average deviation of 0.06 or about 0.6%. Both equations were used with Prgms. Cheb CHSG and CHSG, resrespectively. The values of $-(G-H_{298})/T$ calculated by Prgm. CHSG using the constants from Prgm. LS C_p^* were only slightly lower than those calculated using the constants from Prgm. Cheb-a with an average deviation of 0.002 or 0.02%, thus illustrating the insensitivity of $(G-H_{298})/T$ values to errors in C_p .

It is sometimes useful to go directly from the analytical equations for C_p to an analytical equation for $-(G-H_{298})/T$, which will have six terms derived from the four terms of the original heat capacity equations. Prgm. CgK described below carries out this operation. This program not only calculates values of $-(G-H_{298})/T$, but if the constants for each of the reactants and products of a reaction are combined to obtain an equation for ΔC_p^{o} , the program calculates $-(\Delta G^{o}-\Delta H_{298}^{o})/T$ for a given T. The program can be used for either C_p^{o} or C_p^{o}/R . The program provides for multiplication or division by R to provide consistency. The ability to calculate $-(\Delta G^{\circ} - H_{298}^{\circ})/RT$ allows calculation of $\Delta H_{298}^{\circ}/R$ from a value of K, the equilibrium constant, or the calculation of K from knowledge of $\Delta H_{298}^{\circ}/R$. Normally the C_p equations will be of the same form and $\Delta C_p/R$ will either be

$$\begin{split} & \Delta a_0 + \Delta a_1 T + \Delta a_2 T^2 + \Delta a_3 T^3 \quad \text{or} \quad \Delta a + \Delta b T + \Delta c T^{-2} + \Delta d T^2. \end{split}$$
 Prgm. CgK will handle either form. In addition, if not all of the reactants or product heat capacities are in the same form, the program will handle $\Delta C_p/R = \Delta a_0 + \Delta a_1 T + \Delta a_2 T^2 + \Delta a_3 T^3 + \Delta c T^{-2}. \quad \text{Constants of terms that are not used should be entered as zero.} \\ & \Delta C_p^o/R = \Delta c T^{-2} + \Delta a_0 + \Delta a_1 T + \Delta a_2 T^2 + \Delta a_3 T^3 \quad \text{where b terms are included in } \\ & \Delta a_1 \text{ and d terms are included in } \Delta a_2. \\ & -(\Delta G^\circ - \Delta H_{298}^\circ)/RT = \Delta a_0 \ln T + \Delta a_2 T^{-2} + \Delta a_1 T^{-1} + \Delta a_0^\dagger + \Delta a_1^\dagger T + \Delta a_2^\dagger T^2 + \Delta a_3^\dagger T^3 \\ & \text{where } \Delta a_3^\dagger = \frac{1}{12}\Delta a_3, \quad \Delta a_2^\prime = -\frac{1}{6}\Delta a_2, \quad \Delta a_1^\prime = -\frac{1}{2}\Delta a_1, \quad \Delta a_{-2} = -\frac{1}{2}\Delta c, \\ & \Delta a_0^\circ = \Delta S_1^\circ/R - \Delta a_0(\ln T_1 + 1) + \frac{1}{2}\Delta c T_1^{-2} - \Delta a_1 T_1 - \frac{1}{2}\Delta a_2 T_1^2 - \frac{1}{3}\Delta a_3 T_1^3, \text{ and} \\ & \Delta a_{-1} = \Delta a_0 T_1 + \frac{1}{2}\Delta a_1 T_1^2 + \frac{1}{3}\Delta a_2 T_1^3 + \frac{1}{4}\Delta a_3 T_1^4 - \Delta c T_1^{-1} - (\Delta H_1^\circ - \Delta H_{298}^\circ)/R. \\ & \text{The equations are written for } \Delta H_{298}^\circ \text{ but can be used for } \Delta H_0^\circ \text{ and } (\Delta G - \Delta H_0)/RT \\ & \text{if } (\Delta H_1^\circ - \Delta H_{298}^\circ)/R \text{ is replaced by } (\Delta H_1^\circ - H_0^\circ)/R. \quad T_1 \text{ is the temperature limit} \\ & \text{for range of validity of the } C_p \text{ equations. The following equations are used for calculation of K, } \Delta H_{298}^\circ/R, \text{ or } \Delta G^\circ/R. \\ \end{aligned}$

 $\Delta H_{298}^{\circ}/R = T[-(\Delta G^{\circ}-\Delta H_{298}^{\circ})/RT - lnK].$ The values of ΔH_{298}° from a set of T and K values are averaged and then used for calculation of $\Delta G^{\circ}/R$ and K by $lnK = -\Delta G^{\circ}/RT = -(\Delta G^{\circ} - H_{298}^{\circ})/RT - (\Delta H_{298}^{\circ}/R)T^{-1}.$

When accurate entropies are not available for all reactants and products, both ΔH° and ΔS° can be determined ⁽²¹⁾ through evaluation of

 $\Delta H_{298}^{\circ}/R - T\Delta S_1^{\circ}/R = T[-(\Delta G^{\circ} - H_{298}^{\circ})/RT - \Delta S_1^{\circ}/R - InK] = T[A - InK]$ where A can be calculated from just the constants of the $\Delta C_p^{\circ}/R$ equations by step 10 of Prgm. CgK if zero is entered for $\Delta S_1^{\circ}/R$ in step 3. If steps 14 and 15 are carried out for a set of T and K values with $\Delta S_1^{\circ}/R = 0$, the display of step 15 is $\Delta H_{298}^{\circ}/R - T\Delta S_1^{\circ}/R$ which is stored in the secondary register. These values can then be least-square fit to a linear equation after all data have been inserted in steps 14 and 15 by carrying out step 21. Step 21 displays first $\Delta H_{298}^{\circ}/R$ followed by $\Delta S_1^{\circ}/R$ and incorporates these constants; so that steps 10 and 18 to 20 can then be carried out. Directions: (1) Insert Prgm. CgK (2) Store Δa_0 in RO, Δa_1 in R1, Δa_2 in R2, Δa_3 in R3, Δc in RC, and R = 1.98719 or 8.3144 in RE. Display (3) $T_i \uparrow \Delta S_1^o/R \uparrow (\Delta H_1^o - \Delta H_{298}^o)/R$ A Δa_0 (4) SST Δa_2 (5) SST ∆a_1 (6) SST ∆a¦ (7) SST ∆a' (8) SST ∆az (9) SST ∆az -(\[\G^\-\] \H_2\(\) \RT (10) T B (ll) f a to multiply by R (12) f e to divide by R (13) f c 0.00 before starting set of data in steps 14 and 15 (14) T_i ↑ K_i C $\Delta H_{298}^{\circ}/R$ repeat for all i $\frac{1}{4}\Sigma\Delta H_{298}^{o}/R$ (16) R/S Standard Deviation (17) R/S(18) T D Κ (19) SST ln K $\Delta G^{\circ}/R$ (20) T E $\Delta H_{298}^{o}/R$, $\Delta S_{i}^{o}/R$ (21) f d after step 15 if $\Delta S_i^o/R = 0$ in step 3 Enter 0 for $(\Delta H_1^{\circ} - \Delta H_{298}^{\circ})/R$ if $T_1 = 298.15$ K and enter zero for Δc , NOTE 1: Δa_3 , or Δa_2 if the $\overline{C_p}$ equations do not use these terms. NOTE 2: If it is desired to repeat the calculations using a changed value of T_i , it is necessary to start back at step 2. Steps 4 to 9 can be bypassed if there is no need to check the constants. With $\Delta S_1^{\circ}/R$ known and stored in RA and step 3 completed, steps 10-12, 18

and 20 can be carried out in any order. If S_1°/R is known and has been inserted in step 3, steps 16 and 17 follow the insertion of all of the sets of K and T values. If S_1°/R is not known, it is entered as zero in step 3 and steps 10 and 16-20 cannot be used until $\Delta S_1^\circ/R$ has been determined by carrying out step 21 after the insertion of all of the sets of K and T values. After step 21, steps 10, 18 and 20 can be used in any order.

Prgm.	Cak										
										,	
(3)	fLBLA	STOB	hR↓	STOA	hR↓	1	<u>↑</u>	STOD	RCL3	4	
	<u>.</u>	Х	RCL2	3	÷	+	Х	RCL1	2	÷	
	+	X	RCLO	+	X	RCLC	hR↑	:		RCLB	
		STO4	hR↓	RCL3	3	+	X	RCL2	2	÷	
	+	Х	RCL1	+	X	RCLC	hR↑	gx ²	*	2	
	• •		hR↑	fLN	1	+	RCLO	X	+	RCLA	
()		CHS	ST05	2	STO÷1	6	STO:2	1	2	STO÷3	
(4-7)	RCLC	2	÷	STOC	RCLO	hRTN	RCLC	RCL4	RCL5	RCL1	
(8-9)	RCL2	RCL3		,	-						82
(10)	fLBLB	1	†	1	RCL3	Х	RCL2	+	X	RCL1	
	+	Х	RCL4	hR↑	÷	+	RCLC	hR↑	gx^2	<u>*</u>	
(+	RCL5	+	hR↑	fLN	RCLO	Х	+	hRTN		111
(11)	gLBLa	RCLE	Х	hRTN							115
(12)	gLBLe	RCLE	*	hRTN							119
(13)	gLBLc	fP≵S	0	STO4	ST05	STOG	ST07	ST08	ST09	fP≵S	
(14)	hRTN	fLBLC	fLN	ST06	hR↓						134
(16)	В	RCL6	***	X	hR↑	∑+	hR↓	hRTN	fx	hR↓	149
(17)	ST07	hRTN	gs	hR↓	hRTN		v				
(18-19		B	RCL7	hR↑	*		$\mathrm{ge}^{\mathbf{X}}$	hRTN	fLN		158
(20)	fLBLE	В	Х	CHS	RCL7	+	hRTN				165
(21)	gLBLd	ſ₽ ≵ S	RCL4	RCL6	X	RCL9	RCL8	Х		RCL4	
	gx^2	RCL9	RCL5	Х	-	÷	CHS	STOA	RCL4	Х	
	RCL6	+	RCL9	÷.	fP≵S	ST07	f-x-	RCLA	STO+5	hRTN	195
RP	0	1		2	3	 	E	;	6	7	
	∆a ₀	Δa		∆a ₂	∆a ₃	1	-		Ŭ	I	
	<u> </u>	~~~	-		0]	
		∆a¦	-	∆a'2	∆a'3	∆a_	·l Δε	0	lnK	$\frac{1}{1}\Sigma\Delta H_{298}^{o}/1$	R
RS	4	5		6		7	8	9	udena da un 1965 (1974) un sun a nue a un estadorem		
	ΣТ	ΣT ²	- 5"ΛΈ	0 /R	5(AHO	~/12) ² m	/ AHO / T	<) i			
	L 1	Z1 Z	Δ.()1.	298/1	L(L)1129)8/11/ ZJ	. \ 111298/1	1) 1			
			or S-	H298-'1	ΔS_{1}		с(Дн ₂₉₈ /Г				
				R				an a			
3	А	В		С		D	E				
	$\Delta s_1^o/R$	$\frac{\Delta H_i^{\circ} - \Delta}{R}$	1H° 298	Δe		Ti	R				
				∆a_2							

Test with $C_p = a_0 + a_1T + a_2T^2 + a_3T^3$ for Mo, 1400-2800 K: (2) 1.50365 STO 0, 8.19918×10⁻³ STO 1, -4.13509×10⁻⁶ STO 2, 9.21717×10⁻¹⁰ STO 3, 0 STO C, 1.98719 STO E; (3) 1400 + 16.792 + 7307 A 1.50365; (4) SST 0; (5) SST -63.706; (6) SST -3.874; (7) SST 4.0996×10⁻³; (8) SST -6.892×10⁻⁷; (9) SST 7.681×10⁻¹¹; (10) 1400 B 11.573; (12) f e 5.824; (10) 1800 B 12.956; (10) 2200 B 14.171; (10) 2800 B 15.800; (12) f e 7.951; (11) f a 15.800. Test with $C_p = a + bT + dT^2 + cT^{-2}$ for Mo, 1400-2800 K (2) 9.54656 STO 0, -3.80258×10⁻³ STO 1, 1.6987×10⁻⁶ STO 2, 0 STO 3, -1.35954×10⁵ STO C, 1.98719 STO E; (3) 1400 + 16.792 + 7307 A 9.54656; (4) SST -67 977; (5) SST 3982.51; (6) SST -58.2878; (7) SST -1.90129×10⁻³; (8) SST 2.8312×10⁻⁷; (9) SST 0; (10) 1400 B 11.573; (12) f e 5824; (10) 1800 B 12.955; (10) 2200 B 14.168; (10) 2800 B 15.796; (12) f e 7.949; f a 15.796

Test of complete program with ΔS_1^{9} known (2) 1 STO 0, 10^{-3} STO 1, 10^{-6} STO 2, 0 STO 3, -10^{6} STO C, 1.98719 STO E; (3) $10^{3} + 20 + 100 \text{ A } 1.000$; (4) SST -5×10^{5} ; (5) SST 2733 $\frac{1}{3}$; (6) SST 10.092; (7) SST 5×10^{-4} ; (8) SST $\frac{5}{3} \times 10^{-7}$; (9) 0; (10) 10^{3} B 19.900; (10) 2000 B 20.601; (11) f a 40.939; (12) f e 20.601 (13) f c 0.000; (14) $10^{3} + 0.905$ C 19 999.8; (14) $2 \times 10^{3} + 40$ 194 C 2000.0; (16) R/S 19 999.9; (17) R/S 0.3; (18) 10^{3} D 0.905; (19) SST -0.100 (18) 2000 D 40 196; (19) SST 10.602; (20) 10^{3} E 99.918

Test of program with ΔS_{i}° to be determined (2) same as above; (3) $10^{3} \uparrow 0 \uparrow 100 \text{ A } 1.00$; (4&5) same as above; (6) SST -9.908, (7-9) same as above; (13) f c 0.000; (14) $10^{3} \uparrow 0.905 \text{ C } -0.180$; (14) $2 \times 10^{3} \uparrow 40194 \text{ C } -19 999.98$; (21) f d 19 999.6, 20.00; (10) $10^{3} \text{ B } 19.900$; (10) 2000 B 20.601

As Prgm. CgK would be repeatedly used for equilibrium calculations of the solidus-liquidus equilbria for molybdenum, it is worthwhile to use a variant of Prgm. CgK specifically designed for this type of equilibrium The program can be simplified by calculating $-(\Delta G^{\circ} - \Delta H^{\circ}_{Std})/RT$ calculation. for the fusion process at several evenly spaced temperatures and fitting these values with Prgms. Cheb (f,C_n) and Cheb-a. The resulting four a_n values are then put on a data tape to be stored in registers 0 to 3. Also $\Delta {
m H}^2_{
m St.d}/{
m R}$ for the fusion process would be on the tape to be stored in register 9. The atomic weight of molybdenum, 95.94 would be stored in register C to be used if any data in weight percent are to be converted to mole fraction. With the molefractions of Mo in the solid and liquid phases at a given temperature, the program provides values of γ_L/γ_S and $ln\gamma_L/x_2^2.$ For values at several temperatures, a least-square fit is made to estimate the excess entropy and enthalpy of mixing in the liquid phase. With these constants and for systems with moderate solid solubilities, the liquidus curve can then be calculated by the present program.

The choice of standard temperature to serve as a reference for enthalpy values and the choice of liquid standard state is sometimes found to be confusing. A number of possible choices for molybdenum were discussed in detail in Ref.(13) to demonstrate that the final calculations do not depend upon the arbitrary choices that can be made. A summary of the different types of calculations from Ref.(13) is presented here. The heat capacity data for Mo(s) at 200 K intervals from 1400 to 2800 K were fit by Prgms. Cheb (f,C_n) and Cheb-a. The heat capacity of Mo(1) was extrapolated to lower temperatures by $C_p^o/R = 2.832 + 5.923 \times 10^{-4} \text{T}$ to join the solid C_p^o at 800 K. For Mo(s) = Mo(1), $\Delta C_p^o/R = 2.075 - 3.534 \times 10^{-3} \text{T} + 2.0809 \times 10^{-6} \text{T}^2$ $-4.638 \times 10^{-10} T^3$ for 1400-2800 K. This choice of C_p for Mo(1) fixes the metastable liquid standard state between 1400 and the melting If 298.15 K is chosen as the reference temperature and $\Delta H_{\rm Std}^{0}$ is point. the enthalpy of fusion at 298 K, the method of extrapolating the liquid thermodynamic properties from 1400 K to 298 K can be completely arbitrary with no effect upon the calculations in the operating range of 1400 to 2890 K. This is most simply seen from the equation

 $\ln K = -\Delta G^{\circ}/RT = -(\Delta G^{\circ} - \Delta H_{std}^{\circ})/RT - \Delta H_{std}^{\circ}/RT = \emptyset - (\Delta H_{i}^{\circ} - \Delta H_{std}^{\circ})/RT - \Delta H_{std}^{\circ}/RT$

where \emptyset represents those terms of $-(\Delta G^{\circ}-\Delta H^{\circ}_{\rm Std})/RT$ that depend only upon the heat capacity contributions and ${\mathbb T}_{{\mathfrak f}}$, the bound to the range of validity of the heat capacity equation, and not upon T_{Std}. It is apparent that any change in the standard reference temperature or in the methods of extrapolating to the standard reference temperature cancels out due to an equal change in $(\Delta G - \Delta H_{Std}^{\circ})/RT$ and in $\Delta H_{Std}^{\circ}/RT$. This was illustrated in Ref.(13) by using different heat capacity equations for liquid Mo between 298 and 1400 K and as expected from the above equation, identical values of ln K are obtained in the range 1400-2890 K. Also the use of 2890 K instead of 298 K as the standard reference temperature gave identical ln K values. Likewise, $T_{Std} = 0$ K changes $(\Delta H_1^0 - \Delta H_{Std}^0)/RT$ and $\Delta H_{Std}^0/RT$ equally. If the same standard reference state, e.g. Mo(s) at 298.15 K, is used for both the solid and liquid phases such that $\Delta H_{298}^{2} = 0$ for Mo(s) = Mo(l), $\Delta H_{1400}^{0} - \Delta H_{Std}^{0}$ is changed correspondingly to yield the same $\Delta G^{\circ}/RT$. The various quantities are shown for comparison where A corresponds to extrapolation of the liquid C_p equation from 1400 to 298 K; B takes ΔC_p = 0 from 800 to 298 K, C refers both liquid and solid to the solid reference state at 298 K and D uses 2890 K as the standard reference temperature.

	А	В	С	D
^T Std	298.15	298.15	298.15	2890 К
$\Delta H_{Std}^{o}/R$	5305	5323	0	4303 K
$(\Delta H_{1400}^{\circ}-\Delta H_{Std}^{\circ})/R$	2	17	5307	1005 K
$-(\Delta G_{1400}^{\circ}-\Delta H_{std}^{\circ})/RT$	1.904	1.916	-1.885	1.188
$-\Delta G_{1400}^{\circ}/RT$	-1.886	-1.886	-1.886	-1.885

Within the limit of rounding errors, the final results are identical.

Since there is no difference, Prgm. Ωxb is set-up to use the conventional standard reference temperature of 298.15 K but a wider temperature range is handled than previously (13). For the range 900 K to the melting point, $C_p^o/R = 2.18 + 2.0484 \times 10^{-3}T - 1.1049 \times 10^{-6}T^2 + 0.31558 \times 10^{-9}T^3$ was used for the solid. For the liquid, $C_p^o/R = 2.832 + 5.923 \times 10^{-4}T$ was used to extrapolate from the value of 4.544 at 2890 K to join the solid value of 3.306 at 800 K. For Mo(s) = Mo(1), $\Delta C_p^o/R = 0.651 - 1.456 \times 10^{-3}T + 1.1049 \times 10^{-6}T^2$ -0.31558 $10^{-9}T^3$ from 900 K to the melting point. ΔC_p^o was taken as zero below 900 K. $\Delta H_{2890}^{2}/R = 4303$, $\Delta H_{900}^{8}/R = \Delta H_{298}^{2}/R = 5329$ K and $\Delta S_{900}^{6}/R = -(\Delta G_{900}^{8} - \Delta H_{298}^{2})/RT = 1.921$.

The constants of the ΔC_p° equation were inserted in Prgm. CgK with $T_i = 900$, $\Delta S_i^{\circ}/R = 1.921$, and $(\Delta H_1^{\circ} - \Delta H_{298}^{\circ})/R = 0$. Values of $-(\Delta G^{\circ} - \Delta H_{298}^{\circ})/RT$ were calculated ranging from 900 K to the melting point and fitted with Prgms. Cheb. $C_n - (f, C_n)$ and Cheb-a3 to obtain

 $-(\Delta G^{\circ} - \Delta H_{298}^{\circ})/RT = 1.9357 - 3.544 \times 10^{-5}T + 3.123 \times 10^{-8}T^{2} - 1.042 \times 10^{-11} T^{3}$ for Mo(s) = Mo(1) between 900 K and the melting point with an average deviation of less than 0.001 compared to the values from Prgm. CgK.

$$\begin{split} & \mathrm{K} = \ \mathrm{x}_{Mo}^{\mathrm{L}} \gamma_{Mo}^{\mathrm{S}} / \mathrm{x}_{Mo}^{\mathrm{S}} \gamma_{Mo}^{\mathrm{S}} \text{ and } \gamma_{Mo}^{\mathrm{L}} / \gamma_{Mo}^{\mathrm{S}} \text{ are calculated from the equation} \\ & \mathrm{lnK} = \ -\Delta \mathrm{G}^{\circ} / \mathrm{RT} = -(\Delta \mathrm{G}^{\circ} - \Delta \mathrm{H}_{298}^{2}) / \mathrm{RT} - \Delta \mathrm{H}_{298}^{2} / \mathrm{RT} = \ \mathrm{ln} \gamma_{Mo}^{\mathrm{L}} / \gamma_{Mo}^{\mathrm{S}} + \ \mathrm{ln} \ \mathrm{x}_{Mo}^{\mathrm{L}} / \mathrm{x}_{Mo}^{\mathrm{S}}. \end{split}$$
From just the liquidus curve, it is not possible to unambigously fix the enthalpy of mixing and the excess entropy of mixing for the solution. However, the approximation of assuming that both the enthalpy of mixing and the excess entropy of mixing vary as $(1-x_{MO})^2$ is made for preliminary treatment of the data $\Delta \overline{H}_{MO}/R = b_h (1-x_{MO}^L)^2$, $\Delta \overline{S}_{MO}^E/R = +b_s (1-x_{MO}^L)^2$, and $\ln \gamma_{MO}^{L} = \Delta \overline{G}_{MO}^{E}/RT = (b_{h}/T - b_{s})(1 - x_{MO}^{L})^{2}$ where the partial molal quantities are for the mixing of liquid Mo with the other component. At a given temperature, $b_h/T - b_s$ will be designated as Ω . When x_{Mo}^S and γ_{Mo}^S are close to unity, $\ln \gamma_{Mo}^{S}$ will be very small and will be approximated by $\ln \gamma_{Mo}^{S} = [\ln(\gamma_{Mo}^{L}/\gamma_{Mo}^{S})](1-x_{Mo}^{S})^{2}/(1-x_{Mo}^{L})^{2}$. If the resulting $\gamma_{Mo}^{S}x_{Mo}^{S} > 1$, γ_{Mo}^{S} is approximated by $\frac{1}{x_{Mo}^{S}}$. Then multiplication of $\gamma_{Mo}^{L}/\gamma_{Mo}^{S}$, determined in the equation above from the liquidus and solidus points and $\Delta G^{\circ}/RT$ of fusion, by γ_{MO}^{S} will yield γ_{MO}^{L} which can then be converted to Ω or $b_{h}/T - b_{s}$. From values at several temperatures, a least square fit will yield the values of b_h and b_s from which the first approximation to the enthalpy and excess entropy of mixing of liquid Mo with the other component will be available. Then using bonding theory to fix reasonable enthalpy and entropy values, a second approximation to $\mathbf{b}_{\mathbf{h}}$ and $\mathbf{b}_{\mathbf{s}}$ is obtained that fits both the liquidus data and provides reasonable enthalpy and entropy values. An approximate liquidus curve can be calculated using the initial b_h and b_s with an estimated solid solubility. The program also provides for converting Celsius temperatures to Kelvin and converting weight percent to molefraction.

Directions:

(1)	Insert Prgm. Ωxb
(2)	Insert data tape with constants of equation for $-(\Delta G^{\circ}-\Delta H^{2}_{298})/RT$ and ΔH°_{298} for the fusion of Mo and the atomic weight of Mo.
(3)	If any weight percent data are to be treated, store atomic weight of second component in register C.
(4)	E (to preceed each new set of data)Display0.000
(5)	$T \uparrow (1-x_{MO}^{S}) \uparrow (1-x_{MO}^{L}) A$ $\Omega, \gamma_{MO}^{L} / \gamma_{MO}^{S}$
	Repeat step 5 for data at each T
(6)	R/S after all data entered in step 5 b_h, b_s
(7)	$(1-x_{M_{O}}^{S})^{\uparrow}(1-x_{M_{O}}^{L})B$ values will flash until T within 2°
(8)	$T f a -\Delta G_{f}^{o}/RT$
(9)	t C T
(10)	wt.%A D $x_A = (1-x_{MO})$
NOTE	1: Step 9 is used to preceed steps 5, 7 and 8 when t°C is to be converted to T°K.
NOTE	2: Step 5 calls for $x_A = (1-x_{MO})$. When wt.%A is available, step (10) converts to x_A . Starting with wt.%Mo, key 100 wt.%Mo - D for step 5.
NOTE	by step 4 which clears the secondary registers. At any time after completion of step 5 and before step 4 is carried out to prepare for a new set of data, step 6 can be repeated by keying f c.
NOTE	rough estimate can be used. $x_A^{s=10^{-4}}$ or 10^{-5} can be used for all very small solid solubilities. The calculation is inadequate for large
NOTE	5: Solid solubilities. The values of b _h and b _s obtained in step (6) are sensitive to experimental error. After modification to match reasonable enthalpies and excess entropies, the revised b _h is stored in R5 and the revised b _s is stored in RI; step 7 will then use these values for calculation of the liquidus curve.
NOTE	6: To use for another component, store $\Delta H_1^o/R$ in A, $\Delta S_1^o/R$ in O, and O in 1-3 if $\Delta C_p = 0$. In steps 5 and 7, x_{MO} is replaced by x_2 . When $\Delta C_p \neq 0$, the next level of approximation is to use the average values of $\Delta H_1^o/R$ and $\Delta S_1^o/R$ over the temperature range. For Mo, one can use $\Delta H_1^o/R = 5127$ and $\Delta S_1^o/R = 1.774$ as averages for 1400-2890 K.

Prgm	. Ωxb								1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 - 1997 -		
(9)	fLBLC	RCLD	+	hRTN							4
10)	(LBLD	EEX	\hat{c}	:	hl/x	1	-	RCLC	Х	RCLB	
	•	1	+	hl/x	hR'I'N						19
(4)	f'LBLE	f₽ ≵ S	0	STO4	ST05	ST06	STO8	STO9	f₽ ∠ S	hRTN	
(5)	fLBLA	STOE	1		hx ≵y	hSTI	1	-	-	fLN	
	STQ6	hR↓	STO4	fa	RCL6	-	STO9	hRCI	RCLE	÷	
	gx^2	Х	个	1	1	hRCI	-	fLN	+	fx>0	
	fGSB1	hR↓	STO8	RCL9	+	RCLE	gx^2	÷	ST07	RCL4	
	Х	hLSTx	∑+	hLSTx	÷	f-x-	RCL9	ge^X	hRTN		78
(6)	gLBLfc	f₽ ≵ S	RCL4	rcl6	Х	rcl8	RCL9	X	-	RCL4	
	gx ²	RCL9	RCL5	Х		÷	CHS	hSTI	RCL4	Х	
	rcl6	+	RCL9	÷	f₽żS	ST05	f-x-	hRCI	hRTN		107
(7)	fLBLB	1	STO8	STO9	hR↓	ST0-9	gx ²	hx≠y	STO-8	gx ²	
		↑	1	RCL5	X	RCLA	+	ST07	RCL8	RCL9	
	÷	fLN	hR↑	hRCI	X	+	STOG	RCLO	+	*	
	ſLBL3	f-x-	fa	RCLA	hR↑	÷	+	RCL7	hx ≑ y	RCL6	
	+		STO4	_	hABS	2	gx≤y	GTO4	RCL4	hRTN	
	fLBL4	RCL4	GTO3				0				16
(8)	gLBLfa	↑	 ↑	↑	RCL3	Х	RCL2	+	Х	RCL1	101
	+	Х	RCLO	+	RCLA	hR↑	÷.		hRTN		179
	fLBL1	1	hRCI	-	fLN	CHS	†	hRTN			187
RP	0	1	2	3	4	5	6	7	8	9 .	
	a ₀	0	0.5	6	T	Ъ	л " ^х ́Мо	Ω	ln $\gamma^{ m S}_{ m Mo}$	$\gamma_{\rm N}^{\perp}$	Ío
	all	al	a2	a3	Ţ	b _h	ln <u>Mo</u> xMo	. 90	In YMo	$\ln \frac{r_{\rm N}}{\gamma_{\rm N}^2}$	3
							useđ	use	d x _{Mo} S	\mathbf{x}_{M}^{I}	i Io
RS	4	5	6	7	8	9			1.10	I ^v .	
	ΣT	ΣT^2	ΣΩΤ	•	T ² ΣΩΤ						
R	A	B	C	D	E	I					
	ΔH ^o 298/R	95.9	4 M _A	273.	15 x	A x	S A				
						ď	g				

Test:

- (3) 157.26 STO C
- (9) 150 C 423.15; (10) .09 D 5.49×10⁻⁴, 1.53 D 9.39×10⁻³;
- (4) E 0.000; (5) $2800 + 10^{-3} + .0553 \text{ A } 1.715$, 1.005;
- (5) $2700 + 10^{-4} + .132 \text{ A} 1.741$, 1.031; (5) $2600 + 10^{-5} + .2522 \text{ A} 1.769$, 1.119;
- (5) $2500 + 10^{-6} + .54202 \text{ A} 1.800, 1.697;$ (5) $2400 + 10^{-6} + .7315 \text{ A} 1.833, 2.667;$
- (6) R/S 1977.8, -1.009; (7) 10^{-3} + .0553 B 2683, 2783, 2797, 2800, 2800;
- (7) 10^{-5} + .2522 B 2523, 2592,2599,2600; (8) 2700 f a -0.111;
- (6) f c 1977.8, -1.009

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Calculation of Composition Boundaries of Saturating Phases

Such calculations require that Gibbs energy of solution be expressed in analytical form as a function of temperature and composition. The functional form of the regular solution equation is the most practical for most systems.

The regular solution derivation^{12,20} gives the Gibbs energy in terms of volume fraction. If the volume fraction is expanded as a truncated power series in mole fraction, one obtains

 $\Delta \overline{G}_1^E/RT = b_g x_2^2 + c_g x_2^3 \qquad \text{and} \qquad \Delta \overline{G}_2^E/RT = [b_g + \frac{3}{2}c_g]x_1^2 - c_g x_1^3$ $\Delta \overline{G}_1^E \text{ is the excess partial molal Gibbs energy of component 1 and V_1 is the molal volume of component 1. As the regular solution equation is derived under constant volume conditions, an excess entropy term arises from the volume change upon mixing. In addition, since <math>b_g = (V_2^2/V_1)[\Delta E_1/V_1)^{\frac{1}{2}} - (\Delta E_2/V_2)^{\frac{1}{2}}]^2/RT$, the temperature coefficients of the molal volume and of the energy of vaporization, ΔE , also result in an excess entropy.¹² The enthalpy and excess entropy can be expressed as similar functions of mole fraction. Their combination to give the partial molal Gibbs energy equation yields for each component in its standard state dissolving in the solution

$$\ln \gamma_{1} = \Delta \overline{G}_{1}^{E} / RT = (b_{h} / T - b_{s}) x_{2}^{2} + (c_{h} / T - c_{s}) x_{2}^{3} = \Omega_{1} x_{2}^{2}$$

$$\ln \gamma_{2} = \Delta \overline{G}_{2}^{E} / RT = [b_{h} / T - b_{s} + \frac{3}{2} (c_{h} / T - c_{s})] x_{1}^{2} - (c_{h} / T - c_{s}) x_{1}^{3} = \Omega_{2} x_{1}^{2}$$

where the signs of the b_h , c_h , b_s and c_s terms are the same as the signs of the corresponding enthalpy and entropy terms in the Gibbs energy equation. Thus, b_gT at a temperature T equals b_h-b_sT , etc. These equations average out the contributions of ΔC_p values to the Gibbs energy by using the average enthalpy and entropy values over the temperature range of interest. Analytical equations of this form are found to reproduce, within the range of experimental uncertainty, the thermodynamic properties of many metallic solutions over a considerable range of temperature and composition.

When a miscibility gap exists in the solution at a given temperature, the partial molal Gibbs energies of both components are equal in both saturating phases. If the mole fractions are expressed as $x_1 = 1-y$ and $x_2 = y$ in the phase with excess component 1 and as $x_1 = x$ and $x_2 = 1-x$ in the phase with excess component 2, the equilibrium condition can be expressed as

$$\ln x + b(1-x)^{2} + c(1-x)^{3} = \ln(1-y) + by^{2} + cy^{3}$$

$$\ln y + (b + \frac{3}{2}c)(1-y)^{2} - c(1-y)^{3} = \ln(1-x) + (b + \frac{3}{2}c)x^{2} - cx^{3}$$

where $b = (b_h/T-b_s)$ and similarly for c. When the same values of b and c are not applicable for both phases, the equations become

 $\ln x + b_{x}(1-x)^{2} + c_{x}(1-x)^{3} + e_{x} = \ln(1-y) + (b_{y} + \frac{3}{2}c_{y})y^{2} - c_{y}y^{3} = -d_{x}$ $\ln y + b_{y}(1-y)^{2} + c_{y}(1-y)^{3} + e_{y} = \ln(1-x) + (b_{x} + \frac{3}{2}c_{x})x^{2} - c_{x}x^{3} = -d_{y}$

The subscript x indicates the constants apply in the phase region for which $x = x_1$ is small and y indicates applicability in the phase region for which $y = x_2$ is small. These general equations may be reduced to the first set if $b_x = b$, $b_y = b + \frac{3}{2}c$, $c_x = c$, and $c_y = -c$. When the solubilities are small, the cubic terms can be dropped with the constants of the quadratic terms being $\Omega_x = b_x + c_x$ and $\Omega_y = b_y + c_y$. The e_x and e_y terms are discussed below.

For a symmetrical immiscibility gap when the molal volumes of the two components are closely the same, c = 0 and $b_x = b_y$ resulting in equal values of x and y. The two equations reduce to $\ln[x/(1-x)] + b(1-2x) = 0$ which is readily solved by Prgm. Im. If the excess entropy is neglected, b will be inversely proportional to T and the calculations can be repeated at a new temperature T_2 by multiplying b by T_1/T_2 . This can be done by $T_1 \uparrow T_2 \div$ STO X 4 RCL 4 f b to obtain the value of x = y at T_2 . The temperature that brings b down to 2 will be the critical temperature. For b < 2, the system will be miscible.

Program Im carries out the simultaneous solution of these two equations for x and y to yield the compositions of the saturating phases. The iterative solution using the secant method is based in part on a portion of Prgm. SD-11 of the HP-67 Standard Pac(18). Program Im first assumes that the solubilities are small enough to neglect d_x and d_y and to approximate 1-x by 1. This yields as the first approximation $x' = e^{-(b_x+c_x)}$. x' is used to calculate the first approximation of d_y .

 $\ln y + b_y(1-y)^2 + c_y(1-y)^3 + d_y + e_y = 0$ is then solved by iteration to yield y". If flag 0 is set, the value of y obtained by each iteraction is flashed. The process continues until the fractional change in y in the last iterative step rounded to the number of places after the decimal designated in step 2 is displayed as zero. The same process is then used in step 7 to calculate x". The value of x" is used to calculate a new value of d_y and the equation for y is solved again to obtain y". Steps 6 and 7 are alternately repeated until the values of x and y show no additional change. To speed convergence, step 2 should be set initially at DSP 2 and no change in x and y will be noted after $\Delta x/x$ or $\Delta y/y < 0.005$. For x and y ~0.1, the values of x and y will be accurate to 0.0005 and can be read to 4 places by keying DSP 4. If it is desired to calculate x and y more accurately, steps 7 and 6 can be repeated with DSP 3 or 4. If there is no interest in the progress of the iteration at each step, hCFO will stop the flashing of successive x or y values. hSFO will change back to flashing.

If the same standard states are not used for both phases, a constant term would be added to d_x or d_y or both corresponding to the Gibbs energy difference between the standard states divided by RT. The equilibrium between the solidus and liquidus portions of a phase diagram will be considered as an example. For solid with largely component 1 in equilibrium with the liquid phase, then $e_x = \Delta G_{f,1}^{\circ}/RT = \Delta H_{f,1}^{\circ}/RT - \Delta S_{f,1}^{\circ}/R$ and $e_y = -\Delta G_{f,2}^{\circ}/RT = -\Delta H_{f,2}^{\circ}/RT + \Delta S_{f,2}^{\circ}/R$, where ΔH_f° and ΔS_f° are the average enthalpies and entropies of fusion of the two components over the temperature range of interest. In the equations given on the previous page, b_x and b_y will be considerably different, in general, for the example of solid-liquid equilibria.

Program Im provides for insertion of the enthalpy and entropy terms to allow calculation of b_x , c_x , e_x , b_y , c_y and e_y at various temperatures and then to solve for x and y values for the saturating phases at the given temperature. When the b, c and e values are already known for the desired temperature, the enthalpy values are inserted as zero and the b, c and e values are inserted with reversed sign for the entropy terms as illustrated in one of the test examples.

For mutual solubilities of two liquid phases or two solid phases for which there are no e_x or e_y terms due to differences in standard states, $\Delta H_{f,1}^{0}\Delta S_{f,1}^{0}\Lambda H_{f,2}^{0}$ and $\Delta S_{f,2}^{0}$ are entered as zero. When the difference in standard states corresponds to the solid phase transition, then ΔH_{f}^{0} and ΔS_{f}° are replaced by the corresponding values for the transition. For liquid immiscibility with small enough solubilities for the approximation $\ln \gamma_{1} = (b_{x}+c_{x})y^{2}$ and $\ln \gamma_{2} = (b_{x}+\frac{1}{2}c_{x})(1-y)^{2}$ in the phase rich in component 2, the regular solution theory predicts $(b_{x}+c_{x})/(b_{x}+\frac{1}{2}c_{x}) = (V_{2}/V_{1})$.

This is sometimes useful, but, in practice, the value of c is often more dependent upon change in character of the bonding across the solution range than upon the volumes of the pure materials, especially with change in electronic configuration upon solution. One can sometimes obtain a useful value of c by comparison of $(\ln\gamma_1)/x_2^2 = b+c$ at large x_2 with $(\ln\gamma_2)/x_1^2 = b+\frac{1}{2}c$ at large x_1 . If c = 0 is used in Prgm. Im, $(\ln\gamma_1)/x_2^2$ and $(\ln\gamma_2)/x_1^2$ are taken equal for a given solution, but their value can be greatly different on opposite sides of the miscibility gap.

Directions:

(1)	Insert	Prgm. Im				
(2)	DSP n,	usually n=2 init	cially		Display	
(3)	ΔH° _{f,l} /R	$\wedge \Delta S_{f,1}^{o}/R \wedge b_{h}^{x}$	↑ b ^x _s A		$\frac{\text{Display}}{\Delta H_{f,1}^2/R}$	
		$R \uparrow -\Delta S_{f,2}^{2}/R \uparrow t$		fa) -	$\Delta H_{f,2}^{o}/R$	
		$\uparrow e_h^x \uparrow e_s^x R/S$			c_{h}^{y}	
(6)	ТВ (or t ^o c D)		x	, y' y"	
(7)	С	Repeat 7 and 8	alternately un	til x"	x ^{***}	
(8)	R/S	x and y show no sive steps; rea	o change in suce ad final values	ces- y"		
		with DSP n+1				
	lf y is	known; after st	tep 7, DSP n to	desired	accuracy followed	by
(81)	y sto A	fd			$d_{\mathbf{x}}$	
(7 ')	С			x'	' x _{final}	
	<u>If x is</u>	known; after st	ep 6, DSP n to	desired	accuracy followed	by
(7")	x STO A	fd			dy	
(8")	С			у'		
NOTE	l: The	alternatives fa	a and fb for R/	S in ster	os 3 and 4 allows a	alteratio

- NOTE 1: The alternatives fa and fb for R/S in steps 3 and 4 allows alteration of part of the data without need to reinsert all of the data. For step 6, temperature in °C can be inserted followed by D to convert to °K and initiate the calculations.
- NOTE 2: The display control of step 2 can be changed at any time, but is best started with only two places. If higher than 3 place accuracy is desired, n can be increased at each successive repeat of steps 7 and 8. For each new T, start at step 6 following step 8. If step 7 was the last step, key fP₂S before starting again at step 6. cont'd.

- NOTE 3: C can be used in place of R/S for step 8, but alternating between C for the x calculation and R/S for the y calculation helps one keep track of which value is being calculated.
- NOTE 4: If it is desired to repeat step 7' to obtain a value to a larger number of places, change DSP n and repeat steps 8' and 7'. Similarily, if x is known, repeat steps 7" and 8".
- NOTE 5: After T STOD, E calculates $\Delta G_{1,1}^o/RT$, $\sqrt{calculates} b_x$ and fc calculates c_x for component 1 if the component 1 constants are in the primary registers. If $d_x = -\ln a_1$ is desired, it may be obtained by RCL5 after step 8.
- NOTE 6: Step 6 assumes x and y are small. If they are not, replace step 6 by T STO D x' [†] y' GTO 8 R/S, where x' and y' are estimates of the solution. Then continue with steps 7 and 8. This procedure is also used for calculating solubilities for a sequences of temperatures. If the temperature intervals are small, only T STOD is needed for step 6.

TEST:

- (3) $5127 + 1.774 + 4 \times 10^3 + .3 \text{ A } 5127.00;$
- (4) -1371 ↑ -.762 ↑ 8×10³ ↑ .2 R/S -1371.00;
- (5) 500 ↑ .1 ↑ 500 ↑ .1 R/S 500.00;
- (6) 2×10^3 B (or 1726.85 D) x=.07, y=.02, .02, .02, .02;
- (7) C x=.10, .10, .10, .10;
- (8) R/S y=.02, .02, DSP 4, .0190;
- (7) C x=.1014, .1014, .1014, DSP 5, .10138;
- (8) R/S y=.01902, .01902, .01902, DSP 6, .019022.

y = .019022 known

Steps 2-7 the same as above; (8') .019022 STO A fd, d_x =.0177;

(7') C .1014, .1014, .1014, DSP 5, .10138

x = .10138 known

Steps 2-6 the same as above; (7") .10138 STO A fd, $d_v = .0873$;

(8") C .0190, .0190, .0190, .0190, DSP 6, .019022

$$e_x = e_y = 0, b_g^x = 2.5, b_g^y = 3, c_g^x = .5, c_g^y = .1$$

- (3) 0 ↑ 0 ↑ 0 ↑ -2.5 A 0.00
- (4) $0 \uparrow 0 \uparrow 0 \uparrow -3 R/S 0.00$
- (5) $0 \uparrow -.1 \uparrow 0 \uparrow -.5 \text{ R/S} 0.00$
- (6) 1 B x=.05, y=.04, .06,.06,.06
- (7) C.08, .08, .08, DSP 3, .075; (8) R/S.062, .062, .062, DSP 4, .0618
- (7) C .0755, .0755, .0755, <u>.0755</u>

Dragm	Tm										
Prgm											
(3)	ſLBLA	STO7	hR↓	ST06	hR↓	ST04	hR↓	ST03	hRTN		9
(4)	gLBLfa	f₽ ≵ S	А	fP≵S	hRTN				0		14
(5)	gLBLfb	STO9	hR↓	STO8	hR↓	fP≵S	STO9	hR↓	STO8	fP≵S	25
(C)	hRTN GLDID	amon	0			111		atra	х	1 000	2 3
(6)	fLBLB	STOD	fc	fGSB6	+	Е	+	CHS	$\operatorname{ge}^{\mathbf{x}}$	hPSE	
	fGSBO +	fGSB1 CHS	RCL1 ge ^x	STOA hPSE	fd fGSB0	fc fGSB1	+ GTOC	fgsb6	+	Ε	52
	fLBL0	t t	ge ↑	nrse 2	IGBBU ÷	STOO	6100 +	STOL	hRTN		
(7,7		fc	hSTI	rgsb6	STOE	fGSB2	fGSB3	STOL	fGSB1	RCLA	61
(8,8)		RCLA	hRTN	GTOC	010D	TUDDE	TODDO	DION	TOODT	nom	75
(6')	fLBLD	2	7	3	•	1	5	+	GTOB		84
. ,	fLBL1	RCL1	RCLO	-	STO2	2	÷	-	STOO	hRTN	
	gLBLfd	fc	X.	hLSTx	1	•	5	Х	fGSB6	+	
	-	RCLA	gx ²	Х	1	RCLA	-	fLN		fP≵S	
	ST05	hRTN									116
	fLBL6	RCL6	RCLD	÷	RCL7		hRTN				123
	gLBLfc	RCL8	RCLD	÷	RCL9	-	hRTN				130
	fLBL2	RCL1	fe	STOC	RCLO	fe	STOB		hRTN		138
	fLBL3	RCLC	GTO4	fLBL5		fe	STOC	fLBL4		RCL1	
	-	RCLC	RCLB		÷	X	STO+1	RCL1	hF?O	hPSE	164
	÷ gLBLfe	fRND fr N	fx≠0	GTO5 hLSTx	RCL1	hRTN gx ²	hLSTx	hRCI	Х		104
	RCLE	fLN +	l X	пгетх +	- RCL5	gx- +	nlisi x E	4	n hRTN		182
	fLBLE	RCL3	RCLD	÷	RCL4		hRTN	T	1117.17.14		189
	fLBL8	STOB	hR↓	fGSB9		fGSB9	hRTN				
	fLBL9	↑	STOA	STOO	EEX	CHS	4	ST02	+	ST01	196
	hR↓	fd	hRTN	0100		0110	1				2 U 9
	fLBL7	fc	RCLA	Х	fgsb6	+	RCLA	gx ²	Х	ge^{X}	205
	ŔŢŊ							0		0	220
R	0	1	2	3	<u>ц</u>	5		6	7	8	9
р	x'	т х"	Δx		∆s°/)			x bh	b _s ^x	c_h^x	c _s x
				-1-		•				∽h v	∽s v
S	у'	у"	Δy	-ΔH ₂ /R	-Δs ₂ /1	R d	y	Þh	b ^y s	e _h y	e ^y s
А	B (C I)	E	I	Flag	0 is se	t			
x ***	f(y') f(y	r")	Ľ	by	ey						
y III	f(x') f(x)			•	•						
у 		× /		b _x	°x	Shing to got the state of the s					

Gibbs Energy from Phase Boundaries

When the compositions of saturating phases are known, two constants can be directly determined. If the same b and c values are assumed for both phases for which the same standard states are used,

 $\begin{aligned} &\ln x + b(1-x)^2 + c(1-x)^3 = \ln(1-y) + by^2 + cy^3 = \ln a_1 = \Delta \overline{G}_1 / \text{RT} \\ &\ln y + (b + \frac{3}{2}c)(1-y)^2 - c(1-y)^3 = \ln(1-x) + (b + \frac{3}{2}c)x^2 - cx^3 = \ln a_2 = \Delta \overline{G}_2 / \text{RT} \end{aligned}$ The mole fraction of component 1 is x in the phase rich in 2 and is 1-y in the phase rich in 1. The simultaneous solution of these equations expressed as $\alpha_1 b + \beta_1 c = \varepsilon_1$ and $\alpha_2 b + \beta_2 c = \varepsilon_2$ is discussed on pg. 114 of reference 13 using the quantities

 $\begin{array}{rcl} \alpha_{1} &=& 1\,-\,x^{2}/(1-y)^{2}, & \alpha_{2} &=& 1\,-\,x^{2}/(1-y)^{2}, \\ \beta_{1} &=& 1\,-\,x\,-\,y^{3}/(1-x)^{2}, & \beta_{2} &=& \frac{1}{2}\,+\,y\,-\,\frac{3}{2}x^{2}/(1-y)^{2}\,+\,x^{3}/(1-y)^{2}, \\ \epsilon_{1} &=& [\ln(1-y)/x]/(1-x)^{2}, & \epsilon_{2} &=& [\ln(1-x)/y]/(1-y)^{2}, \\ \text{and} & D &=& \alpha_{1}\beta_{2}\,-\,\alpha_{2}\beta_{1}. \end{array}$

When the same b and c values are not expected to be applicable for both phases or if the solubilities are small enough to combine the cubic terms with the square terms, then one uses $\ln\gamma_1/(1-x)^2 = b_x$ and $\ln\gamma_2/(1-y)^2 = b_y$. In addition to calculating values of b and c or b_x and b_y , values of $\ln\gamma_1 = \Delta \overline{G}_1^E/RT$ and $\ln a_1 = \Delta \overline{G}_1/RT$ and the corresponding values for the other components are calculated for both phases. If a set of b and c values or b_x and b_y values have been obtained at several temperatures, steps 4-6 of Prgm. Ωxb can be used to obtain enthalpy and excess entropy terms corresponding to $b = b_h/R + b_s$ or $c = c_h/T + b_s$ and similarily for b_x and b_y . The steps are (4) E, (5) T \uparrow b for each b value, then (6) f c will yield b_h and b_s . The same steps are repeated to obtain c_h and c_s values from the set of c values.

Directions:

(1)	Insert Prgm. yx-bc	Display
(2)	y ↑ x A	b
(3)	SST	c
(4)	у↑хВ	b _x
(5)	SST	b _y
(6)	С	$ln\gamma_1, lna_1, ln\gamma_2, lna_2$ for phase rich in 2
(7)	D	$ln\gamma_2$, lna_2 , $ln\gamma_1$, lna_1 for phase rich in l

NOTE 1: b and c can be obtained after step (5) by RCLE hRCI.

NOTE 2: Steps 6 and 7 can be used in either order following step 3 or

step 5.

Prgm. yx-bc Test:

(2) .095 ↑ .07 A 2.4837; (3) SST 0.5397;
(4) .04 ↑ .03 B 3.6889; (5) SST 3.4635;
(6) C 3.4709, -0.0357, 0.0036, -0.0269;
(7) D 3.1920, -0.0269, 0.0051, -0.0357;
after (2), (6) C 2.5823, -0.0769, 0.0160, -0.0566;
(7) D 2.2973, -0.0566, 0.0229, -0.0769.

Prgm. yx-bc

(2)	PT DT A	CITIO IT	٦		and	omok]	C III C IIII C III	-1	
(2)	fLBLA CHS	STO5 STO4	l RCL5		CHS ÷	STO6 STO8	hx ≵y hR	STO3 RCL3	l hx≠y	Walt Wat S
	STO7	RCL6		hx≯y 5	•	RCL8	gx^2	X	RCL4	+
	CHS	1	٠	5	+	STOD	1	RCL7	3	hy^x
		RCL6	x	STOC	RCL7	fLN	CHS	RCL4	g_{x}^{2}	11.y ÷
	ST02	1	RCL7	gx^2		STOA	RCL8	fLN	CHS	RCL6
	gx ²	. <u></u>	STO1	1	RCL8	gx ²		STOB	X	hx _y
	RCL2	X	-	RCLC	RCLB	X	RCLA	RCLD	X	
	STO9		hSTOI	RCLC	RCL2	X	RCL1	RCLD	X	-
(3)	RCL9	<u>.</u>	STOE	hRTN	hx≵y	hRTN				96
(4)	fLBLB	А	+	hRCI	RCL5	X	***	STOC	RCL3	•
(5)	5	+	hRCI	Х	RCLE	+	STOD	RCLC	hRTN	hx≵y
(6)	fLBLC	hRCI	RCL6	Х	RCLE	+	RCL6	gx^2	Х	STOO
	RCL5	fLN	+	fa	RCL5	Х	e Sant	RCL5	gx^2	Х
	个	STOl	rcl6	fLN	÷	RCLO	hR↓	gSTK	hRTN	145
(7)	fLBLD	fa	RCL4	Х		RCL4	gx^2	Х	STO2	RCL3
	fLN	+	hRCI	RCL3	Х	RCLE	+		gx2	X
	个		RCL4	fLN	+	RCL2	hR↓	gSTK	hRTN	174
	gLBLfa	RCLE	1	٠	5	hRCI	X	+	hRC1	hRTN ¹⁸⁴
0	1	2		3	λį	5	6	7	8	9
	εl	ε	2	У	l-y	x	l-x	y/(1-x) x/(1-y) D
$ln\gamma_1^{(2)}$	2) $\ln \gamma_2^2$	2)	(1)							$ln\gamma_1^{(1)}$
ınγı	τnγ2	In	Y2							$\ln \gamma_{1}$
A	В	С		D	E	I				
0.7										
α	α2	β	1_	β2	Ъ	С				
		b	x	by						

Calculation of Thermodynamic Properties of Liquid Solution from Properties of Solid Solution and Solidus-Liquidus Data and Effect of Errors

Program Im presents the equations for the solidus-liquidus equilibrium. The same equations and symbols will be used here for Prgm. Ω yxbc. This program starts with enthalpies and entropy of fusion of both components and the partial molal enthalpies and entropies in the solid phase and calculates

$$\Omega_1^{\ell} = \ln \gamma_1^{\ell} / x_2^2$$
 and $\Omega_2^{\ell} = \ln \gamma_2^{\ell} / x_1^2$

for the liquid phase of composition $x_1^{\ell} = x$ in equilibrium with the solid phase of composition $x_2^{S} = y$ and then calculates b_x and c_x of $\Omega_1^{\ell} = b_x + c_x(1-x)$. It is often of interest to determine the influence of error in x and y on the calculated values of Ω_1 and Ω_2 . This program provides values of $d\Omega/dx$ and $d\Omega/dy$ for both components. The variation of each Ω with the change of b_y and c_y can also be calculated. An error in the entropy of fusion of a component causes an error only in Ω^{ℓ} for the same component. The program gives the error in Ω_2^{ℓ} for an error in $\Delta S_{f,2}/R$.

$$\begin{split} &\ln a_{2}^{S} = \ln y + b_{y}(1-y)^{2} + c_{y}(1-y)^{3}, \quad \ln a_{2}^{\ell} = \ln(1-x) + (b_{x} + \frac{3}{2}c_{x} - c_{x}x)x^{2} \\ &\ln a_{2}^{\ell} - \ln a_{2}^{S} = -\Delta G_{fus,2}^{\circ}/RT = e_{y} \\ &\ln a_{2}^{S} + e_{y} - \ln(1-x) = \ln a_{2}^{\ell} - \ln(1-x) = (b_{x} + \frac{3}{2}c_{x} - e_{x}x)x^{2} = \Omega_{2}^{\ell}x^{2} \\ &\Omega_{2}^{\ell} = [\ln y + b_{y}(1-y)^{2} + c_{y}(1-y)^{3} + e_{y} - \ln(1-x)]/x^{2} \\ &\ln a_{1}^{\ell} - \ln a_{1}^{S} = -\Delta G_{fus,1}^{\circ}/RT = -e_{x} \\ &\ln a_{1}^{\ell} = \ln x + b_{x}(1-x)^{2} + c_{x}(1-x)^{3} \\ &\ln a_{1}^{S} = \ln(1-y) + (b_{y} + \frac{3}{2}c_{y})y^{2} - c_{y}y^{3} \\ &\ln a_{1}^{S} - e_{x} - \ln x = \ln a_{1}^{\ell} - \ln x = (b_{x} + c_{x} - c_{x}x)(1-x)^{2} = \Omega_{1}^{\ell}(1-x)^{2} \\ &\Omega_{1}^{\ell} = [\ln(1-y) + (b_{y} + \frac{3}{2}c_{y})y^{2} - c_{y}y^{3} - e_{x} - \ln x]/(1-x)^{2} \\ &c_{x} = 2(\Omega_{2} - \Omega_{1}), \quad b_{x} = 3\Omega_{1} - 2\Omega_{2} + 2(\Omega_{2} - \Omega_{1})x = \Omega_{1} - c_{x}(1-x) \\ &\frac{d\Omega_{2}^{\ell}}{d_{y}} = \left[\frac{1}{y} - 2 b_{y}(1-y) - 3c_{y}(1-y)^{2}\right]/x^{2} \\ &\frac{d\Omega_{1}^{\ell}}{d_{y}} = \left[\frac{1}{1-y} + (2b_{y} + 2c_{y})y - 3c_{y}y^{2}\right]/(1-x)^{2} \\ \end{split}$$

cont'd.

$$\begin{array}{rcl} \frac{dx_{x}^{2}}{dx} &=& \frac{2\ln(1-x)}{x^{3}} + \frac{1}{x^{2}(1-x)} &-& \frac{2\left[x^{2}\partial_{y}^{2}+\ln(1-x)\right]}{x^{3}} &=& \frac{1}{x^{2}(1-x)} - \frac{2h_{x}^{2}}{x} \\ \frac{d\Omega_{x}^{2}}{dx} &=& -\frac{1}{x(1-x)^{2}} - & \frac{2\ln x}{(1-x)^{3}} + & \frac{\left[2(1-x)^{2}\Omega_{y}^{2}+\ln x\right]}{(1-x)^{3}} &=& -\frac{1}{x(1-x)^{2}} + & \frac{2\Omega_{1}^{2}}{(1-x)} \\ \frac{d\Omega_{2}}{dx} &=& \frac{1}{x(1-x)^{2}} - & \frac{d\Omega_{2}}{(1-x)^{2}} &=& \frac{1}{x^{2}} \\ \frac{d\Omega_{2}}{dx} &=& \frac{1}{x^{2}} - & \frac{d\Omega_{2}}{dx} \\ \frac{d\Omega_{1}}{dx} &=& \frac{1}{x^{2}} - & \frac{d\Omega_{2}}{dx} \\ \frac{d\Omega_{2}}{dx} &=& \frac{d\Omega_{1}}{x^{2}} - & (1-x)\frac{d\alpha_{x}}{dx} = & \frac{d\Omega_{1}}{dx} - & 2(1-x)\left[\frac{d\Omega_{2}}{dx} - & \frac{d\Omega_{1}}{dx}\right] \\ \frac{d\Omega_{2}^{2}}{d(\Delta S_{1}^{2}, 2/R)} &=& \frac{T^{-T}_{R}_{MP2}}{Tx^{2}} , & \frac{dC_{x}}{d(\Delta S_{1}^{2}, 2/R)} = & \frac{2(T^{-T}_{MP2})}{Tx^{2}} , & \frac{db_{x}}{d(\Delta S_{1}^{2}/R)} = & \frac{(2x-2)(T^{-T}_{MP2})}{Tx^{2}} \\ \end{array}$$

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- NOTE 1: After step 6, steps 7, 8 and 9 can be carried out in any order.
- NOTE 2: After step 9, 2X gives $dc_x/d(\Delta S_{f,2}/R)$ which yields $db_x/d(\Delta S_{f,2}/R)$ by multiplication by x-l.
- NOTE 3: If liquid thermodynamic data are available and it is desired to calculate solid data from x,y values, the same program can be used with reversal of 1 and 2 and x and y with the following directions:

Prgm. Ωyxbc

(2) (3) (4)	fLBLA STO7 STO9	fP ≵ S hR↓ hR↓	ST04 ST06 ST08	hR↓ hR↓ hRTN	STO3 STO4	hP ≵ S hR↓	hRTN STO3	hRTN			7 15 19
(5)	fLBLB ÷	STOD RCL9	h 1/x - STO0	RCL6 STO5 fP≵S	X fP ; S hRTN	RCL7 CHS	- STO5	STOO 3	rcl8 X	RCLD 2	44
(6)	fLBLC l fLN X fPZS	fP≵S hx≵y RCL1 - CHS hRTN	STO1 gx ² RCL2 STO7	l STOB ÷ gx ² f-x-	hLSTx STO9 ÷ RCL2	CHS fe f-x- STO8 X	STO2 + RCLB f-x- RCL8	fP≵S fP≵S fe RCL9	hR↓ RCL2 _ CHS	STOA flN RCL1 2 STO6	y 6
(7)	gLBLe + fLBLD	fLN X RCLB	l + RCLB	hLSTx RCL3 RCL5	RCLD X	gx ² ÷ 3	hLSTx RCL ¹ 4 X	RCL5 RCLO	X hRTN 2	RCLO X	115
() /	+ hLSTx X f-x-	X f-x- + RCL9	RCLA RCLA X 2	h l/x RCLA RCLB X	RCL5 h l/x RCL1	CHS X - RCL2	fP≵S 3 RCL2 X X	RCL1 X gx ² h 1/x RCL1	gx ² RCL0 ÷ RCL2	÷ 2 STOE RCL1 X	
(8)	÷ h l/x fLBLE f-x-	CHS - fP≵S RCLA	STOC RCL2 RCLB RCL2	f-x- ÷ RCL1 ÷	RCL8 fP≵s ÷ gx ²	2 STO1 gx ² f-x-	RCLC STOC 1	roll STO2 f-x-	RCL2 hR↓ RCLB 5	hRTN X RCLA	185 209
(9)	− gLBLa fP‡S	X RCLD hRTN	fP\$S -	hRTN hLSTx	<u>.</u>	f₽ <mark></mark> \$	RCL1	gx ²	÷.	CHS	221
RP	0	1	2	3	4	5	6	7	8	9	
	by	$\frac{d\Omega_{l}}{d_{x}}$	$\frac{d\Omega_2}{d_x}$ -A	.H ₂ /R -	∙∆S _l /R	cy	bh	b ^y s	e_{h}^{y}	e_s^y	
RS	0	1	2 3	3 4	5	6	7	8	9		
	$b_y + \frac{3}{2} c_y$	х	l-x ∆H _l	/R AS _l /	′R −e _y	b _x	c _x	Ω_{l}	Ω2		
		A		В	and a second	С	D	E		I	
		У		1 - y	dΩ/đ _x	,dΩ2/db	y T	$\frac{d\Omega_{l}}{dy}$		$\frac{d\Omega_2}{dy}$	

If it is desired to determine the effect upon b_x and c_x of the liquid phase of errors in b_y and c_y of the solid phase, Prgm. bc calculates dc_x/db_y , dc_x/dc_y , db_x/db_y , and db_x/db_y .

$$\begin{aligned} \frac{dc_x}{db_y} &= 2\left(\frac{d\Omega_2}{db_y} - \frac{d\Omega_1}{db_y}\right) = 2\left[\frac{(1-y)^2}{x^2} - \frac{y^2}{(1-x)^2}\right] \\ \frac{dc_x}{dc_y} &= 2\left(\frac{d\Omega_2}{dc_y} - \frac{d\Omega_1}{dc_y}\right) = 2\left[\frac{(1-y)^3}{x^2} - (\frac{3}{2}-y)\frac{y^2}{(1-x)^2}\right] \\ \frac{db_x}{db_y} &= \frac{d\Omega_1}{db_y} - (1-x)\frac{dc_x}{db_y} = \frac{d\Omega_1}{db_y} - 2(1-x)\left[\frac{d\Omega_2}{db_y} - \frac{d\Omega_1}{db_y}\right] \\ &= \frac{y^2}{(1-x)^2} - 2(1-x)\left[\frac{(1-y)^2}{x^2} - \frac{y^2}{(1-x)^2}\right] \\ \frac{db_x}{dc_y} &= \frac{d\Omega_1}{dc_y} - (1-x)\frac{dc_x}{dc_y} = \frac{d\Omega_1}{dc_y} - 2(1-x)\left[\frac{d\Omega_2}{dc_y} - \frac{d\Omega_1}{dc_y}\right] \\ &= (\frac{3}{2}-y)\frac{y^2}{(1-x)^2} - 2(1-x)\left[\frac{(1-y)^3}{x^2} - (\frac{3}{2}-y)\frac{y^2}{(1-x)^2}\right] \end{aligned}$$

Directi	lons:			Displ	av							
(l) Ir	(1) Insert tape bc											
(2) y	(2) $y \uparrow x A$ dc_x/db_y dc_x/dc_y											
	db _x /db _y											
	$\frac{d\mathbf{b}_{\mathbf{x}}}{d\mathbf{b}_{\mathbf{y}}}$ $\frac{d\mathbf{b}_{\mathbf{x}}}{d\mathbf{b}_{\mathbf{y}}}$											
Prgm. 1	oc					na an amhainn an sna mar an ag a' annan Mar Marainn - a sang ann annag ann 2000ir			alanan ay an baanay ay ay ay ah 1990 - Tan ay ah ah			
fLBLA	STOA	l	-	CHS	STOB	hR↓	STOC	1	-			
CHS	STOD	RCLA	÷	gx^2	STO9	RCLC	RCLB	*	gx ²			
STO8		2	Х	ST02	f-x-	RCL9	RCLD	Х	RCL8			
l	•	5	RCLC	-	Х	STO7		2	Х			
STO3	f-x-	RCL8	1	RCLA	-	RCL2	Х		ST00			
f-x-	RCL7	1	RCLA	NAMES.	RCL3	X		STOl	hRTN	60		

cont'd.

0	1	2	3	7	8	9
b _x b _y	$\frac{db_x}{dc_y}$	$\frac{dc_y}{db_y}$	$\frac{\mathrm{d} \mathbf{c}_{\mathbf{x}}}{\mathrm{d} \mathbf{c}_{\mathbf{y}}} \left(\frac{3}{2} - \mathbf{y}\right)$	$\left(\frac{y}{1-x}\right)^2$	$\left(\frac{y}{1-x}\right)^2$	$\left(\frac{1-y}{x}\right)^2$
		A	B	С	D	
		X.	1X	У	1 - y	

There is no general procedure for applying the various programs that have been described for treating a phase diagram. Liquidus data usually would be treated first by Prgm. Ωxb which neglects the effect of solid solubility to obtain a rough idea of the partial molal enthalpies and entropies in the liquid phase. Neglect of solid solubility or experimental error in the data can yield values substantially in error and the initial values may have to be somewhat modified to be within the range of reasonable values indicated by various theoretical models. Program Im can be used to calculate the boundaries of the solidus and liquidus from the selected thermodynamic values. There are, in general, an infinite combination of thermodynamic quantities that will fit a set of experimental data within experimental error. Comparison of the results from Prgm. Im with the original experimental data will indicate what changes in the thermodynamic data might be necessary to give a better fit and still be in the range of reasonable entropy and enthalpy values. Program yx-bc is of use of obtain Gibbs energy values from boundary values at a given temperature or Prgm. Dyxbc will yield the molefraction dependence of the thermodynamic properties in one phase knowing the properties in the saturating phase. This last program can be of value in adjusting the thermodynamic data through analysis of the effect of error in various quantities upon the phase boundaries and the effect of error in one thermodynamic quantity upon related thermodynamic For intermediate phases with relatively limited solid solubility, values. the next program, Prgm. SP, is often of value. Because of the experimental error in most observations and the resulting wide spread of thermodynamic quantities that could be used to fit the data, it is important that this fitting process not be carried out in a mechanical manner. Considerable judgment using theoretical models with useful predictive value is essential to narrow down the range of thermodynamic values consistent with the experimental data.

Solubility Product Calculations for Liquidus of Intermediate Phases

$$\begin{array}{l} \mbox{Combination of the equations $\Delta G^{\circ}/RT = (\Delta G^{\circ} - \Delta H^{\circ}_{298})/RT + \Delta H^{\circ}_{298}/RT$ for $A(s) + nB(s) = AB_{n}(s), A(s) = A(\ell), and nB(s) = nB(\ell)$ gives for $AB_{n}(s) = A(\ell) + nB(\ell), $-\Delta G^{\circ}/RT = -\Delta H^{\circ}_{298}/RT - (\Delta G^{\circ} - \Delta H^{\circ}_{298})/RT$ = 1nK = $1nx_{A} + n1nx_{B} + 1n\gamma_{A} + n1n\gamma_{B} = 1nx_{A} + n1nx_{B} - (b_{s} - b_{h}/T)(1 - x_{A})^{2} - (c_{s} - c_{h}/T)(1 - x_{A})^{3} \\ - n[b_{s} + \frac{3}{2}c_{s} - (b_{h} + \frac{3}{2}c_{h})/T]x_{A}^{2} + n(c_{s} - c_{h}/T)x_{A}^{3}. \\ T = \frac{\Delta H_{298}/R + [b_{h} + c_{h}(1 - x_{A})](1 - x_{A})^{2} + n[b_{h} + \frac{3}{2}c_{h} - c_{h}x_{A}]x_{A}^{2} \\ - \left(\frac{\Delta G - \Delta H^{2}_{298}}{RT}\right) - 1nx_{A} - nln(1 - x_{A}) + [b_{s} + c_{s}(1 - x_{A})](1 - x_{A})^{2} + n(b_{s} + \frac{3}{2}c_{s} - c_{s}x_{A})x_{A}^{2} \end{array}$$

$$-\Delta G^{\circ}/RT = \ln x_{A} + n\ln(1-x_{A}) + \frac{[b_{h}+c_{h}(1-x_{A})](1-x_{A})^{2} + n[b_{h}+\frac{3}{2}c_{h}-c_{h}x_{A}]x_{A}^{2}}{T}$$
$$- [b_{s}+c_{s}(1-x_{A})](1-x_{A})^{2} - n[b_{s}+\frac{3}{2}c_{s}-c_{s}x_{A}]x_{A}^{2}$$

Program SP can be used to calculate the liquidus curve when $\Delta H_{298}^o/R$ is known or to calculate $\Delta H_{298}^o/R$ when the liquidus curve is known.

Directions:

(1)	Insert Pr	gm. SP									
(2)) b _h ↑ c _h ↑ b _s ↑ c _s E) ΔH ₂₉₈ /R STOA, -(ΔG-ΔH ₂₉₈)/RT STOB n STOE, 273.15 STO3										
(3)	ΔH ^o 298/R	STOA, $-(\Delta G - \Delta H_{298}^{\circ})$)/RT STOB	n STOE, 273.	15 STO3						
(4)	$x_A \land T$										
(5)	$R/S \rightarrow t$										
(6)	хд ↑ Т В	→ -∆G°/RT									
(7)	R/S → ∆H°	/R									
		nangen standen eine seite der seiten an der stellen in der Stellen ander stellen andere andere stelle Stellen andere stelle	Karlanda ya Mita a Managa Mita Muna ata inter ata ya managa ata inter ata ya managa ata ya managa ata ya managa	anna agus aine an san san san san san san san san san							
0		10	~								
А		В	С	D	Е						
А Дн°298	g/R	Β -(ΔG°-ΔH ^o 298)/RT		D l-x _A	E n						
	g/R										
	g/R										
∆H°298	₉ /R	-(ΔG°-ΔH ₂₉₈)/RT	xA	l-x _A	n						

 $c_{\rm S}$

cont'd.

b_s

RS

harangener a a sanganga (*										·
Prgm. S	SP									
fLBLE	f₽ ≵ S	STOl	hR↓	STOO	hR↓	f₽ ≵ S	STOL	hR↓	ST00	
hRTN										11
fLBLA	STOC	1	-	CHS	STOD	fGSBO	RCLA	+	ST02	
fP ∠ S	fGSBO	f₽ <mark>≵</mark> S	RCLB	+	RCLD	fLN	RCLE	Х	-	
RCLC	fLN		RCL2	hx≵y	*	hRTN	RCL3	-	hRTN	
RCL3	-	hRTN								44
fLBL0	1	٠	5	RCLC	-	RCL1	Х	RCLO	+	
RCLE	Х	RCLC	gx ²	Х	RCLD	RCL1	Х	RCLO	+	
RCLD	gx ²	Х	+	hRTN						69
fLBLB	STO4	hR↓	STOC	1		CHS	STOD	fLN	RCLE	
Х	RCLC	fLN	+	fGSBO	RCL4	÷	+	f₽≵S	fGSBO)
fP≵S		hRTN	RCLB		RCL4	Х	CHS	hRTN		98

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Chapter IV

Regular Solution Calculations

The previous chapters have frequently referred to regular solution equations and their application to thermodynamic calculations. (12,19,20) In Chapter III, it was pointed out that the partial molal Gibbs energies expressed in terms of volume fractions as composition variables can be conveniently expanded in terms of mole fraction truncated at the cubic term. In its simplest form, the regular solution theory is a pair-interaction model with interactions between unlike atoms or molecules being taken as the geometric mean of the interactions between like atoms of molecules and excess entropy contributions comes only from the temperature coefficients of the volume and energy of sublimation terms. For most actual systems, the geometric mean assumption must be modified and additional entropy contributions must be considered. For metallic atoms, there is also the question of the appropriate valence state to use for evaluating the cohesive energy(7).

The programs tabulated below can use for ΔH either the enthalpy of sublimation to the ground electronic state or to any valence state in obtaining the first approximations to the solubility parameters. Subsequent approximations depend upon the types of atoms being mixed and an appraisal of the types of chemical interactions that can take place. Thus strain contributions due to size disparity can add positive contributions to the Gibbs energy of mixing for solids with equivalent lattice positions and can add negative contributions for appropriate size differences for structures such as those of the Laves phases, liquids, and other systems with nonequivalent sites that achieve better space filling with appropriate size The very substantial acid-base interactions that can take disparities. place when atoms with non-bonding electrons are mixed with atoms with unoccupied low energy orbitals must also be considered.(8) However, the simplest form of internal pressure calculation is the starting point for For this level of approximation, the energy of vaporization most systems. is taken as $\Delta H-RT_i$ where T_i is the average temperature of the range of interest. $T\Omega_{l} = T\ln\gamma_{l}/(l-x_{l})^{2} = V_{l}(\delta_{n}-\delta_{l})^{2}/R = V_{l}\left\{\left[\left(\frac{\Delta H_{n}}{R} - T_{i}\right)/V_{n}\right]^{\frac{1}{2}} - \left[\left(\frac{\Delta H_{l}}{R} - T_{i}\right)/V_{l}\right]^{\frac{1}{2}}\right\}^{2}$ $T\Omega_n = T \ln \gamma_n / x_1^2 = V_n (\delta_n - \delta_1)^2 / R$ $T_{critical} = [2V_{n}(\delta_{n}-\delta_{l})^{2}/R]/[1+(V_{n}/V_{l})^{\frac{1}{2}}]^{2}$

Program RS can store data for eight binary systems of component 1.

Dire	ctions:	Display
(1)	Insert Prgm. RS	Dat blay
(2)	R ↑ T _i ↑ ΔH _l ↑ V _l A	$\left[\left(\Delta H_{1}/R-T_{i}\right)/V_{1}\right]^{\frac{1}{2}}$
(3)	В	2
(4)	$\Delta H_n \uparrow V_n C$; repeat for all n	$V_{l}(\delta_{n}-\delta_{l})^{2}/R$, $V_{n}(\delta_{n}-\delta_{l})^{2}/R$
(5)	up to $n = 9$ T \uparrow n D	$\Omega_1 = \ln \gamma_1 / x_n^2$, $\Omega_n = \ln \gamma_n / x_1^2$
(6)	n E	T _{critical} for 1-n system
NOTE	1: Step 3 is required before a new	w set of data in step 4.

NOTE 2: Steps 5 and 6 can be used in any order. If it is desired to add more values in step 4 after step 5 or 6 has been used, carry out step 6 for last value of n used in step 4. Then key f b and step 4 will accept values for n+1, n+2, etc. up to the maximum of 9.

Prgm	. RS	, <u>, , , , , , , , , , , , , , , , , , </u>			,., ,, ,, ,, ,, ,, ,, ,, ,, ,, ,, ,, ,,			<u>,</u>	~ <u>~</u>	*	
(2)	fLBLA	STOA	hR↓	STOB	hR↓	STOC	hR↓	STOE	RCLB	hx≠y	
	÷	RCLC	-	RCLA	÷	$f\sqrt{x}$	STOD	hRTN			18
(3)	fLBLB	2	hSTI	hRTN							22
(4)	fLBLC	ST00	hR↓	STOL	RCLE	÷	RCLC		RCLO	÷.	
	$f\sqrt{x}$	RCLD		gx^2	个	STO(i) fa	RCLO	STO(i)	Х	
	fb	hx≵y	RCLA	Х	f-x-	hR↓	hRTN				49
(5)	fLBLD	hSTI	hR↓	RCL(i)	hx≵y	÷	↑	↑	fa	RCL(i)	
	Х	hx≵y	RCLA	Х	f-x-	hR↓	hRTN				66
(6)	fLBLE	hSTI	RCL(i)	2	X	fa	RCL(i) X	RCL(i)	RCLA	
	÷	f√x	1	+	gx^2	÷	hRTN				83
	gLBLfa		1	0	+	hSTI	hR↓	hRTN			91
	gLBLfb	hRCI	9	-	hSTI	hR↓	hRTN				98
			~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~				an gan, a a sa gan ya gan da		an a		
	А	В	С	/ A 11	D /pm.	$\frac{1}{2}$	Ε	I			
	Vl	∆Hl	T.		<u>l/R-Ti</u> V;	-) ²	R	Index			

	<u>_</u>	-de,	1 (	V _i /			
P	0	1	2	3	• • • • • • • • • • • •	8	9
	Vn	ΔHn	(δ ₁ -δ ₂ ) ² /R	$(\delta_1 - \delta_3)^{t}$	² /R(8	$(1-\delta_8)^2/R$	$(\delta_1 - \delta_9)^2/R$
S	A	2	3.		•••••	8	9
		V ₂	v ₃ .	•••••	• • • • • • • • • • • • • • • • • •	v ₈	v ₉

### Use of Regular Solution Equations for Fitting of Data

When there is substantial solid and liquid solubility and Prgm. Im of Chapter III is required, an exact analytical solution cannot be given, but quite accurate fitting of the results can be obtained in terms of a polynomial fit of the deviation from simple limiting equations.

If the solid and liquid solubilities are small, the liquidus of component A of an A-B binary system is given by  $\ln x_A^L = \alpha + a_{-1}T^{-1}$  where  $\alpha = \Delta S_f^o/R + b_s^L + c_s^L$  and  $a_{-1} = -\Delta H_f^o/R - b_h^L - c_h^L$ 

and all quantities refer to component A. The solidus can be expressed as

$$\ln x_{B}^{S} = \alpha^{S} + a_{-1}^{S} T^{-1}$$

where

$$\alpha^{S} = -\Delta S_{f}^{o}/R + b_{s}^{S} + c_{s}^{S}$$
 and  $a_{-1}^{S} = \Delta H_{f}^{o}/R - b_{h}^{S} - c_{h}^{S}$ 

and all quantities refer to component B.

Actual data are fitted to the equation  $\ln x = a_{-1}T^{-1} + \Delta$ , where  $\Delta$  is a power series in T. It is sometimes useful to fit  $\ln \frac{x}{T_{f}-T} = a_{-1}T^{-1} + \Delta$ . When the liquidus or solidus is of very limited extent near the melting point, a simpler equation can be used. The limiting liquidus near the melting point of B becomes

$$\frac{\mathbf{x}_{A}^{L}}{(\mathbf{T}_{f}-\mathbf{T})} = \frac{\Delta \mathbf{S}_{f}^{\circ}/\mathbf{R}}{\mathbf{T}(1-\mathbf{x}_{A}^{S}/\mathbf{x}_{A}^{L})}$$

from Eq.(19-14) of ref.(12) which reduces to

$$\frac{x_{A}^{L}}{(T_{f}-T)} = \frac{\Delta S_{f}^{\circ}/R}{T}$$

when solid solubility is small.

The various modifications of these limiting equations can be used to obtain accurate least square fits for liquidus or solidus boundaries to a polynomial with a minimum number of terms. Prg. LSA is designed to handle the variety of situations that might arise by evaluating the constants of  $f(x,T) = a_1/T + a_0 + a_1T + a_2T^2$ , where f(x,T) may be  $\ln[x/(T_f-T)]$ ,  $\ln x$ , or  $x/(T_f-T)$ .  $a_{-1}$  is fixed as indicated below by the type of equilibrium and  $\Delta = f(x,T) - a_{-1}/T$  is fit to a power series in T or under some circumstances in  $T_f-T$ .

Prgm. LSA for  $f(x,T) = \frac{a_{-1}}{T} + \Delta$ ,  $\Delta = a_0 + a_1T + a_2T^2$ Display (1) Insert tape LSA(2)Store a_1 in B a_1 (3) fa to clear registers a_7 Modify prgm. to fit  $f(x,T) = x/(T_f-T)$  or lnx in place of ln  $\frac{x}{2800-T}$ (4)  $T_1^2 \Delta$ (5) х, ТрВ If any of T values are evenly spaced, key xA after 2nd of regularly spaced values if  $10^{-2}$  has been replaced by proper I. (6) R/S an (7)SST SST aj, ag (8)ΤD Х E after 2nd of regularly spaced T x (9) C, SST  $a_1$ ,  $a_0'$  for  $T_f-T$  in place of T (10) T fd Х fe for evenly spaced T x Prgm. is set for  $f(x,T) = \ln \frac{x}{2890-T}$ , although usually a simpler NOTE 1: function is used. NOTE 2: For  $f(x,T) = \ln x$ , remove fGSB 1- in 3rd line (steps 22-3) and fGSB1 + near end of LBL D (steps 132-3 minus 2) and if  $T_{f}$ -T power series be used, remove hRt fLN + of gLBLd, step 192. NOTE 3: If x is keyed in as  $10^{n}$ x, add EEX n ÷ in 3rd line (after hRt of step 20) NOTE 4: For melting point of metal other than Mo, replace 2890 by  $T_{f}$  after fLBL1, fLBLC, and end of gLBLd. For even spaced temp, at intervals I other than 100°, replace EEX 2 NOTE 5: by I after fLBLA, fLBLE, and gLBLd. If x is to be used instead of lnx, remove fLN in 3rd line (step 21), NOTE 6: and remove ge^x at end of fLBLE and end of gLBLe. For an A-B system with  $\ln \gamma_A = bx_B^2 + cx_B^3$  and  $\ln_B = (b + \frac{3}{2}c)x_A^2 - cx_A^3$ , there are the following combinations of f(x,T) and  $a_{-1}$ . NOTE 7:

Solid-Liquid Equilibrium A liquidus, low $x_A^L$ A solidus, low $x_A^L$ A liquidus, high $x_A^L$ A solidus, high $x_A^L$ B liquidus, high $x_B^L$ B solidus, high $x_B^L$	$\frac{f(x)}{\ln x}$ $\ln x$ $\ln x_{B}/(T_{f})$ $\ln x_{B} \text{ or } \ln[x_{B}/x_{A}/(T_{f})]$ $\ln x_{A} \text{ or } \ln[x_{A}/x_{A}/(T_{f})]$	^ζ A - ^ζ B ^ζ ,A ^{-T} ) ^ζ (T _f ,A-T)] ζ	$ \begin{array}{c} -\Delta H_{f,A}^{a} - b^{L} - c^{L} \\ \Delta H_{f,B}^{} - b^{S} - c^{S} \\ \Delta S_{f,A} \\ \Delta H_{f,B}^{-} (b^{S} - b^{L}) - \frac{1}{2} (c^{S} - c^{L}) \\ \Delta S_{f,B} \\ \Delta H_{f,A}^{} - (b^{S} - b^{L}) - (c^{S} - c^{L}) \end{array} $
<u>Solid-Solid Equilibrium</u> A solidus	ln x	(T)	$-b^{\mathrm{S}}-\frac{1}{2}c^{\mathrm{S}}$
B solidus	ln x	2	_b ^S _c ^S
Prgm. LS∆			
fLBLARCLEEEX $\Sigma^+$ hLSTxhR^1fPZS $\Sigma^+$ hLSTxRCL4RCL8fPZSgx2RCL9 $\div$ RCL8-ST03-ST05RCL1RCL4RCL6Xgx2-RCL8:ST0CRCL8RCL7XRCL6-ST0AhRTNfLBLEhR^1EEX+fGSB1+fLBL128gLBLb $\uparrow$ $\uparrow$ +hRTNfLBLC282XRCLC	$fP_{\star}$ 'S $fCLREG$ 2+ $fLN$ $fGSB1$ $RCLE$ X $STO2$ $hR_{\star}$ $RCL7$ - $hx_{\star}y$ $STO8$ $RCL2$ $RCL4$ $hR^{\star}$ - $RCL5$ XX $RCLE$ $RCLC$ X $RCLC$ X $RCLC$ $RCLD$ 2+ $ge^{x}$ $hRTN$ 90 $\uparrow$ $RCLD$ 90+ $CHS$ 2+	STOB hRTN fLBLB STOE - RCLB STO+1 fP _₹ S STO0 RCL9 hLSTx RCL6 RCL7 gx ² RCL7 X STOE RCL8 RCL3 hRC1 hx _₹ y - + RCL9 fLBLD fb hR↑ - X RCLC fb STO3 STO4 hRTN gLBLe ↑	$\uparrow \uparrow gx^{2}$ $RCLE \div -$ $hRTN 35$ $\div STO4 RCL6$ $RCL9 \div X$ $RCL9 \div RCL5$ $hx \neq y - hSTI$ $RCL5 X RCL3$ $X - hx \neq y$ $RCL3 \div STOD$ $\div RCL4 hx \neq y$ $RCL3 \div STOD$ $\div RCL4 hx \neq y$ $RCL3 + hRTN 1^{20}$ $RCLB hR \uparrow \div 1^{20}$ $RCLB hR \uparrow \pm 1^{35}$ $fLN hRTN 1^{44}$ $+ X RCL4$ $156$ $hR \uparrow RCLD X$ $RCL3 + 174$
gLBLd hR↑ EEX X RCL4 + 8 9 0 hRTN	2 + X RCL3 hR↑ -	gLBLe ↑ + hR↑ RCLB hx≠y	↑ ↑ RCLD fLN + 2 ÷ + ge ^x 205

	0	1	2	3	4	5		7	8	9	
RS		ΣτΔ			ΣT2	$\Sigma T^{\lambda_{\mu}}$	ΣΤ	ΣT2	$\Sigma T^3$	i	
then			2		$\frac{1}{i}\Sigma\Delta$	r			đ		
RP	ΣΔ		$\Sigma T^2 \Delta$	S	1	Т			ų		
				$a'_1$	a'						
					4	5	6	7	8	9	
RP					ΣΔ	$\Sigma \Delta_5$	$\Sigma T^2$	$\Sigma T^{\lambda}$	$\Sigma T^2 \Delta$	i	
then											
RS											
	A	В		С	D	E		Ι			
		a_1	~			Т					
	ao			al	a2	t		u			

A fifth chapter entitled "Philosophy of Critical Evaluation and Compilation" has been written which describes the compilation procedures and use of some of the programs described in the earlier chapters and which particularily reviews some of the bonding models used to predict thermodynamic values for systems lacking data. Because of the length of the present report, the decision was made to issue the fifth chapter as a separate report (LBL-7666) and the present summary of calculation programs will be ended with a program used to estimate enthalpies of formation using the Rittner model for gaseous halides with a large ionic component.

### Bonding Energies of Gaseous Halides Using the Rittner Model

Rittner (24) has used an ionic model including polarizability to calculate bonding energies of gaseous halides. There are a number of difficulties in the model having to do with the proper treatment of overlap repulsion, the appropriate value of the polarizability to be used, and the contribution of higher order terms. In spite of the limitations, Cubicciotti (25,27) and Hildenbrand (26) have reported success in using the model for compounds of highly charged cations and criteria have been developed to predict range of applicability of the model.

It is of interest to determine how far the model can be extended among the gaseous transition metal halides, for which thermodynamic data are very scanty. It would be useful to compare the calculated values with the available experimental data to determine if the model or some modification of the model might yield useful data. The HP-67 calculator program given below follows the calculations of Cubicciotti, Hildenbrand, <u>et al.</u> (27). Enthalpies of formation or dissociation energies to the gaseous ions can be readily calculated for MX, linear MX₂, planar MX₃, and tetrahedral MX₄ given a value of the cation-anion distance and the polarizabilities. These molecules will be described as MX_n with x=1 to 4. The operation of the program will be illustrated using typical parameters for Zr and Mo.

H₀^o/R = 167 103[
$$a_n/r + \alpha b_n/r^4$$
] for  
1) Mⁿ⁺(g) + nX⁻(g) = MX_n(g)

(

where  $\alpha$  is the anion polarizability in A³ for the symmetrical molecules and is the sum of the anion and cation polarizabilities for MX molecules. The constant 167 103, given by  $e^2 N_0 / 10^{-8} R = e^2 / 10^{-8} k = (4.80324)^2 10^4 / 1.38065$ , yields values of  $\Delta H_0^0 / R$  in Kelvin units when r is given in Å. The values of  $a_n$  and  $b_n$  are (27):

$$\begin{array}{rcl} a_{1} & = & -\frac{10}{11} & & b_{1} & = & -\frac{7}{11} \\ a_{2} & = & -\frac{10}{11}(3.5) & & b_{2} & = & -\frac{7}{11}(\frac{49}{16}) \\ a_{3} & = & -\frac{10}{11}(9-\sqrt{3}) & & b_{3} & = & -\frac{7}{11}(8.17) \\ a_{4} & = & -\frac{10}{11}(15-3\sqrt{1.5}) & & b_{4} & = & -\frac{7}{11}(16.48) \end{array}$$

where  $\frac{10}{11} = 1 - \frac{1}{p}$  and  $\frac{7}{11} = 1 - \frac{4}{p}$  with p, the power of the overlap repulsion term, taken equal to 11.

The value of  $\Delta H_0^o/R$  for  $M(s) + \frac{n}{2}X_2 = M^{n+}(g) + nX^-(g)$  is given by the sum of  $\Delta H_0^o$  for the following reactions.

$$(2) \qquad M(s) = M(g)$$

(3)  $\frac{n}{2}X_2 = \frac{n}{2}X(g)$ 

where X₂ is taken as gaseous for F₂ and Cl₂ at OK and solid for Br₂ and I₂.

$$(l_{4}) \qquad nX(g) + ne^{-}(g) = nX^{-}(g)$$

(5)  $M(g) = M^{n+}(g) + ne^{-}(g).$ 

For reaction (2),  $\Delta H_0^0/R = 78970\pm250K$  was used (29) for Mo and 72560 $\pm$ 500K was used (28) for Zr. For reaction (3),  $\Delta H_0^0/R$  was taken (30-35) as 9 240 $\pm$ 70 for F, 14 386 $\pm$ 1 for Cl, 14 185 $\pm$ 5 for Br, and 12 889 $\pm$ 2 K for I. For reaction (4),  $\Delta H_0^0/R$  was taken (36,37) as -39 440 $\pm$ 20 for F, -41 930 $\pm$ 30 for Cl, -39 030 $\pm$ 30 for Br, and -35 550 $\pm$ 30K for I. The ionization potentials in cm⁻¹ from Moore (38) for

are used to calculate  $\Delta H_0^{\circ}/R$  for reaction(5). The sum of  $\Delta H_0^{\circ}/R$  values for reactions (1) to (5) yields  $\Delta H_0^{\circ}/R$  for the formation from the elements.

(6) 
$$M(s) + \frac{n}{2} X_{2}(g) = MX_{n}(g)$$
$$\Delta H_{298}^{\circ}/R = \Delta H_{0}^{\circ}/R + (H_{298}^{\circ}-H_{0}^{\circ})/R_{MX_{n}} - (H_{298}^{\circ}-H_{0}^{\circ})/R_{M} - \frac{n}{2}(H_{298}^{\circ}-H_{0}^{\circ})/R_{X_{2}}.$$
For M(s),  $(H_{298}^{\circ}-H_{0}^{\circ})/R = 552K$  for Mo and 665K for Zr. For gaseous X₂,  
 $\frac{1}{2}(H_{298}^{\circ}-H_{0}^{\circ})/R = 531$  for F₂ and 552 for Cl₂. For Br₂( $\ell$ ) at 298K,  $\frac{1}{2}(H_{298}^{\circ}-H_{0}^{\circ})/R = 1474K$  and for I₂(s),  $\frac{1}{2}(H_{298}^{\circ}-H_{0}^{\circ})/R = 794K$ . Values of  $(H_{298}^{\circ}-H_{0}^{\circ})/R$  for MX_n(g) are tabulated in Table A.

The calculation of  $\Delta H_{298}^{\circ}/R$  for reaction(6) can be expressed as  $\Delta H_{298}^{\circ}/R = U/R + M_n + nX_x + (H_{298}^{\circ}-H_0^{\circ})/R_{MX_n}$ , where  $\Delta H_0^{\circ}/R$  for reaction(1) is expressed as  $U/R = A_n/r + \alpha B_n/r^4$  with  $A_n = 167 \ 103 \ a_n$  and  $B_n = 167 \ 103 \ b_n$ ,  $M_n$  is the sum of the  $\Delta H_0^{\circ}/R$  values of reactions(2) and (5) minus  $(H_{298}^{\circ}-H_0^{\circ})/R$  for M(s), and  $X_x$  is the sum of  $\Delta H_0^{\circ}/R$  values of reactions(3) and (4) minus  $(H_{298}^{\circ}-H_0^{\circ})/R$  for  $\frac{1}{2}X_2$ . The  $M_n$  values are calculated by step 5 of the program. The other values, along with  $\alpha_x$  values, are tabulated as follows:

n	A _n	B _n		Xx	$\alpha, \hat{A}^3$
1	-151 912	- 53 169	F	-30 730	1.2
2	-531 691	-325 660	Cl	-28 100	3.0
3	-1 104 087	-868 800	Br	-26 320	4.5
4	-1 872 429	-1 752 500	Ι	-23 450	6.42

The  $\alpha$  value for I is from Cubicciotti, Hildenbrand, <u>et al.(27)</u>. The other values are from Dalgarno (39). The calculation is carried out in two stages. After storing A_n and B_n, U/R is calculated from the M-X distance r by U/R = A_n/r +  $\alpha$ B_n/r⁴. Then the value of U/R is combined with (H²₂₉₈-H⁰₀)/R of MX_n, to calculate  $\Delta$ H²₂₉₈/R of reaction(6). Constants are stored in registers or are inserted with a data card as indicated.

RP	0 r	l A _l	2 A ₂	3 A ₃	կ А <u>կ</u>	5 1.4388	б Х _F	7 X ₍	8 Cl X _{Br}	. 9 . X _I	
RS		l B _l		2 ^B 2		3 ³ 3	կ Bչլ	б М _l	7 M ₂	8 M ₃	9 M _{l4}
		A X _x		в М _п		n An	$^{\rm D}_{\alpha_{\rm x}}{}^{\rm B}_{\rm n}$	E R	I 6-9 n n+10 n+15		
TABL	EA		Valu	es of	(H ² 98	3 <b>-</b> H8)/R	in K	94.4949-00-04 <u>-04-04-04-04</u>	****		
n	MoF	'n	MoCl	n N	10Br _n	Mol _n	ZrFn	Zr(	Cl _n ZrB	Brų ZrI)	4
1	107	'9	1141	. ]	L176	1221	1092	11	32 117	'2 121 <u>'</u>	3
2	162	24	1948	2	2080	2200	1423	178	83 187	1842	2
3	176	58	2104	. 4	2395	2570	1817	21	19 232	21 2930	С
4	227	<b>7</b> 4	2678	2	2987	.3168	2278	27	13 300	)1 318	5

Directions:

(1)	Insert Prgm. I	
	Insert Data Tape I	Display
	α ↑ n ↑ r ↑ 6 for F, 7 for Cl	
	8 for Br, or 9 for I A	U/R of MX _n (g)
(3a)	If calculation is to be repeated	
	for different n with same X,	
	α↑n↑r B	U/R of MX _n (g)
(4)	With U/R displayed, key	
	$(H_{298}-H_0)/R$ of $MX_n$ R/S	$\Delta H_{298}^2/R$ of $MX_n(g)$
(5)	To insert data for different M,	
	$\Delta H_0^{R}$ subl. of M $\uparrow$ (H ₂₉₈ -H ₀ )/R of M(s) E	1.44
	$IP_4 \uparrow IP_3 \uparrow IP_2 \uparrow IP_1 \uparrow R/S$	M)4
(6)	fa to multiply any value by R	
(7)	for the drivide entry traduce by P	

(7) fe to divide any value by R

NOTE		is $\alpha_x$ of bion and c				orn=	l where	sum of	$\alpha$ valu	aes of	
NOTE	2: St	tep 2 Leav in go dire	res reg	isters	prepare						è
Prgm.	. I										
(3)	fLBLA	hSTI	hR↓	RCLi	STOA	hR↓	fLBLB	ST00	hR↓	hSTI	
	hR↓	RCLi	STOC	bR↓	hRCT	1	0	+	hSTI	hR↓	
	RCLi	Х	STOD	hRCI	5	+	hS'TI	RCLi	STOB	RCLO	
	$h\frac{1}{x}$ hRTN	RCLC	hx≑y	X	RCLD	hLSTx	}4	$hy^X$	X	4	
(4)	÷	hRCI	1	5	\$ <b>57%</b>	RCIA	Х	÷	RCLB	+	44
( )	hRTN	and the star				110.1111	21	·	TOTD	,	5 2
(5)	fLBLE	1	9	hSTI	hR↓	~	STOi	fDSZ	STOi	fDSZ	
	STOi	fDSZ	STOi	RCL5	STOO	hRTN					68
	STOX0	hR↓	RCLO	STO+i	fISZ	hR↓	RCL5	Х	RCLO	÷	
	fGSBO	fGSB0	STO+i	RCLi	STOB	hRTN					84
	fLBLO	STO+i	fISZ	hx≓y	RCL5	Х	+	hRTN			92
(6)	gLBLa	RCLE	Х	hRTN							96
(7)	gLBLe	RCLE	<u>*</u>	hR'I'N							100
	RP EG, 6	0 0 -1 7	1 51 912 {		691 9	3 -1 101	+ 087 A	ц -1 872 В		5 1.4388 D	EI
	730	-28 100		320	-23 45	io3	A 30 730	0			8719 0
I	rs - 5	1 3 169 -	2 32560	3 8688	00 -17	) ₄ 752500	6 160802	7 348275	8 6635	10 1	9 201871
(3a) (3a) (3) (3a) (3) (5) (6)	8.2 † 1 3 † 4 · 10 † 1 6.42 † 78970 ·	4 † 1.82 1 † 1.82 † 2.26 † † 2.26 B 3 † 2.66 † 552 E 1 3347 (7)	B -1232 7 A -10 -87599 ↑ 9 A .44, 37	204, 030041, , -52648 74180 ↑	(4) 107 (4) 267 (4) 114 1, (4)	9 R/S 7 8 R/S 6 1 R/S 4 2570 R/	947 52108 6245 's 69249	) 7260 R/S	12018	71,	

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This report was done with support from the United States Energy Research and Development Administration. Any conclusions or opinions expressed in this report represent solely those of the author(s) and not necessarily those of The Regents of the University of California, the Lawrence Berkeley Laboratory or the United States Energy Research and Development Administration.